# GENERALIZED CONNECTION ON $T\left(T^{2} M\right)$ 

## Irena Čomić

## Communicated by Mileva Prvanović


#### Abstract

The geometry of some manifolds fibered over a given manifold $M$ is in the first place characterized by the group of allowable coordinate transformations. For the tangent manifold $T M$ these are given by $x^{i^{\prime}}=x^{i^{\prime}}(x) y^{i^{\prime}}=\frac{\partial x^{i^{i}}}{\partial x^{i}} y^{i}$, $\operatorname{rank}\left[\frac{\partial x^{i^{\prime}}}{\partial x^{i}}\right]=n$, and for the total space of a vector bundle $E \rightarrow M$, we have $x^{i^{\prime}}=x^{i^{\prime}}(x), y^{a^{\prime}}=M_{a}^{a^{\prime}}(x) y^{a}, \operatorname{rank}\left(M_{a}^{a^{\prime}}\right)=m=$ dimension of type fiber.

In the last years R. Miron, Gh. Atanasiu and others examined the Osc ${ }^{k} M$ spaces, $[\mathbf{1 0}],[\mathbf{1 1}],[\mathbf{1 2}]$. Here the case $k=2$ will be investigated. Instead of $\mathrm{Osc}^{2} M$ the notation $T^{2} M$ will be used ( $\mathrm{Osc}^{1} M$ coincides with $T M$ ). Instead of $d$-connection used in $[\mathbf{1 0}],[\mathbf{1 1}]$, $[\mathbf{1 2}]$, we consider here the generalized connection and determine its torsion tensor. As a special case the known $d$-connection is obtained.


1. Adapted basis in $T\left(T^{2} M\right)$. Let $T^{2} M$ be a $3 n$ dimensional $C^{\infty}$ manifold. A point $u \in T^{2} M$ in the local charts $(U, \varphi)$ and $\left(U^{\prime}, \varphi^{\prime}\right)$ has coordinates $\left(x^{i}, y^{i}, z^{i}\right)$ and $\left(x^{i^{\prime}}, y^{i^{\prime}}, z^{i^{\prime}}\right)$ respectively. In $U \cap U^{\prime}$ the allowed coordinate transformations are given by the equations:
(a) $x^{i^{\prime}}=x^{i^{\prime}}(x)$
(b) $y^{i^{\prime}}=\frac{\partial x^{i^{\prime}}}{\partial x^{j}} y^{j}$
(c) $z^{i^{\prime}}=\frac{1}{2} \frac{\partial^{2} x^{i^{\prime}}}{\partial x^{k} \partial x^{j}} y^{k} y^{j}+\frac{\partial x^{i^{\prime}}}{\partial x^{j}} z^{j}$.

If rank $\left[\frac{\partial x^{i^{1}}}{\partial x^{i}}\right]=n$, then the inverse transformation of (1.1) exists:
(a) $x^{i}=x^{i}\left(x^{i^{\prime}}\right)$
(b) $y^{i}=\frac{\partial x^{i}}{\partial x^{j^{\prime}}} j^{j^{\prime}}$
(c) $z^{k}=\frac{1}{2} \frac{\partial^{2} x^{k}}{\partial x^{i^{\prime}} \partial x^{j^{\prime}}} y^{i^{\prime}} y^{j^{\prime}}+\frac{\partial x^{k}}{\partial x^{i^{\prime}}} z^{i^{\prime}}$.
(1.2) (a) and (1.2) (b) are obvious. To obtain (1.2) (c) we start from

$$
\begin{equation*}
\frac{\partial x^{i^{\prime}}}{\partial x^{h}} \frac{\partial x^{k}}{\partial x^{i^{\prime}}}=\delta_{h}^{k} \tag{1.2}
\end{equation*}
$$

From (1.2) it follows:

$$
\frac{\partial^{2} x^{i^{\prime}}}{\partial x^{h} \partial x^{j}} \frac{\partial x^{k}}{\partial x^{i^{\prime}}}+\frac{\partial x^{i^{\prime}}}{\partial x^{h}} \frac{\partial^{2} x^{k}}{\partial x^{i^{\prime}} \partial x^{j^{\prime}}} \frac{\partial x^{j^{\prime}}}{\partial x^{j}}=0 .
$$

The multiplication of the above equation with $y^{h} y^{j}$ gives

$$
\begin{equation*}
\frac{\partial^{2} x^{i^{\prime}}}{\partial x^{h} \partial x^{j}} y^{h} y^{j} \frac{\partial x^{k}}{\partial x^{i^{\prime}}}=-\frac{\partial^{2} x^{k}}{\partial x^{i^{\prime}} \partial x^{j^{\prime}}} y^{i^{\prime}} y^{j^{\prime}} \tag{1.3}
\end{equation*}
$$

If we multiply (1.1) (c) by $\partial x^{k} / \partial x^{i^{\prime}}$, use (1.3) and (1.4) we obtain (1.2) (c).
The identity transformation is a special case of (1.1), namely if we put $x^{i^{\prime}}=x^{i}$ in (1.1) (a), then $y^{i^{\prime}}=y^{i}, z^{i^{\prime}}=z^{i}$ follows.

If in the local chart $\left(U^{\prime \prime}, \varphi^{\prime \prime}\right)$ the point $u$ has coordinates $\left(x^{i^{\prime \prime}}, y^{i^{\prime \prime}}, z^{i^{\prime \prime}}\right)$, then in $\left(U^{\prime \prime} \cap U^{\prime}, \varphi^{\prime \prime}\right)(1.1)$ are valid, if the index $i^{\prime}$ is substituted by $i^{\prime \prime}$ and the indices without ' obtain '. After some calculation it can be obtained, that the connection between $\left(x^{i^{\prime \prime}}, y^{i^{\prime \prime}}, z^{i^{\prime \prime}}\right)$ and $\left(x^{i}, y^{i}, z^{i}\right)$ in $U \cap U^{\prime} \cap U^{\prime \prime}$ is given by (1.1) if the index $i^{\prime}$ is substituted by $i^{\prime \prime}$.

From the above follows:
THEOREM 1.1. The transformations of type (1.1) form a group.
In $T\left(T^{2} M\right)$ the natural bases are:

$$
\begin{equation*}
\bar{B}=\left\{\frac{\partial}{\partial x^{i}}, \frac{\partial}{\partial y^{i}}, \frac{\partial}{\partial z^{i}}\right\} \text { and } \bar{B}^{\prime}=\left\{\frac{\partial}{\partial x^{i^{\prime}}}, \frac{\partial}{\partial y^{i^{\prime}}}, \frac{\partial}{\partial z^{i^{\prime}}}\right\} . \tag{1.4}
\end{equation*}
$$

The bases vectors of $\bar{B}$ and $\bar{B}^{\prime}$ are connected by:

$$
\begin{array}{rlr}
\frac{\partial}{\partial x^{i}} & =\frac{\partial x^{i^{\prime}}}{\partial x^{i}} \frac{\partial}{\partial x^{i^{\prime}}}+\frac{\partial y^{i^{\prime}}}{\partial x^{i}} \frac{\partial}{\partial y^{i^{\prime}}}+\frac{\partial z^{i^{\prime}}}{\partial x^{i}} \frac{\partial}{\partial z^{i^{\prime}}}, \\
\frac{\partial}{\partial y^{i}} & \frac{\partial y^{i^{\prime}}}{\partial y^{i}} \frac{\partial}{\partial y^{i^{\prime}}}+\frac{\partial z^{i^{\prime}}}{\partial y^{i}} \frac{\partial}{\partial z^{i^{\prime}}},  \tag{1.5}\\
\frac{\partial}{\partial z^{i}}= & \frac{\partial z^{i^{\prime}}}{\partial z^{i}} \frac{\partial}{\partial z^{i^{\prime}}} .
\end{array}
$$

From (1.1) (c) it follows

$$
\begin{equation*}
\frac{\partial}{\partial z^{i}}=\frac{\partial x^{i^{\prime}}}{\partial x^{i}} \frac{\partial}{\partial z^{i^{\prime}}} \tag{1.6}
\end{equation*}
$$

Let us introduce the notation

$$
\begin{equation*}
\frac{\delta}{\delta y^{i}}=\frac{\partial}{\partial y^{i}}-\mathcal{H}_{i}^{j}(x, y) \frac{\partial}{\partial z^{i}} \tag{1.7}
\end{equation*}
$$

Proposition 1.1. $\frac{\delta}{\delta y^{i}}$ defined by (1.8) is transformed as tensor, i.e.

$$
\begin{equation*}
\frac{\delta}{\delta y^{i}}=\frac{\partial x^{i^{\prime}}}{\partial x^{i}} \frac{\delta}{\delta y^{i^{\prime}}} \tag{1.8}
\end{equation*}
$$

if $\mathcal{H}_{j}^{i}(x, y)$ is transformed in the following way:

$$
\begin{equation*}
\mathcal{H}_{i^{\prime}}^{j^{\prime}}=\mathcal{H}_{i}^{j} \frac{\partial x^{i}}{\partial x^{i^{\prime}}} \frac{\partial x^{j^{\prime}}}{\partial x^{j}}-\frac{\partial^{2} x^{j^{\prime}}}{\partial x^{i} \partial x^{j}} y^{j} \frac{\partial x^{i}}{\partial x^{i^{\prime}}} . \tag{1.9}
\end{equation*}
$$

The proof is obtained by direct calculation.
From (1.9) and (1.10) we obtain the first connection coefficient of Berwald type:

$$
\begin{equation*}
\mathcal{H}_{i^{\prime}}^{j^{\prime}} k^{\prime}=\frac{\partial \mathcal{H}_{i^{\prime}}^{j^{\prime}}}{\partial y^{k^{\prime}}}=\frac{\partial \mathcal{H}_{i}^{j}}{\partial y^{k}} \frac{\partial x^{i}}{\partial x^{i^{\prime}}} \frac{\partial x^{j^{\prime}}}{\partial x^{j}} \frac{\partial x^{k}}{\partial x^{k^{\prime}}}-\frac{\partial^{2} x^{j^{\prime}}}{\partial x^{i} \partial x^{k}} \frac{\partial x^{i}}{\partial x^{i^{\prime}}} \frac{\partial x^{k}}{\partial x^{k^{\prime}}} \tag{1.10}
\end{equation*}
$$

Let us introduce the notation:

$$
\begin{equation*}
\frac{\delta}{\delta x^{i}}=\frac{\partial}{\partial x^{i}}-\mathcal{M}_{i}^{j}(x, y) \frac{\partial}{\partial y^{j}}-\mathcal{N}_{i}^{j}(x, y, z) \frac{\partial}{\partial z^{j}} \tag{1.11}
\end{equation*}
$$

Proposition 1.2. $\frac{\delta}{\delta x^{i}}$ defined by (1.12) is transformed in the form:

$$
\begin{equation*}
\frac{\delta}{\delta x^{i}}=\frac{\partial x^{i^{\prime}}}{\partial x^{i}} \frac{\delta}{\delta x^{i}}, \tag{1.12}
\end{equation*}
$$

if $\mathcal{M}_{i}^{j}(x, y)$ and $\mathcal{N}_{i}^{j}(x, y, z)$ are transformed in the following way:

$$
\begin{align*}
\mathcal{M}_{i^{\prime}}^{j^{\prime}}= & \mathcal{M}_{i}^{j} \frac{\partial x^{i}}{\partial x^{i^{\prime}}} \frac{\partial x^{j^{\prime}}}{\partial x^{j}}-\frac{\partial^{2} x^{j^{\prime}}}{\partial x^{i} \partial x^{j}} y^{j} \frac{\partial x^{i}}{\partial x^{i^{\prime}}},  \tag{1.13}\\
\mathcal{N}_{i^{\prime}}^{j^{\prime}=}= & \mathcal{N}_{i}^{j} \frac{\partial x^{i}}{\partial x^{i^{\prime}}} \frac{\partial x^{j^{\prime}}}{\partial x^{j}}-\frac{\partial^{2} x^{j^{\prime}}}{\partial x^{i} \partial x^{j}} z^{j} \frac{\partial x^{i}}{\partial x^{i^{\prime}}} \\
& -\frac{1}{2} \frac{\partial^{3} x^{j^{\prime}}}{\partial x^{i} \partial x^{h} \partial^{k}} y^{h} y^{k} \frac{\partial x^{i}}{\partial x^{i^{\prime}}}+\frac{\partial^{2} x^{j^{\prime}}}{\partial x^{k} \partial x^{j}} y^{k} \mathcal{M}_{h}^{j} \frac{\partial x^{h}}{\partial x^{i^{\prime}}} . \tag{1.14}
\end{align*}
$$

Remark. $\mathcal{M}_{i}^{j}$ and $\mathcal{N}_{i}^{j}$ used here in [10], [11], [12] are denoted by (1) $\mathcal{N}_{i}^{j}$ and (2) $\mathcal{N}_{i}^{j}$ respectively.

Proof. If we add the following equations (which follow from (1.6), (1.1) and (1.3)):

$$
\begin{aligned}
\frac{\partial}{\partial x^{i}}= & \left(\frac{\partial x^{i^{\prime}}}{\partial x^{i}}\right) \frac{\partial}{\partial x^{i^{\prime}}}+\left(\frac{\partial^{2} x^{j^{\prime}}}{\partial x^{k} \partial x^{j}} y^{j} \frac{\partial x^{k}}{\partial x^{i^{\prime}}} \frac{\partial x^{i^{\prime}}}{\partial x^{i}}\right) \frac{\partial}{\partial y^{j^{\prime}}} \\
& +\left(\frac{1}{2} \frac{\partial^{3} x^{j^{\prime}}}{\partial x^{h} \partial x^{k} \partial x^{j}} y^{k} y^{j} \frac{\partial x^{h}}{\partial x^{i^{\prime}}} \frac{\partial x^{i^{\prime}}}{\partial x^{i}}+\frac{\partial^{2} x^{j^{\prime}}}{\partial x^{k} \partial x^{j}} z^{j} \frac{\partial x^{k}}{\partial x^{i^{\prime}}} \frac{\partial x^{i^{\prime}}}{\partial x^{i}}\right) \frac{\partial}{\partial z^{j^{\prime}}}, \\
-\mathcal{M}_{i}^{j} \frac{\partial}{\partial y^{j}}= & -\mathcal{M}_{h}^{j} \frac{\partial x^{h}}{\partial x^{i^{\prime}}} \frac{\partial x^{i^{\prime}}}{\partial x^{i}}\left(\frac{\partial x^{j^{\prime}}}{\partial x^{j}} \frac{\partial}{\partial y^{j^{\prime}}}+\frac{\partial^{2} x^{j^{\prime}}}{\partial x^{k} \partial x^{j}} y^{k} \frac{\partial}{\partial z^{j^{\prime}}}\right) \\
-\mathcal{N}_{i}^{j} \frac{\partial}{\partial z^{j}}= & -\mathcal{N}_{h}^{j} \frac{\partial x^{j^{\prime}}}{\partial x^{j}} \frac{\partial x^{h}}{\partial x^{i^{\prime}}} \frac{\partial x^{i^{\prime}}}{\partial x^{i}} \frac{\partial}{\partial z^{j^{\prime}}},
\end{aligned}
$$

we obtain $\frac{\delta}{\delta x^{i}}=\frac{\partial x^{i^{\prime}}}{\partial x^{i}}\left(\frac{\partial}{\partial x^{i^{\prime}}}-\mathcal{M}_{i^{\prime}}^{j^{\prime}} \frac{\partial}{\partial y^{j^{\prime}}}-\mathcal{N}_{i^{\prime}}^{j^{\prime}} \frac{\partial}{\partial z^{j^{\prime}}}\right)=\frac{\partial x^{i^{\prime}}}{\partial x^{i}} \frac{\delta}{\delta x^{i^{\prime}}}$, where $\mathcal{M}_{i^{\prime}}^{j^{\prime}}$ is given by (1.14) and $\mathcal{N}_{i^{\prime}}^{j^{\prime}}$ by (1.15).

The basis

$$
\begin{equation*}
B=\left\{\frac{\delta}{\delta x^{i}}, \frac{\delta}{\delta y^{i}}, \frac{\partial}{\partial z^{i}}\right\} \tag{1.15}
\end{equation*}
$$

is called adapted basis for $T\left(T^{2} M\right)$. Its elements are transformed by (1.7), (1.9) and (1.13).

From (1.14) and (1.15) we obtain two other connection coefficients of Berwald type, namely

$$
\begin{align*}
\mathcal{M}_{i^{\prime} k^{\prime}}^{j^{\prime}} & =\frac{\partial \mathcal{M}_{i^{\prime}}^{j^{\prime}}}{\partial x^{k^{\prime}}}=\mathcal{M}_{i}^{j}{ }_{k} \frac{\partial x^{i}}{\partial x^{i^{\prime}}} \frac{\partial x^{k}}{\partial x^{k^{\prime}}} \frac{\partial x^{j^{\prime}}}{\partial x^{j}}-\frac{\partial^{2} x^{j^{\prime}}}{\partial x^{i} \partial x^{k}} \frac{\partial x^{i}}{\partial x^{i^{\prime}}} \frac{\partial x^{k}}{\partial x^{k^{\prime}}}  \tag{1.16}\\
\mathcal{N}_{i^{\prime} k^{\prime}}^{j^{\prime}} & =\frac{\partial \mathcal{N}_{i^{\prime}}^{j^{\prime}}}{\partial z^{k^{\prime}}}=\mathcal{N}_{i}^{j}{ }_{k} \frac{\partial x^{i}}{\partial x^{i^{\prime}}} \frac{\partial x^{k}}{\partial x^{k^{\prime}}} \frac{\partial x^{j^{\prime}}}{\partial x^{j}}-\frac{\partial^{2} x^{j^{\prime}}}{\partial x^{i} \partial x^{k}} \frac{\partial x^{i}}{\partial x^{i^{\prime}}} \frac{\partial x^{k}}{\partial x^{k^{\prime}}} \tag{1.17}
\end{align*}
$$

2. The adapted basis in $T^{*}\left(T^{2} M\right)$. The natural basis of $T^{*}\left(T^{2} M\right)$ is

$$
\begin{equation*}
\bar{B}^{*}=\left\{d x^{i}, d y^{i}, d z^{i}\right\} \tag{2.1}
\end{equation*}
$$

where the following relations are valid with respect to (1.1):

$$
\begin{equation*}
\text { (a) } d x^{i^{\prime}}=\frac{\partial x^{i^{\prime}}}{\partial x^{i}} d x^{i} \quad \text { (b) } d y^{i^{\prime}}=\frac{\partial^{2} x^{i^{\prime}}}{\partial x^{j} \partial x^{i}} y^{i} d x^{j}+\frac{\partial x^{i^{\prime}}}{\partial x^{i}} d y^{i} \tag{2.2}
\end{equation*}
$$

(c) $d z^{i^{\prime}}=\left(\frac{1}{2} \frac{\partial^{3} x^{i^{\prime}}}{\partial x^{k} \partial x^{h} \partial x^{j}} y^{k} y^{h}+\frac{\partial^{2} x^{i^{\prime}}}{\partial x^{i} \partial x^{j}} z^{i}\right) d x^{j}+\left(\frac{\partial^{2} x^{i^{\prime}}}{\partial x^{k} \partial x^{j}} y^{k}\right) d y^{j}+\left(\frac{\partial x^{i^{\prime}}}{\partial x^{j}}\right) d z^{j}$.

From (2.2) (b) and (2.2) (c) it is obvious, that the transformations of $d y^{i^{i}}$ and $d z^{i^{\prime}}$ are not tensorial. If we put

$$
\begin{equation*}
\delta y^{i}=d y^{i}+\mathcal{H}_{j}^{i} d x^{j} \tag{2.3}
\end{equation*}
$$

then from (1.10) and (2.1) we get

$$
\begin{equation*}
\delta y^{i^{\prime}}=\frac{\partial x^{i^{\prime}}}{\partial x^{i}} \delta y^{i} \tag{2.4}
\end{equation*}
$$

Let us introduce the notation

$$
\begin{equation*}
\delta z^{i}=d z^{i}+\mathcal{M}_{j}^{i}(x, y) d y^{j}+\mathcal{N}_{j}^{i}(x, y, z) d x^{j}+\mathcal{G}_{j}^{i}(x, y) d x^{j} \tag{2.5}
\end{equation*}
$$

where the transformation laws of $\mathcal{M}_{j}^{i}$ and $\mathcal{N}_{j}^{i}$ are prescribed by (1.14) and (1.15).
Proposition 2.1. If $\mathcal{G}_{j}^{i}(x, y)$ has the following law of transformation (with respect to (1.1)):

$$
\begin{equation*}
\mathcal{G}_{j}^{i}=\mathcal{G}_{j^{\prime}}^{i^{\prime}} \frac{\partial x^{i}}{\partial x^{i^{\prime}}} \frac{\partial x^{j^{\prime}}}{\partial x^{j}}+\mathcal{M}_{j^{\prime}}^{i^{\prime}} \frac{\partial^{2} x^{j^{\prime}}}{\partial x^{k} \partial x^{j}} y^{k} \frac{\partial x^{i}}{\partial x^{i^{\prime}}}+\mathcal{M}_{j}^{k} \frac{\partial^{2} x^{j^{\prime}}}{\partial x^{k} \partial x^{h}} y^{h} \frac{\partial x^{i}}{\partial x^{j^{\prime}}} \tag{2.6}
\end{equation*}
$$

then $\delta z^{i}$ defined by (2.5) satisfies the equation

$$
\begin{equation*}
\delta z^{j^{\prime}}=\frac{\partial x^{j^{\prime}}}{\partial x^{j}} \delta z^{j} \tag{2.7}
\end{equation*}
$$

Proof. The substitution of (2.2), (1.14) and (1.15) into $\delta z^{j^{\prime}}=d z^{j^{\prime}}+$ $\mathcal{M}_{i^{\prime}}^{j^{\prime}} d y^{i^{\prime}}+\mathcal{N}_{i^{\prime}}^{j^{\prime}} d x^{i^{\prime}}+\mathcal{G}_{i^{\prime}}^{j^{\prime}} d x^{i^{i^{\prime}}}$ gives (2.7) if (2.6) is true.

The adapted basis of $T^{*}\left(T^{2} F\right)$ is

$$
\begin{equation*}
B^{*}=\left\{d x^{i}, \delta y^{i}, \delta z^{i}\right\} \tag{2.8}
\end{equation*}
$$

where the elements of $B^{*}$ defined by (2.3) and (2.5) satisfy the transformation law prescribed by (2.2) (a), (2.4) and (2.7).

Theorem 2.1. The bases $B$ (1.16) and $B^{*}(2.8)$ of $T\left(T^{2} M\right)$ and $T^{*}\left(T^{2} M\right)$ respectively are dual to each other if

$$
\begin{align*}
\mathcal{H}_{j}^{i}(x, y) & =\mathcal{M}_{j}^{i}(x, y)  \tag{2.9}\\
\mathcal{G}_{j}^{i}(x, y) & =\mathcal{M}_{k}^{i}(x, y) \mathcal{M}_{j}^{k}(x, y) \tag{2.10}
\end{align*}
$$

and if $\bar{B}$ (1.6) and $\bar{B}^{*}$ (2.1) are dual to each other.
Proof. By direct calculation using (1.8), (1.12), (2.3) and (2.5) we obtain:

$$
\begin{gather*}
\left\langle d x^{j}, \frac{\delta}{\delta x^{i}}\right\rangle=\delta_{i}^{j}, \quad\left\langle d x^{j}, \frac{\delta}{\delta y^{i}}\right\rangle=0, \quad\left\langle d x^{j} \frac{\partial}{\partial z^{i}}\right\rangle=0 \\
\left\langle\delta y^{j}, \frac{\delta}{\delta x^{i}}\right\rangle=\mathcal{H}_{i}^{j}-\mathcal{M}_{i}^{j}, \quad\left\langle\delta y^{j}, \frac{\delta}{\delta y^{i}}\right\rangle=\delta_{i}^{j}, \quad\left\langle\delta y^{j}, \frac{\partial}{\partial z^{i}}\right\rangle=0  \tag{2.11}\\
\left\langle\delta z^{j}, \frac{\delta}{\delta x^{i}}\right\rangle=\mathcal{G}_{i}^{j}-\mathcal{M}_{k}^{j} \mathcal{M}_{i}^{k}, \quad\left\langle\delta z^{j}, \frac{\delta}{\delta y^{i}}\right\rangle=-\mathcal{H}_{i}^{j}+\mathcal{M}_{i}^{j}, \quad\left\langle\delta z^{j}, \frac{\partial}{\partial z^{i}}\right\rangle=\delta_{i}^{j} .
\end{gather*}
$$

The duality follows from (2.11) and (2.9).
Proposition 2.2. $\mathcal{G}_{i}^{j}(x, y)$ defined by (2.10) satisfies (2.6).
Proof. Using (1.4) and (1.14), it can be proved that

$$
\begin{equation*}
\mathcal{M}_{r}^{s}=\mathcal{M}_{i^{\prime}}^{j^{\prime}} \frac{\partial x^{i^{\prime}}}{\partial x^{r}} \frac{\partial x^{s}}{\partial x^{j^{\prime}}}-\frac{\partial^{2} x^{s}}{\partial x^{j^{\prime}} \partial x^{t^{t^{\prime}}}} y^{t^{\prime}} \frac{\partial x^{j^{\prime}}}{\partial x^{r}} \tag{2.12}
\end{equation*}
$$

From (2.12) and (2.10) after some calculation we obtain (2.6).
Remark. It is important that the bases $B$ and $B^{*}$ be dual to each other, because if they are not, then the contraction of tensors doesn't result tensors. If $\bar{B}$ and $\bar{B}^{*}\left((1.6)\right.$ and (2.1)) are dual to each other, it doesn't follow that $\bar{B}^{\prime}$ and $\bar{B}^{*^{\prime}}$ are dual.

Now we have:
Theorem 2.2. The bases $B(\mathcal{M}, \mathcal{N})=\left\{\frac{\delta}{\delta x^{i}}, \frac{\delta}{\delta y^{i}}, \frac{\partial}{\partial z^{i}}\right\}$ and $B^{*}(\mathcal{M}, \mathcal{N})=$ $\left\{d x^{i}, \delta y^{i}, \delta z^{i}\right\}$, where their elements are given by

$$
\begin{gather*}
\frac{\delta}{\delta x^{i}}=\frac{\partial}{\partial x^{i}}-\mathcal{M}_{i}^{j} \frac{\partial}{\partial y^{j}}-\mathcal{N}_{i}^{j} \frac{\partial}{\partial z^{j}}, \quad \frac{\delta}{\delta y^{i}}=\frac{\partial}{\partial y^{i}}-\mathcal{M}_{i}^{j} \frac{\partial}{\partial z^{j}}  \tag{2.13}\\
\delta y^{i}=d y^{i}+\mathcal{M}_{j}^{i} d x^{j}, \quad \delta z^{i}=d z^{i}+\mathcal{M}_{j}^{i} \delta y^{j}+\mathcal{N}_{j}^{i} d x^{j}
\end{gather*}
$$

are adapted basis for $T\left(T^{2} M\right)$ and $T^{*}\left(T^{2} M\right)$ respectively, dual to each other, and they satisfy the law of transformation:

$$
\begin{align*}
\frac{\delta}{\delta x^{i}} & =\frac{\partial x^{i^{\prime}}}{\partial x^{i}} \frac{\delta}{\delta x^{i^{\prime}}}, \frac{\delta}{\delta y^{i}}=\frac{\partial y^{i^{\prime}}}{\partial y^{i}} \frac{\delta}{\delta y^{i^{\prime}}}, \frac{\partial}{\partial z^{i}}=\frac{\partial z^{i^{\prime}}}{\partial z^{i}} \frac{\partial}{\partial z^{i^{\prime}}}  \tag{2.14}\\
d x^{i} & =\frac{\partial x^{i}}{\partial x^{i^{\prime}}} d x^{i}, \quad \delta y^{i}=\frac{\partial y^{i}}{\partial y^{i^{\prime}}} \delta y^{i^{\prime}}, \delta z^{i}=\frac{\partial z^{i}}{\partial i^{\prime}} \delta z^{i^{\prime}}
\end{align*}
$$

It must be noted that there exist as many adapted bases as many functions $\mathcal{M}_{i}^{j}(x, y)$ and $\mathcal{N}_{i}^{j}(x, y, z)$ can be found, satisfying (1.14) and (1.15) respectively.

If we denote by $T_{H}, T_{V_{1}}, T_{V_{2}}$ the subspaces of $T\left(T^{2} M\right)$ spanned by $\left\{\frac{\delta}{\delta x^{2}}\right\}$, $\left\{\frac{\delta}{\delta y^{i}}\right\},\left\{\frac{\partial}{\partial z^{i}}\right\}$, and by $T_{H}^{*}, T_{V_{1}}^{*}, T_{V_{2}}^{*}$ the subspaces of $T^{*}\left(T^{2} M\right)$ spanned by $\left\{d x^{i}\right\}$, $\left\{\delta y^{i}\right\},\left\{\delta z^{i}\right\}$ respectively, then

$$
T\left(T^{2} M\right)=T_{H} \oplus T_{V_{1}} \oplus T_{V_{2}}, \quad T^{*}\left(T^{2} M\right)=T_{H}^{*} \oplus T_{V_{1}}^{*} \oplus T_{V_{2}}^{*}
$$

For the further examinations it is useful to introduce different kinds of indices. Indices $i, j, h, k, l=\overline{1, n}$ will be used in $T_{H}$ and $T_{H}^{*}, a, b, c, d, e, f=\overline{n+1,2 n}$ in $T_{V_{1}}$ and $T_{V_{1}}^{*}, p, q, r, s, t=\overline{2 n+1,3 n}$ in $T_{V_{2}}$ and $T_{V_{2}}^{*}$. The Greek letters as indices will take values from 1 to $3 n$. Using this notation the adapted bases have the form:

$$
\begin{equation*}
B=\left\{\frac{\delta}{\delta x^{i}}, \frac{\delta}{\delta y^{a}}, \frac{\partial}{\partial z^{p}}\right\}=\left\{\delta_{\alpha}\right\}, \quad B^{*}=\left\{d x^{j}, \delta y^{b}, \delta z^{q}\right\}=\left\{\delta^{\beta}\right\} \tag{2.15}
\end{equation*}
$$

where

$$
\begin{gather*}
\frac{\delta}{\delta x^{i}}=\frac{\partial}{\partial x^{i}}-\mathcal{M}_{i}^{b} \frac{\partial}{\partial y^{b}}-\mathcal{N}_{i}^{q} \frac{\partial}{\partial z^{q}}, \quad \frac{\delta}{\delta y^{a}}=\frac{\partial}{\partial y^{a}}-\mathcal{M}_{a}^{q} \frac{\partial}{\partial z^{q}}  \tag{2.16}\\
\delta y^{a}=d y^{a}+\mathcal{M}_{i}^{a} d x^{i}, \quad \delta z^{q}=d z^{q}+\mathcal{M}_{a}^{q} \delta y^{a}+\mathcal{N}_{j}^{q} d x^{j}
\end{gather*}
$$

If $i=a(\bmod n)$ and $j=b=q(\bmod n)$, then $\mathcal{M}_{i}^{j}=\mathcal{M}_{i}^{b}=\mathcal{M}_{a}^{q}, \mathcal{N}_{i}^{q}=\mathcal{N}_{i}^{j}$, in (2.16).

Some tensor field $T$ expressed in the bases $B$ and $B^{*}((2.15))$ has the form:

$$
T=T_{\ldots j \ldots b \ldots s \ldots}^{\ldots i \ldots a \ldots} \frac{\delta}{\delta x^{i}} \otimes d x^{j} \ldots \frac{\delta}{\delta y^{a}} \otimes \delta y^{b} \cdots \frac{\partial}{\partial z^{r}} \otimes \delta z^{s} \ldots
$$

The components of the tensor $T$, with respect to the coordinate transformations (1.1) are transformed in the following way:

For some vector field $X \in T\left(T^{2} M\right)$ and some 1-form $\omega \in T^{*}\left(T^{2} M\right)$ we have:

$$
\begin{gather*}
X=X^{i} \frac{\delta}{\delta x^{i}}+X^{a} \frac{\delta}{\delta y^{a}}+X^{p} \frac{\partial}{\partial z^{p}}=X^{\alpha} \delta_{\alpha}  \tag{2.17}\\
\omega=\omega_{j} d x^{j}+\omega_{b} \delta y^{b}+\omega_{q} \delta z^{q}=\omega_{\beta} \delta^{\beta}
\end{gather*}
$$

With respect to (1.1) the coordinates of $X$ and $\omega$ transform in the following way:

$$
\begin{array}{ll}
X^{i^{\prime}}=X^{i} \frac{\partial x^{i^{\prime}}}{\partial x^{i}}, \quad X^{a^{\prime}}=X^{a} \frac{\partial y^{a^{\prime}}}{\partial y^{a}}, \quad X^{p^{\prime}}=X^{p} \frac{\partial z^{p^{\prime}}}{\partial z^{p}} \\
\omega_{j^{\prime}}=\omega_{j} \frac{\partial x^{j}}{\partial x^{j^{\prime}}}, & \omega_{b^{\prime}}=\omega_{b} \frac{\partial y^{b}}{\partial y^{b^{\prime}}},
\end{array} \omega_{q^{\prime}}=\omega_{q} \frac{\partial z^{q}}{\partial z^{q^{\prime}}},
$$

because for $i=a=p(\bmod n)$ we have:

$$
\begin{equation*}
\frac{\partial x^{i^{\prime}}}{\partial x^{i}}=\frac{\partial y^{a^{\prime}}}{\partial y^{a}}=\frac{\partial z^{p^{\prime}}}{\partial z^{p}} \tag{2.18}
\end{equation*}
$$

3. Generalized covariant derivatives. The generalized connection in Lagrange and Hamilton spaces was studied among others in [2]-[6]. In $T\left(T^{2} M\right)$ it is introduced in the following way. Let $\nabla: T\left(T^{2} M\right) \times T\left(T^{2} M\right) \rightarrow T\left(T^{2} M\right)(\times$ is the Descartes product) be a linear connection, such that $\nabla:(X, Y) \rightarrow \nabla_{X} Y \in$ $T\left(T^{2} M\right), \forall X, Y \in T\left(T^{2} M\right)$. The operator $\nabla$ is called generalized connection. It is called $d$-connection if $\nabla_{X} Y$ is in $T_{H}, T_{V_{1}}$ or $T_{V_{2}}$ if $Y$ is in $T_{H}, T_{V_{1}}$ or $T_{V_{2}}$, respectively, $\forall X \in T\left(T^{2} M\right)$. It has been studied by many authors, mostly romanian geometers.

We shall not make that restriction on $\nabla$ here. In the following we shall use the abbreviations: $\delta_{k}=\frac{\delta}{\delta x^{k}}, \delta_{a}=\frac{\delta}{\delta y^{a}}, \partial_{k}=\frac{\partial}{\partial x^{k}}, \partial_{a}=\frac{\partial}{\partial y^{a}}, \partial_{p}=\frac{\partial}{\partial z^{p}}$.

Definition 3.1. The generalized connection $\nabla$ is defined by

$$
\begin{align*}
& \nabla_{\delta_{i}} \delta_{j}=\underline{F}_{j}^{k}{ }_{i} \delta_{k}+F_{j}{ }_{j} \delta_{i}+F_{j}^{r}{ }_{i} \partial_{r}, \quad \nabla_{\delta_{i}} \delta_{b}=F_{b}{ }^{k} \delta_{k}+\underline{F}_{b}{ }^{c} \delta_{c}+F_{b}{ }^{r} \partial_{r} \\
& \nabla_{\delta_{i}} \partial_{q}=F_{q}{ }_{i}^{k} \delta_{k}+F_{q i}^{c} \delta_{c}+\underline{F}_{q}^{r}{ }_{i} \partial_{r}, \quad \nabla_{\delta_{a}} \delta_{j}=\underline{C}_{j}{ }^{k} \delta_{k}+C_{j a}^{c} \delta_{c}+C_{j a}^{r} \partial_{r}, \\
& \nabla_{\delta_{a}} \delta_{b}=C_{b}{ }_{b}^{k} \delta_{k}+\underline{C}_{b}{ }_{a}^{c} \delta_{c}+C_{b}^{r}{ }_{a}^{r} \partial_{r}, \quad \nabla_{\delta_{a}} \partial_{q}=C_{q}^{k}{ }_{a}^{k} \delta_{k}+C_{q}^{c}{ }_{a}^{c} \delta_{c}+\underline{C}_{q}^{r}{ }_{a}^{r} \partial_{r},  \tag{3.1}\\
& \nabla_{\partial_{p}} \delta_{j}=\underline{L}_{j p}^{k} \delta_{k}+L_{j p}^{c} \delta_{c}+L_{j p}^{r} \partial_{r}, \quad \nabla_{\partial_{p}} \delta_{b}=L_{b p}^{k} \delta_{k}+\underline{L}_{b p}^{c} \delta_{c}+L_{b p}^{r} \partial_{r}, \\
& \nabla_{\partial_{p}} \partial_{q}=L_{q p}^{k} \delta_{k}+L_{q}^{c}{ }_{p} \delta_{c}+\underline{L}_{q}^{r}{ }_{p} \partial_{r} .
\end{align*}
$$

The $d$-connection is defined if in (3.1) all terms on the right-hand side vanish, except the underlined ones.

For the vector field $X$ defined by (2.17) we have

$$
\begin{aligned}
& \nabla_{\delta_{i}}\left(X^{j} \delta_{j}+X^{b} \delta_{b}+X^{q} \partial_{q}\right) \\
&=\left(\delta_{i} X^{j}\right) \delta_{j}+X^{j} \nabla_{\delta_{i}} \delta_{j}+\left(\delta_{i} X^{b}\right) \delta_{b}+X^{b} \nabla_{\delta_{i}} \delta_{b}+\left(\delta_{i} X^{q}\right) \partial_{q}+X^{q} \nabla_{\delta_{i}} \partial_{q} \\
&=\left(\delta_{i} X^{k}+F_{j}{ }^{k}{ }_{i} X^{j}+F_{b}{ }^{k}{ }_{i} X^{b}+F_{q}{ }^{k}{ }_{i} X^{q}\right) \delta_{k}+\left(\delta_{i} X^{c}+F_{j}{ }^{c} X^{j}+F_{b}{ }_{i} X^{b}+F_{q}{ }^{c} X^{q}\right) \delta_{c}+ \\
&\left(\delta_{i} X^{r}+F_{j}{ }^{r} X^{j}+F_{b}{ }^{r} X^{b}+F_{q}{ }^{r}{ }_{i} X^{q}\right) \partial_{r}
\end{aligned}
$$

From the above equation it follows

$$
\begin{equation*}
\nabla_{\delta_{i}} X=X_{\mid i}^{k} \delta_{k}+X_{\mid i}^{c} \delta_{c}+X_{\mid i}^{r} \partial_{r} \tag{3.2}
\end{equation*}
$$

where

$$
\begin{equation*}
X_{\mid i}^{x}=\delta_{i} X^{x}+F_{j}^{x}{ }_{i} X^{j}+F_{b i}^{x} X^{b}+F_{q}{ }_{i}^{x} X^{q}, \quad x \in\{k, c, r\} \tag{3.3}
\end{equation*}
$$

or shorter

$$
\begin{equation*}
X_{\mid i}^{x}=\delta_{i} X^{x}+F_{\alpha i}^{x} X^{\alpha}, \quad x \in\{k, c, r\} . \tag{3.4}
\end{equation*}
$$

The summation over $\alpha$ is the sum of summations over $j, b$ and $q$ as is written in (3.3). The sign $\mid i$ is the covariant derivative in direction of the basis vector $\delta_{i}$.

The covariant derivative of $X$ in the direction of $\delta_{a}$ has the form:

$$
\begin{equation*}
\nabla_{\delta_{a}} X=\left.X^{k}\right|_{a} \delta_{k}+\left.X^{c}\right|_{a} \delta_{c}+\left.X^{r}\right|_{a} \partial_{r} \tag{3.5}
\end{equation*}
$$

where

$$
\begin{equation*}
\left.X^{x}\right|_{a}=\delta_{a} X^{x}+C_{j a}^{x} X^{j}+C_{b a}^{x} X^{b}+C_{q}^{x} X^{q}, \quad x \in\{k, c, r\} \tag{3.6}
\end{equation*}
$$

or shorter

$$
\begin{equation*}
\left.X^{x}\right|_{a}=\delta_{a} X^{x}+C_{\alpha a}^{x} X^{\alpha} \tag{3.7}
\end{equation*}
$$

The covariant derivative of the vector field $X$ in the direction of $\partial_{p}$ is given by

$$
\begin{equation*}
\nabla_{\partial_{p}} X=X^{k}\left\|_{p} \delta_{k}+X^{c}\right\|_{p} \delta_{c}+X^{r} \|_{p} \partial_{r} \tag{3.8}
\end{equation*}
$$

where

$$
\begin{equation*}
X^{x} \|_{p}=\partial_{p} X^{x}+L_{j p}^{x} X^{j}+L_{b p}^{x} X^{b}+L_{q p}^{x} X^{q}, \quad x \in\{k, c, r\} \tag{3.9}
\end{equation*}
$$

or abbreviated

$$
\begin{equation*}
X^{x} \|_{p}=\partial_{p} X^{x}+L_{\alpha p}^{x} X^{\alpha} \tag{3.10}
\end{equation*}
$$

In (3.7) and (3.10) the summation over $\alpha$ is the sum of summations over $j, b$ and $q$ (as in (3.4)).

Theorem 3.1. If $X$ and $Y$ are vector fields in $T\left(T^{2} M\right), \nabla$ the generalized connection defined by (3.1), then the following equation is valid:

$$
\begin{align*}
& \nabla_{Y} X=\left(X_{\mid i}^{k} Y^{i}+\left.X^{k}\right|_{a} Y^{a}+X^{k} \|_{p} Y^{p}\right) \delta_{k}  \tag{3.11}\\
& \quad+\left(X_{\mid i}^{c} Y^{i}+\left.X^{c}\right|_{a} Y^{a}+X^{c} \|_{p} Y^{p}\right) \delta_{c}+\left(X_{\mid i}^{r} Y^{i}+\left.X^{r}\right|_{a} Y^{a}+X^{r} \|_{p} Y^{p}\right) \partial_{r}
\end{align*}
$$

Proof. The proof follows from (3.2)-(3.10) and the bilinearity of $\nabla$.
The equation (3.11) can be written in the abbreviated form as follows

$$
\begin{gather*}
\nabla_{Y} X=X_{\mid \beta}^{\alpha} Y^{\beta} \delta_{\alpha}  \tag{3.12}\\
X_{\mid \beta}^{\alpha}=\delta_{\beta} X^{\alpha}+\Gamma_{\gamma \beta}^{\alpha} X^{\gamma} \tag{3.13}
\end{gather*}
$$

If $\beta=i$, then $\Gamma=F$; if $\beta=a$, then $\Gamma=C$; if $\beta=p$ then $\Gamma=L$.
Theorem 3.2. All covariant derivatives $X_{\mid i}^{\alpha},\left.X^{\alpha}\right|_{a}, X^{\alpha} \|_{p}(\alpha=k$, or $\alpha=$ $c$ or $\alpha=r$ ) from (3.11) are transformed as tensors with respect to (1.1) if all connection coefficients from (3.1) are transformed as tensors, except the following, which have the form

$$
\begin{align*}
& F_{j}{ }^{k}=F_{j^{\prime}}^{k^{\prime}}{ }_{i^{\prime}} \frac{\partial x^{i^{\prime}}}{\partial x^{i}} \frac{\partial x^{k}}{\partial x^{k^{\prime}}} \frac{\partial x^{j^{\prime}}}{\partial x^{j}}+\frac{\partial^{2} x^{k^{\prime}}}{\partial x^{i} \partial x^{j}} \frac{\partial x^{k}}{\partial x^{k^{\prime}}} \\
& F_{b i}^{c}=F_{b^{\prime}}^{c^{\prime}}{ }^{\prime}  \tag{3.14}\\
& \frac{\partial x^{i^{\prime}}}{\partial x^{i}} \frac{\partial y^{b^{\prime}}}{\partial y^{b}} \frac{\partial y^{c}}{\partial y^{c^{\prime}}}+\frac{\partial^{2} y^{c^{\prime}}}{\partial x^{i} \partial y^{b}} \frac{\partial y^{c}}{\partial y^{c^{\prime}}} \\
& F_{q}^{r}{ }_{i}=F_{q^{\prime}}^{r^{\prime}}{ }_{i^{\prime}} \frac{\partial x^{i^{\prime}}}{\partial x^{i}} \frac{\partial z^{q^{\prime}}}{\partial z^{q}} \frac{\partial z^{r}}{\partial z^{r^{\prime}}}+\frac{\partial^{2} z^{r^{\prime}}}{\partial x^{i} \partial z^{q}} \frac{\partial z^{r}}{\partial z^{r^{\prime}}}
\end{align*}
$$

The proof is obtained by direct calculation.

Remark. From (2.17) it follows that $F_{j}{ }^{k}{ }_{i}=F_{b}{ }^{c}{ }_{i}=F_{q}{ }^{r}$ if $k=c=r(\bmod n)$, $j=b=q(\bmod n)$. The connection coefficients $\mathcal{H}_{j}^{i}, \mathcal{M}_{i}{ }_{k}$ and $\mathcal{N}_{i}{ }^{j}{ }_{k}$ defined by (1.11), (1.17) and (1.18) respectively, satisfy the transformation laws prescribed by (3.14).
4. The torsion tensor of the generalized connection. The torsion tensor $T(X, Y)$ is defined in the usual way by:

$$
T(X, Y)=\nabla_{X} Y-\nabla_{Y} X-[X, Y]
$$

Theorem 4.1. The torsion tensor for the generalized connection has the form $T(X, Y)=T^{k} \delta_{k}+T^{c} \delta_{c}+T^{r} \partial_{r}$, where

$$
\begin{aligned}
& T^{k}=T_{j}{ }_{i} Y^{j} X^{i}+T_{j}{ }_{b} Y^{j} X^{b}+T_{j}{ }_{q} Y^{j} X^{q} \\
& +T_{b}{ }_{i} Y^{b} X^{i}+T_{b}{ }_{a}^{k} Y^{b} X^{a}+T_{b}{ }_{q}^{k} Y^{b} X^{q} \\
& +T_{p}{ }_{i} Y^{p} X^{i}+T_{p b}^{k} Y^{p} X^{b}+T_{p}^{k} Y^{p} X^{q}, \\
& T_{j}{ }^{k}=F_{j}{ }^{k}{ }_{i}-F_{i}{ }_{j}^{k}, \quad T_{j}{ }^{k}=C_{j}{ }^{k}{ }_{b}-F_{b}{ }_{j}, \quad T_{j}{ }^{k}{ }_{q}=L_{j}^{k}{ }_{q}-F_{q}{ }^{k}, \\
& T_{b}{ }_{i}=F_{b}{ }^{k}-C_{i}{ }^{k}, \quad T_{b}{ }_{a}{ }_{a}=C_{b}{ }^{k}{ }_{a}-C_{a b}{ }^{k}, \quad T_{b}{ }^{k}=L_{b}{ }^{k}{ }_{q}-C_{q}{ }^{k} \text {, } \\
& T_{p i}^{k}=F_{p i}^{k}-L_{i}^{k}, \quad T_{p b}^{k}=C_{p b}^{k}-L_{b p}^{k}, \quad T_{p q}^{k}=L_{p q}^{k}-L_{q p}^{k}, \\
& T^{c}=T_{j}{ }_{i}^{c} Y^{j} X^{i}+T_{j}{ }^{c} Y^{j} X^{b}+T_{j}{ }^{c}{ }_{r} Y^{j} X^{r}+T_{b}{ }_{i} Y^{b} X^{i}+T_{a}{ }^{c}{ }_{b} Y^{a} X^{b}+ \\
& T_{b}{ }_{r} Y^{b} X^{r}+T_{q}{ }^{c}{ }_{i} Y^{q} X^{i}+T_{q}{ }^{c}{ }_{b} Y^{q} X^{b}+T_{q}{ }^{c}{ }_{r} Y^{q} X^{r}, \\
& T_{j}{ }^{c}{ }_{i}=F_{j}{ }^{c}{ }_{i}-F_{i}{ }^{c}{ }_{j}-K_{i}{ }^{c}{ }_{j}, \quad T_{a}^{c}{ }_{b}=C_{a b}^{c}-C_{b}{ }^{c}{ }_{a}, \quad T_{q}{ }^{c}{ }_{i}=F_{q}{ }^{c}{ }_{i}-L_{i}{ }^{c}{ }_{q}, \\
& T_{j b}{ }^{c}=C_{j b}{ }^{c}-F_{b}{ }^{c}{ }_{j}+K_{j}{ }^{c}{ }_{b}, \quad T_{b r}{ }^{c}=L_{b r}{ }^{c}-C_{r}{ }^{c}{ }_{b}, \quad T_{q b}{ }^{c}=C_{q b}{ }^{c}-L_{b}{ }^{c}{ }_{q}, \\
& T_{b}{ }_{i}=F_{b}{ }^{c}{ }_{i}-C_{i}{ }^{c} b-K_{i}{ }^{c}, \quad T_{j}{ }^{c}{ }_{r}=L_{j r}{ }^{c}{ }^{c}-F_{r}{ }^{c}, \quad T_{q r}{ }^{c}=L_{q r}{ }^{c}-L_{r}^{c}{ }_{q}, \\
& T^{r}=T_{j}^{r}{ }_{i} Y^{j} X^{i}+T_{j}^{r}{ }_{b} Y^{j} X^{b}+T_{j}{ }_{q}^{r} Y^{j} X^{q}+T_{b}^{r}{ }_{i} Y^{b} X^{i}+T_{b}{ }_{a}^{r} Y^{b} X^{a}+ \\
& T_{b}{ }_{q} Y^{b} X^{q}+T_{p}^{r}{ }_{i} Y^{p} X^{i}+T_{p}{ }_{b} Y^{p} X^{b}+T_{p}{ }_{q} Y^{p} X^{q} \\
& T_{j}^{r}{ }_{i}=F_{j}^{r}{ }_{i}-F_{i}{ }_{j}{ }^{r}-K_{i}{ }_{j}, \quad T_{j}{ }^{r}{ }_{b}=C_{j}{ }^{r} b-F_{b}{ }^{r}{ }_{j}+K_{j}{ }^{r}, \quad T_{b}{ }^{r}{ }_{q}=L_{b}^{r}{ }_{q}-C_{q}^{r}, \\
& T_{j}^{r}{ }_{q}=L_{j}{ }^{r}{ }_{q}-F_{q}{ }^{r}{ }_{j}+K_{j}{ }^{r} q, \quad T_{b}{ }^{r}{ }_{i}=F_{b}{ }^{r}{ }_{i}-C_{i}^{r}{ }_{b}-K_{i}{ }^{r}{ }_{b}, \quad T_{p}{ }^{r}{ }_{b}=C_{p}{ }^{r}{ }_{b}-L_{b}{ }^{r}{ }_{p}, \\
& T_{b}{ }^{r}{ }_{a}=C_{b}{ }^{r}{ }_{a}-C_{a b}^{r}-K_{a b}^{r}, \quad T_{p}{ }^{r}=F_{p}{ }^{r}-L_{i}^{r}{ }_{p}-K_{i}{ }^{r}, \quad T_{p}{ }^{r}{ }_{q}=L_{p}^{r}{ }_{q}-L_{q}^{r}{ }_{p} .
\end{aligned}
$$

Proof. The proof is obtained by direct calculation using (3.12), (3.13) and the relations:

$$
\begin{gathered}
{[X, Y]=\left[X\left(Y^{j}\right)-Y\left(X^{j}\right)\right] \delta_{j}+\left[X\left(Y^{b}\right)-Y\left(X^{b}\right)\right] \delta_{b}+} \\
{\left[X\left(Y^{q}\right)-Y\left(X^{q}\right)\right] \partial_{q}+A+B+C+D}
\end{gathered}
$$

where

$$
\begin{gathered}
A=X^{i} Y^{j}\left(\delta_{i} \delta_{j}-\delta_{j} \delta_{i}\right), \quad B=\left(X^{i} Y^{b}-Y^{i} X^{b}\right)\left(\delta_{i} \delta_{b}-\delta_{b} \delta_{i}\right) \\
C=\left(X^{i} Y^{q}-Y^{i} X^{q}\right)\left(\delta_{i} \partial_{q}-\partial_{q} \delta_{i}\right), \quad D=X^{a} Y^{b}\left(\delta_{a} \delta_{b}-\delta_{b} \delta_{a}\right)
\end{gathered}
$$

Using (2.16) we obtain:

$$
A=X^{i} Y^{j}\left(K_{i}{ }_{j}^{c} \partial_{c}+\bar{K}_{i j}^{r} \partial_{r}\right)=X^{i} Y^{j}\left(K_{i}{ }_{j} \delta_{c}+K_{i j}^{r} \partial_{r}\right),
$$

where

$$
\begin{gathered}
K_{i}{ }^{c}{ }_{j}=\delta_{j} \mathcal{M}_{i}^{c}-\delta_{i} \mathcal{M}_{j}^{c}, \quad K_{i}{ }^{r}{ }_{j}=\bar{K}_{i j}^{r}+\mathcal{M}_{c}^{r} K_{i}{ }^{c}, \bar{K}_{i j}^{r}=\delta_{j} \mathcal{N}_{i}^{r}-\delta_{i} \mathcal{N}_{j}^{r} . \\
B=\left(X^{i} Y^{b}-Y^{i} X^{b}\right)\left(K_{i}{ }^{c}{ }_{b} \delta_{c}+K_{i}^{r} \partial_{r}\right) \\
K_{i b}{ }^{c}=\delta_{b} \mathcal{M}_{i}^{c}, K_{i}^{r}{ }_{b}=\bar{K}_{i b}^{r}+\mathcal{M}_{c}^{r} K_{i}{ }^{c}, \bar{K}_{i}^{r}{ }_{b}=\delta_{b} \mathcal{N}_{i}^{r}-\delta_{i} \mathcal{M}_{b}^{r} . \\
C=\left(X^{i} Y^{q}-Y^{i} X^{q}\right) K_{i}{ }^{r} \partial_{r}, K_{i}{ }^{r}{ }_{q}=\frac{\partial \mathcal{N}_{i}^{r}}{\partial z^{q}} \\
D=X^{a} Y^{b} K_{a}^{r}{ }_{b} \partial_{r}, \quad K_{a}^{r}{ }_{b}=\delta_{b} \mathcal{M}_{a}^{r}-\delta_{a} \mathcal{M}_{b}^{r} .
\end{gathered}
$$

Remark. As $\mathcal{M}=\mathcal{M}(x, y)$, in all above formulae $\partial_{r} \mathcal{M}=0$.
5. Special cases. As mentioned in Definition 3.1 the special connection $\bar{\nabla}$ (the so called $d$-connection) is obtained, if in (3.1) only the underlined terms are left. More precisely:

Definition 5.1 The $d$-connection $\bar{\nabla}$ is defined by:

$$
\begin{array}{ccc}
\bar{\nabla}_{\delta_{i}} \delta_{j}=\bar{F}_{j}^{k} \delta_{k}, & \bar{\nabla}_{\delta_{i}} \delta_{b}=\bar{F}_{b i}^{c} \delta_{c}, & \bar{\nabla}_{\delta_{i}} \partial_{q}=\bar{F}_{q i}^{r} \partial_{r} \\
\bar{\nabla}_{\delta_{a}} \delta_{j}=\bar{C}_{j}^{k} \delta_{k}, & \bar{\nabla}_{\delta_{a}} \delta_{b}=\bar{C}_{b}^{c}{ }_{a} \delta_{c}, & \bar{\nabla}_{\delta_{a}} \partial_{q}=\bar{C}_{q a}^{r} \partial_{r}  \tag{5.1}\\
\bar{\nabla}_{\partial_{p}} \delta_{j}=\bar{L}_{j p}^{k} \delta_{k}, & \bar{\nabla}_{\partial_{p}} \delta_{b}=\bar{L}_{b}^{c}{ }_{p} \delta_{c}, & \bar{\nabla}_{\partial_{p}} \partial_{q}=\bar{L}_{q p}^{r} \partial_{r}
\end{array}
$$

From (5.1) the following property of $d$-connection is obvious:

$$
\bar{\nabla}_{X}: T_{H} \rightarrow T_{H}, \bar{\nabla}_{X}: T_{V_{1}} \rightarrow T_{V_{1}}, \bar{\nabla}_{X}: T_{V_{2}} \rightarrow T_{V_{2}}
$$

for any vector field $X$ from $T\left(T^{2} M\right)$.
Theorem 5.1. If $X$ and $Y$ are vector fields in $T\left(T^{2} M\right)$ expressed in the basis $B$ (2.15), then

$$
\begin{aligned}
\left.\bar{\nabla}_{X} Y=Y_{\overline{\lceil } i}^{k} X^{i}+Y^{k} \overline{\mid}_{a} X^{a}+Y^{k} \bar{\Pi}_{p} X^{p}\right) \delta_{k} & +\left(Y_{\overline{\lceil } i}^{c} X^{i}+Y^{c} \overline{\mid}_{a} X^{a}+Y^{c} \bar{\Pi}_{p} X^{p}\right) \delta_{c} \\
& +\left(Y_{\overline{\lceil } i}^{r} X^{i}+Y^{r} \bar{\Pi}_{a} X^{q}+Y^{r} \bar{\Pi}_{p} X^{p}\right) \partial_{r} \\
Y_{\bar{\Gamma} i}^{x}=\delta_{i} Y^{x}+\bar{F}_{y i}^{x} Y^{y}, \quad Y^{x} \overline{\lceil }_{a}=\partial_{a} Y^{x}+ & \bar{C}_{y a}^{x} Y^{y}, \quad Y^{x} \overline{\|}_{p}=\partial_{p} Y^{x}+\bar{L}_{y p}^{x} Y^{y}
\end{aligned}
$$

where either $x=k, y=j$, or $x=c, y=b$, or $x=r, y=q$.
Theorem 5.2. The connection coefficients $\bar{C}_{j}{ }_{a}^{k}, \bar{C}_{b a}^{c}, \bar{C}_{q}^{r}{ }_{a}, \bar{L}_{j p}^{k}, \bar{L}_{b p}^{c}$ and $\bar{L}_{q p}^{r}$ with respect to (1.1) are transformed as tensors. The transformation laws for
$\bar{F}_{j}{ }^{k}, \bar{F}_{b i}{ }_{i}$ and $\bar{F}_{q}{ }^{r}{ }_{i}$ are given by (3.14) if in these formulae the connection coefficients are overlined.

Theorem 5.3. The torsion tensor $T(X, Y)$ of d-connection $\bar{\nabla}$ has the form $T(X, Y)=\bar{T}^{k} \delta_{k}+\bar{T}^{c} \delta_{c}+\bar{T}^{r} \partial_{r}$, where

$$
\begin{aligned}
& \bar{T}^{k}=\bar{T}_{j}^{k}{ }_{i} Y^{j} X^{i}+\bar{T}_{j}^{k}{ }_{b} Y^{j} X^{b}+\bar{T}_{b}^{k} Y^{b} X^{j}+\bar{T}_{j}^{k} Y^{j} X^{q}+\bar{T}_{q}{ }_{j}^{k} Y^{q} X^{j}, \\
& \bar{T}^{c}=\bar{T}_{j}^{c}{ }_{i}^{j} Y^{i} X^{i}+\bar{T}_{i}^{c}{ }_{b} Y^{j} X^{b}+\bar{T}_{b}^{c}{ }_{j}^{b} Y^{b} X^{j}+\bar{T}_{a}^{c}{ }_{b} Y^{a} X^{b}+\bar{T}_{b}^{c}{ }_{q} Y^{b} X^{q}+\bar{T}_{q}^{c}{ }_{b} Y^{q} X^{b}, \\
& \bar{T}^{r}= \\
& \quad \bar{T}_{j}^{r}{ }_{i} Y^{j} X^{i}+\bar{T}_{b}^{r}{ }_{a}^{r} Y^{b} X^{a}+\bar{T}_{j}^{r}{ }_{q} Y^{j} X^{q}+\bar{T}_{q}^{r}{ }_{j} Y^{q} X^{j} \\
& \quad+\bar{T}_{b}^{r}{ }_{q} Y^{b} X^{q}+\bar{T}_{q}^{r}{ }_{b} Y^{q} X^{b}+\bar{T}_{p}^{r}{ }_{q} Y^{p} X^{q}
\end{aligned}
$$

and

$$
\begin{align*}
& \bar{T}_{j}{ }_{i}{ }_{i}=\bar{F}_{j}{ }_{i}{ }_{i}-\bar{F}_{i}{ }_{j}^{k}, \quad \quad \bar{T}_{j}{ }^{k}{ }_{b}=-\bar{T}_{b}{ }_{j}=\bar{C}_{j}{ }^{k}, \\
& \bar{T}_{j}{ }^{k}{ }_{q}=-\bar{T}_{q}{ }_{j}=\bar{L}_{j}{ }^{k}, \quad \bar{T}_{j}{ }_{i}{ }^{\prime}=-K_{i}{ }^{c}{ }_{j}, \\
& \bar{T}_{b}{ }^{c}{ }_{j}=-\bar{T}_{j}{ }^{c}{ }_{b}=\bar{F}_{b}{ }^{c}{ }_{j}-K_{j b}{ }^{c}, \quad \bar{T}_{a}{ }^{c}{ }_{b}=\bar{C}_{a b}{ }^{c}-\bar{C}_{b}{ }^{c}{ }_{a}, \\
& \bar{T}_{b}{ }^{c}{ }_{r}=-\bar{T}_{r}{ }^{c}{ }_{b}=\bar{L}_{b r}{ }^{c}, \quad \bar{T}_{q}{ }^{r}{ }_{j}=-\bar{T}_{j}^{r}{ }_{q}=\bar{F}_{q}{ }^{r}{ }_{j}-K_{j}^{r}{ }_{q},  \tag{5.2}\\
& \bar{T}_{q}{ }^{r}{ }_{b}=-\bar{T}_{b}{ }^{r}{ }_{q}=\bar{C}_{q b}^{r}, \quad \bar{T}_{p}{ }^{r}{ }_{q}=\bar{L}_{p}^{r}{ }_{q}-\bar{L}_{q p}^{r}, \\
& \bar{T}_{j}{ }^{r}{ }_{i}=-K_{i}{ }^{r}, \quad \quad \bar{T}_{b}{ }^{r}{ }_{a}=-K_{a b}^{r} .
\end{align*}
$$

It is easy to see, that all components of the torsion tensor $T$ which appear in (5.2) with respect to (1.1) transform as tensors.

## REFERENCES

1. M. Anastasiei, Geometry of higher order sprays, (preprint).
2. I.Čomić, Generalized Miron's d-connection in the recurrent $K$-Hamilton spaces, Publ. Inst. Math. Beograd 52(66) (1992), 136-152.
3. I.Čomić, Induced generalized connection in vector subbundles, Publ. Inst. Math. Beograd 53(67) (1993), 103-114.
4. I.Čomić, Recurrent Hamilton spaces with generalized Miron's d-connection, Ann. Univ. Craiova 18 (1990), 90-104.
5. I.Čomić, Generalized connections on vector bundles and subbundles Coll. Sci. Papers Fac. Sci. Kragujevac 16 (1994), 13-32.
6. I.Čomić, T. Kawaguchi, H. Kawaguchi, A theory of dual vector bundles and their subbundles for general dynamical systems of the information geometry, Tensor 52 (1993), 286-300.
7. A. Kawaguchi, On the vectors of higher order and the extended affine connections, Ann. Pura Appl. (IV), 55 (1961), 105-118.
8. M. Matsumoto, Foundations of Finsler Geometry and Special Finsler Spaces, Kaiseisha Press, Otsu, Japan, 1986.
9. R. Miron, M. Anastasiei, The geometry of Lagrange space, theory and applications, Kluwer, Amsterdam, 1993.
10. R. Miron, Gh. Atanasiu, Compendium sur les espaces Lagrange d'ordre supérieur, Seminarul de Mecanica 40, Universitatea din Timisoara, 1994.
11. R. Miron, Gh. Atanasiu, Differential geometry of the $k$-osculator bundle, Rev. Roumaine Math. Pures Appl. (to appear).
12. R. Miron, Gh. Atanasiu, Higher order Lagrange spaces, Rev. Roumaine Math. Pures Appl. (to appear).
13. Gh. Munteanu, Gh. Atanasiu, On Miron-connections in Lagrange spaces of second order, Tensor (N.S.) 50 (1991), 241-247.

Fakultet tehničkih nauka
(Received 0702 1996)
Trg Dositeja Obradovića 6 21000 Novi Sad
Yugoslavia

