

**ON THE ESTIMATES OF THE CONVERGENCE
 RATE OF THE FINITE DIFFERENCE SCHEMES
 FOR THE APPROXIMATION OF SOLUTIONS
 OF HYPERBOLIC PROBLEMS, II part**

Boško Jovanović

Abstract. Some new estimates of the convergence rate for hyperbolic initial-boundary value problems are obtained using the interpolation theory of the function spaces.

1. Introduction

For the finite difference approximation to the solution of vibrating string equation, in [1] fractional order convergence rate estimates are obtained, e.g.

$$\|\bar{z}\|_{C_\tau(W_{2,h}^1)} \leq C(h + \tau)^{\frac{2}{3}(s-1)} \|u_0\|_{W_2^s}, \quad 1 \leq s \leq 4$$

(see also [4]). Here we show how the similar estimates in the other discrete norms can be obtained.

In the sequel we use the notations introduced in [1]. When we quote a formula from [1] we add roman I to the corresponding number, e.g. (I.3).

As the model problem we consider the first initial-boundary value problem (IBVP) for the vibrating string equation, in the domain $Q = (0, 1) \times (0, T]$:

$$(1) \quad \begin{aligned} \partial^2 u / \partial t^2 &= \partial^2 u / \partial x^2, & (x, t) \in Q, \\ u(0, t) &= u(1, t) = 0, & t \in [0, t], \\ u(x, 0) &= u_0(x), & \partial u(x, 0) / \partial t = 0, & x \in (0, 1). \end{aligned}$$

In the sequel, we assume that $u_0(x) \in W_2^s(0, 1)$, $s \geq 0$, and can be oddly extended for $x < 0$ and $x > 1$, preserving the class.

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Let us note that the formula (I.3) holds for $k = 0$. Indeed, let us multiply the equation (1) by the function

$$w(x, t) = x \int_0^1 \int_0^\xi u(\eta, t) d\eta d\xi - \int_0^x \int_0^\xi u(\eta, t) d\eta d\xi,$$

satisfying the conditions

$$\frac{\partial^2 w}{\partial x^2} = -u(x, t), \quad w(0, t) = w(1, t) = 0.$$

Integrating the obtained equation, after some simple transformations, we obtain:

$$\max_{t \in [0, 1]} \left(\left\| \frac{\partial^2 w}{\partial t \partial x} \right\|_{L_2}^2 + \|u\|_{L_2}^2 \right) = \|u_0\|_{L_2}^2,$$

and

$$(2) \quad \|u\|_{C(L_2)} \leq \|u_0\|_{L_2}.$$

Let $\bar{\omega}_h$ be a uniform mesh on $[0, 1]$ with the step-size h , and $\bar{\omega}_\tau$ — a uniform mesh on $[-\tau/2, T]$ with the step-size τ . In addition to the norms defined in [1], let us set:

$$\begin{aligned} \|v\|_{W_{2,h}^{-1}} &= \left\{ h \sum_{x \in \omega_h^-} \left[h^2 \sum_{\xi \in \omega_h} \sum_{\substack{\eta \in \omega_h \\ \eta \leq \xi}} v(\eta) - h \sum_{\substack{\eta \in \omega_h \\ \eta \leq x}} v(\eta) \right]^2 \right\}^{1/2}, \\ \|v\|_{W_{2,h}^2} &= (\|v\|_h^2 + \|v_x\|_h^2 + \|v_{x\bar{x}}\|_h^2)^{1/2}, \quad \|v\|_{C_\tau(W_{2,h}^2)} = \max_{t \in \omega_\tau^-} \|v(\cdot, t)\|_{W_{2,h}^2}, \\ \|v\|_{C_\tau(L_{2,h})} &= \max_{t \in \omega_\tau^-} \|v(\cdot, t)\|_{L_{2,h}}, \quad \|v\|_{L_{1,\tau}(L_{2,h})} = \tau \sum_{t \in \omega_\tau^-} \|v(\cdot, t)\|_{L_{2,h}}. \end{aligned}$$

Finally, let C be the positive generic constant, independent of h and τ .

2. Estimates in the second order norm

Let us approximate the IVP (1) with the standard symmetric weighted finite-difference scheme (FDS) [2]

$$(3) \quad v_{t\bar{t}} = [\sigma \hat{v} + (1 - 2\sigma) v + \sigma \check{v}]_{x\bar{x}}, \quad x \in \omega_h, \quad t \in \omega_\tau,$$

$$(4) \quad v(0, t) = v(1, t) = 0, \quad t \in \bar{\omega}_\tau,$$

$$(5) \quad v^0 = v^1 = u_0(x), \quad x \in \omega_h.$$

Using the energy method we easily obtain the equality:

$$(N_2(v))^2 \equiv \|v_{tx}\|_h^2 + \tau^2 (\sigma - 0.25) \|v_{tx\bar{x}}\|_h^2 + \|\bar{v}_{x\bar{x}}\|_h^2 = \|v_{x\bar{x}}^0\|_h^2.$$

The expression $N_2(v)$ is the norm for $\sigma \geq 1/4$, while for $\sigma < 1/4$ it is the norm if

$$\tau \leq h \sqrt{\frac{1 - c_0}{1 - 4\sigma}}, \quad c_0 = \text{const} \in (0, 1).$$

Further,

$$\|v_{x\bar{x}}^0\|_h = \|u_{0,x\bar{x}}\|_h = \left\{ h \sum_{x \in \omega_h} \left[\frac{1}{h} \int_{x-h}^{x+h} \left(1 - \frac{|\xi - x|}{h} \right) u_0''(\xi) d\xi \right]^2 \right\}^{1/2} \leq \frac{2}{\sqrt{3}} \|u_0''\|_{L_2}$$

and

$$(6) \quad \max_{\tau \in \omega_\tau^-} N_2(v) \leq C \|u_0\|_{W_2^2}.$$

The error $z = u - v$ satisfies the conditions

$$(7) \quad z_{t\bar{t}} = [\sigma \hat{z} + (1 - 2\sigma) z + \sigma \check{z}]_{x\bar{x}} + \psi, \quad x \in \omega_h, \quad t \in \omega_\tau,$$

$$(8) \quad z(0, t) = z(1, t) = 0, \quad t \in \overline{\omega}_\tau,$$

$$(9) \quad z^0 = z^1 = u(x, \tau/2) - u_0(x), \quad x \in \omega_h,$$

where $\psi = u_{t\bar{t}} - [\sigma \hat{u} + (1 - 2\sigma) u + \sigma \check{u}]_{x\bar{x}}$.

The a priori estimate

$$(10) \quad \max_{t \in \omega_\tau^-} N_2(z) \leq \|z_{x\bar{x}}^0\|_h + \frac{1}{\sqrt{c}} \|\psi_x\|_{L_{1,\tau}(L_{2,h})}$$

holds, where $c = 1$ for $\sigma \geq 1/4$, and $c = c_0$ for $\sigma < 1/4$.

Using the representations of z^0 and ψ obtained in [1] we easily obtain:

$$z_{x\bar{x}}^0 = \frac{1}{h} \int_{x-h}^{x+h} \int_0^{\tau/2} \int_0^t \left(1 - \frac{|x - \xi|}{h} \right) \frac{\partial^4 u(\xi, \eta)}{\partial t^2 \partial x^2} d\eta dt d\xi$$

and

$$\begin{aligned} \psi_x(x, t) = & \\ & - \frac{1}{h^2 \tau} \int_{x-h}^{x+h} \int_x^\xi \int_\eta^{t+h} \int_{t-\tau}^{t+\tau} (\xi - \eta) \left(1 - \frac{|\xi - x|}{h} \right) \left(1 - \frac{|\chi - t|}{\tau} \right) \frac{\partial^5 u(\zeta, \chi)}{\partial x^3 \partial t^2} d\chi d\zeta d\eta d\xi \\ & + \frac{1}{h^2 \tau} \int_{x-h}^{x+h} \int_\xi^{\xi+h} \int_{t-\tau}^{t+\tau} \int_t^\zeta (\zeta - \chi) \left(1 - \frac{|\xi - x|}{h} \right) \left(1 - \frac{|\zeta - t|}{\tau} \right) \frac{\partial^5 u(\eta, \chi)}{\partial x^3 \partial t^2} d\chi d\zeta d\eta d\xi \\ & - \frac{\sigma \tau}{h^2} \int_{x-h}^{x+h} \int_\xi^{\xi+h} \int_{t-\tau}^{t+\tau} \left(1 - \frac{|\xi - x|}{h} \right) \left(1 - \frac{|\zeta - t|}{\tau} \right) \frac{\partial^5 u(\eta, \zeta)}{\partial x^3 \partial t^2} d\zeta d\eta d\xi. \end{aligned}$$

Consequently

$$(11) \quad \|z_{x\bar{x}}^0\|_h \leq C \tau^2 \max_{t \in [0, T]} \left\| \frac{\partial^4 u}{\partial t^2 \partial x^2} \right\|_{L_2}$$

and

$$(12) \quad \|\psi_x\|_{L_{1,\tau}(L_{2,h})} \leq C (h + \tau)^2 \max_{t \in [0, T]} \left\| \frac{\partial^5 u}{\partial x^3 \partial t^2} \right\|_{L_2(0,1)}.$$

From the inequalities (10)–(12) and (I.3) we immediately obtain the following convergence rate estimate:

$$(13) \quad \max_{t \in \omega_\tau^-} N_2(z) \leq C (h + \tau)^2 \|u_0\|_{W_2^5}.$$

On the other hand, from

$$\begin{aligned} \max_{t \in \omega_\tau^-} N_2(z) &\leq \max_{t \in \omega_\tau^-} N_2(u) + \max_{t \in \omega_\tau^-} N_2(v), \\ \max_{t \in \omega_\tau^-} N_2(u) &\leq C \left(\left\| \frac{\partial^2 u}{\partial t \partial x} \right\|_{C(L_2)} + \left\| \frac{\partial^2 u}{\partial x^2} \right\|_{C(L_2)} \right) \leq C \|u_0\|_{W_2^2}, \end{aligned}$$

and (6) follows

$$(14) \quad \max_{t \in \omega_\tau^-} N_2(z) \leq C \|u_0\|_{W_2^2}.$$

Let D denote the space of functions defined on the mesh $\bar{\omega}_h \times \bar{\omega}_\tau$, with the norm $\max_{t \in \omega_\tau^-} N_2(\cdot)$, and let R be the linear mapping $u_0 \mapsto z$. From the relations (13) and (14) it follows that R is a bounded operator from W_2^5 into D , as well as from W_2^2 into D . Using the interpolation theory of function spaces [3] we conclude that R is a bounded linear operator from $(W_2^2, W_2^5)_{\theta,2}$ into D ($0 < \theta < 1$), and

$$\max_{t \in \omega_\tau^-} N_2(z) \leq C (h + \tau)^{2\theta} \|u_0\|_{(W_2^2, W_2^5)_{\theta,2}}.$$

Further [3], we have

$$(W_2^2, W_2^5)_{\theta,2} = W_2^{2(1-\theta)+5\theta} = W_2^{3\theta+2}.$$

Setting $3\theta + 2 = s$ we finally obtain the required convergence rate estimate for the FDS (3)–(5):

$$(15) \quad \max_{t \in \omega_\tau^-} N_2(z) \leq C (h + \tau)^{\frac{2}{3}(s-2)} \|u_0\|_{W_2^s}, \quad 2 \leq s \leq 5.$$

From (15) we also obtain

$$(16) \quad \|z\|_{C_\tau(W_{2,h}^2)} \leq C (h + \tau)^{\frac{2}{3}(s-2)} \|u_0\|_{W_2^s}, \quad 2 \leq s \leq 5.$$

3. Estimate in L_2 -norm

To obtain the error estimate in the case of less smooth solutions let us approximate the initial conditions with the average values:

$$(17) \quad v^0 = v^1 = S_x u_0,$$

where S_x is the Steklov smoothing operator [1]. The FDS (3), (4), (17) satisfies the relation

$$(18) \quad (N_0(v))^2 \equiv \|v_t\|_{W_{2,h}^{-1}}^2 + \tau^2 (\sigma - 0.25) \|v_t\|_h^2 + \|\bar{v}\|_h^2 = \|v^0\|_h^2 \leq \|u_0\|_{L_2}^2.$$

The expression $N_0(v)$ is a norm if σ satisfies the previous conditions (in section 2).

For $s < 1/2$ the solution of the IBVP (1) may not be continuous, so we define the error by:

$$z = S_x u - v.$$

Such defined error satisfies the conditions (7) and (8), with

$$\psi = S_x \{u_{t\bar{t}} - [\sigma \hat{u} + (1 - 2\sigma) u + \sigma \check{u}]_{x\bar{x}}\},$$

and the initial conditions

$$(19) \quad z^0 = z^1 = S_x u(\cdot, \tau/2) - S_x u_0.$$

The a priori estimate

$$(20) \quad \max_{t \in \omega_\tau^-} N_0(z) \leq \|z^0\|_h + \frac{1}{\sqrt{c}} \tau \sum_{t \in \omega_\tau^-} \|\psi_x\|_{W_{2,h}^{-1}}.$$

holds, where, as before, $c = 1$ for $\sigma \geq 1/4$, and $c = c_0$ for $\sigma < 1/4$.

From the equality (19) we easily obtain

$$z^0 = \frac{1}{h} \int_{x-h/2}^{x+h/2} \int_0^{\tau/2} \int_0^\eta \frac{\partial^2 u(\xi, \zeta)}{\partial t^2} d\xi d\eta d\xi$$

wherfrom we get

$$(21) \quad \|z^0\|_h \leq C \tau^2 \left\| \frac{\partial^2 u}{\partial t^2} \right\|_{C(L_2)} \leq C \tau^2 \|u_0\|_{W_2^2}.$$

The expression $\|\psi_x\|_{W_{2,h}^{-1}}$ can be represented as

$$\|\psi_x\|_{W_{2,h}^{-1}} = \|M + \varphi\|_h \leq |M| + \|\varphi\|_h,$$

where

$$\begin{aligned} M &= M(t) = S_t^2 \left(h \sum_{x \in \omega_h^-} \frac{\partial u(x + h/2, t)}{\partial x} \right) \\ &= \frac{1}{\tau} \sum_{x \in \omega_h^-} \int_{t-\tau}^{t+\tau} \int_x^{x+h} \int_\xi^{x+h/2} \int_{x+h/2}^\eta \left(1 - \frac{|t-\chi|}{\tau} \right) \frac{\partial^3 u(\zeta, \chi)}{\partial x^3} d\zeta d\eta d\xi d\chi \end{aligned}$$

and

$$\begin{aligned}
\varphi = & S_x [\sigma \hat{u} + (1 - 2\sigma) u + \sigma \check{u}]_x - S_t^2 \left[\frac{\partial u(x + h/2, t)}{\partial x} \right] \\
= & \frac{\sigma \tau}{h^2} \int_{x-h/2}^{x+h/2} \int_{\xi}^{\xi+h} \int_{t-\tau}^{t+\tau} \left(1 - \frac{|t-\zeta|}{\tau} \right) \frac{\partial^3 u(\eta, \zeta)}{\partial t^2 \partial x} d\zeta d\eta d\xi \\
& + \frac{1}{h^2} \int_{x-h/2}^{x+h/2} \int_x^{\xi} \int_x^{\eta} \int_{\zeta}^{\zeta+h} \frac{\partial^3 u(\chi, t)}{\partial x^3} d\chi d\zeta d\eta d\xi \\
& + \frac{1}{h\tau} \int_x^{x+h} \int_{t-\tau}^{t+\tau} \int_{\eta}^t \int_t^{\zeta} \left(1 - \frac{|t-\eta|}{\tau} \right) \frac{\partial^3 u(\xi, \chi)}{\partial t^2 \partial x} d\chi d\zeta d\eta d\xi \\
& + \frac{1}{h\tau} \int_x^{x+h} \int_{x+h/2}^{\xi} \int_{x+h/2}^{\eta} \int_{t-\tau}^{t+\tau} \left(1 - \frac{|t-\chi|}{\tau} \right) \frac{\partial^3 u(\zeta, \chi)}{\partial x^3} d\chi d\zeta d\eta d\xi.
\end{aligned}$$

From these representations and (I.3) follows

$$\begin{aligned}
\max_{t \in \omega_{\tau}^-} |M| & \leq C h^2 \left\| \frac{\partial^3 u}{\partial x^3} \right\|_{C(L_2)} \leq C h^2 \|u_0\|_{W_2^3}, \\
\max_{t \in \omega_{\tau}^-} \|\varphi\|_h & \leq C (h + \tau)^2 \left(\left\| \frac{\partial^3 u}{\partial x^3} \right\|_{C(L_2)} + \left\| \frac{\partial^3 u}{\partial t^2 \partial x} \right\|_{C(L_2)} \right) \leq C (h + \tau)^2 \|u_0\|_{W_2^3},
\end{aligned}$$

and

$$(22) \quad \max_{t \in \omega_{\tau}^-} \|\psi\|_{W_{2,h}^{-1}} \leq C (h + \tau)^2 \|u_0\|_{W_2^3}.$$

Finally, from (20)–(22) we obtain the following estimate of the convergence rate for the FDS (3), (4), (17):

$$(23) \quad \max_{t \in \omega_{\tau}^-} \|\bar{z}\|_h \leq \max_{t \in \omega_{\tau}^-} N_0(z) \leq C (h + \tau)^2 \|u_0\|_{W_2^3}.$$

On the other hand, from evident inequality

$$\max_{t \in \omega_{\tau}^-} \|\bar{z}\|_h \leq \max_{t \in \omega_{\tau}^-} \|\bar{v}\|_h + \max_{t \in \omega_{\tau}^-} \|\overline{S_x u}\|_h,$$

and estimates (18) and (2) we obtain

$$\max_{t \in \omega_{\tau}^-} \|\bar{v}\|_h \leq \max_{t \in \omega_{\tau}^-} N_0(v) = \|v^0\|_h = \|S_x u_0\|_h \leq \|u_0\|_{L_2},$$

and

$$\max_{t \in \omega_{\tau}^-} \|\overline{S_x u}\|_h \leq \max_{t \in \omega_{\tau}^-} \|S_x u\|_h \leq \max_{t \in [0, T]} \|u\|_{L_2} \leq \|u_0\|_{L_2}.$$

Consequently

$$(24) \quad \max_{t \in \omega_{\tau}^-} \|\bar{z}\|_h \leq 2 \|u_0\|_{L_2}.$$

From (23) and (24) by interpolation we obtain the following convergence rate estimate for the FDS (3), (4), (17):

$$(25) \quad \|\bar{z}\|_{C_\tau(L_{2,h})} = \max_{t \in \omega_\tau^-} \|\bar{z}\|_h \leq C (h + \tau)^{\frac{2}{3}s} \|u_0\|_{W_2^s}, \quad 0 \leq s \leq 3.$$

4. Fourth-order scheme

For $\tau < h$, in the same way for the fourth-order difference scheme (I.21), (I.5), (I.22) we obtain the following convergence rate estimates:

$$\begin{aligned} \max_{t \in \omega_\tau^-} N_2(z) &\leq C h^{\frac{4}{5}(s-2)} \|u_0\|_{W_2^s}, \quad 2 \leq s \leq 7, \\ \|z\|_{C_\tau(W_{2,h}^2)} &\leq C h^{\frac{4}{5}(s-2)} \|u_0\|_{W_2^s}, \quad 2 \leq s \leq 7, \end{aligned}$$

and

$$\|\bar{z}\|_{C_\tau(L_{2,h})} \leq C h^{\frac{4}{5}s} \|u_0\|_{W_2^s}, \quad 0 \leq s \leq 5.$$

In the last case, initial condition (I.22) should be replaced by

$$v^0 = v^1 = S_x u_0 + \frac{\tau^2}{8} (S_x u_0)_{x\bar{x}}.$$

The obtained results could be extended, without difficulties, to the case of IVP with nonhomogeneous second initial condition.

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Matematički fakultet
11001 Beograd, p.p. 550
Jugoslavija

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