

THE REGULARIZED TRACE FOR NUCLEAR PERTURBATIONS OF DISCRETE OPERATORS

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Abstract. The regularized trace of nuclear perturbations for semibounded (positive) discrete operators is found.

1. Introduction

In [1] the regularized trace was found for a perturbation of a discrete operator by an operator of finite rank. For this result the behaviour of the function $N(\lambda)$ played an important role, where $N(\lambda)$ denotes the distribution of the eigenvalues for the nonperturbed operator.

In this paper this result is generalized for perturbations of a discrete operator by a nuclear operator which satisfies some additional conditions.

2. Statement of the problem and the main result

Let L_0 be a selfadjoint discrete operator acting on a separable Hilbert space H , with domain $\mathcal{D}(L_0)$. We denote by S a nuclear operator with a Hilbert representation:

$$S = \sum_{k=1}^{\infty} s_k(\cdot, f_k)g_k, \quad \left(\sum_{k=1}^{\infty} s_k < \infty \right).$$

Let us consider the operator $L = L_0 + SL_0$. In this case μ_ν and λ_ν are eigenvalues of the operators L and L_0 respectively, arranged in the increasing order of their real parts, each of which is repeated according to its multiplicity. The main result is the following theorem.

THEOREM. *Let L_0 be a selfadjoint, discrete, semibounded (from below) operator acting on H . Suppose that the distribution function of its eigenvalues has the form $N(\lambda) = \sum_{\lambda_n \leq \lambda} 1 = c\lambda + O(\lambda^p)$ where c is a constant, $0 < p < 1$, $n = 1, 2, \dots$, $0 \notin \sigma(L_0)$. Suppose that one of the following conditions is satisfied:*

1° $c = 0$; for some $q \in [0, 1]$ we have that $g_k \in \mathcal{D}(L_0^{1-q})$, $f_k \in \mathcal{D}(L_0^q)$ for every k and in this case $\sum_{k=1}^{\infty} s_k \|L_0^{1-q} g_k\| \cdot \|L_0^q f_k\| < \infty$;

2° $c \neq 0$; for some $q \in [0, 2]$ we have that $g_k \in \mathcal{D}(L_0^{2-q})$, $f_k \in \mathcal{D}(L_0^q)$ for every k and in this case $\sum_{k=1}^{\infty} s_k \|L_0^{2-q} g_k\| \cdot \|L_0^q f_k\| < \infty$.

Then the operator L is discrete and there exists a subsequence of natural numbers k_n ($k_n \rightarrow +\infty$) such that

$$\lim_{n \rightarrow \infty} \sum_{k=1}^{k_n} (\mu_k - \lambda_k) = \sum_{k=1}^{\infty} s_k (L_0^{1-q} g_k, L_0^q f_k). \quad (1)$$

Note that formula (1) gives the regularized trace of the operator L because the right-hand side of this equality can be evaluated.

3. Resolvent of the operator L

For the sake of symmetry, we shall introduce new vectors $f'_k = \sqrt{s_k} f_k$, $g'_k = \sqrt{s_k} g_k$. Then

$$S = \sum_{k=1}^{\infty} (\cdot, f'_k) g'_k, \quad (f'_k, f'_l) = s_k \delta_{kl}, \quad (g'_k, g'_l) = s_k \delta_{kl}.$$

In order to find the resolvent of the operator L it is necessary to examine the solvability of the equation

$$L_0 u + \sum_{k=1}^{\infty} (L_0 u, f'_k) g'_k = \lambda u + f \quad (2)$$

where $u \in \mathcal{D}(L_0)$ and $f \in H$. Suppose $\lambda \notin \sigma(L_0)$. Applying the operator $R_\lambda^0 = (L_0 - \lambda I)^{-1}$ to (2) we have that

$$u + \sum_{j=1}^{\infty} (L_0 u, f'_j) R_\lambda^0 g'_j = R_\lambda^0 f.$$

Let us denote $(L_0 u, f'_j) = \xi_j$; then we obtain

$$u = - \sum_{j=1}^{\infty} R_\lambda^0 g'_j + R_\lambda^0 f.$$

Substituting the expression for u in the formula for ξ_j we shall get an infinite system of algebraic linear equations:

$$\xi_k + \sum_{j=1}^{\infty} \xi_j (L_0 R_\lambda^0 g'_j, f'_k) = (L_0 R_\lambda^0 f, f'_k), \quad k = 1, 2, 3, \dots \quad (3)$$

This system can be used to determine ξ_j . Put $a_{kj} = -(L_0 R_\lambda^0 g'_j, f'_k)$ and $c_k = (L_0 R_\lambda^0 f, f'_k)$; then the system (3) can be written in the form $\xi_k - \sum_{j=1}^{\infty} a_{kj} \xi_j = c_k$, $k = 1, 2, 3, \dots$. We can check directly

$$\sum_{k=1}^{\infty} |a_{kk}| < \infty, \quad \sum_{j,k=1}^{\infty} |a_{kj}|^2 < \infty. \quad (4)$$

Formulas (4) imply that the infinite matrix $I - \mathcal{A}$, where $\mathcal{A} = (a_{kj})_{k,j=1}^\infty$ is a Koch matrix. Therefore, by deleting a fixed column and a row from the matrix $I - \mathcal{A}$ we shall get again a Koch matrix.

To the Koch matrix $I - \mathcal{A}$ we can associate its determinant:

$$\det(I - \mathcal{A}) = 1 - \sum_{j=1}^\infty a_{jj} + \frac{1}{2!} \sum_{j,k=1}^\infty \begin{vmatrix} a_{jj} & a_{jk} \\ a_{kj} & a_{kk} \end{vmatrix} - \frac{1}{3!} \sum + \dots$$

(For more details about the Koch matrix see [2] or [3])

Since the matrix of the system $\xi_{ik} - \sum_{j=1}^\infty a_{kj} \xi_j = c_k$, $k = 1, 2, 3, \dots$ is a Koch matrix, the solution of the system (3) is given by

$$\xi_k = \sum_{j=1}^\infty (-1)^{j+k} \Delta_{jk} c_j / \Delta \tag{5}$$

where Δ_{jk} is the determinant of the matrix obtained from the matrix $I - \mathcal{A}$ by deleting its j -th row and k -th column; $\det(I - \mathcal{A}) = \Delta$. Because of $R_\lambda f = u$ ($R_\lambda = (L - \lambda I)^{-1}$) and $u = -\sum_{j=1}^\infty \xi_j R_\lambda^0 g'_j + R_\lambda^0 f$, (5) implies that

$$R_\lambda f = R_\lambda^0 f - \sum_{k=1}^\infty \left(\sum_{j=1}^\infty (-1)^{j+k} \Delta_{jk} (L_0 R_\lambda^0 f, f'_j) R_\lambda^0 g'_k \right) \Delta^{-1}. \tag{6}$$

The spectrum of the operator L is discrete because of Theorem 10.1, p. 336, in [2]. Suppose $\lambda \notin \sigma(L_0) \cup \sigma(L)$; then

$$L - \lambda I = (I + S L_0 R_\lambda^0)(L_0 - \lambda I).$$

Since the operator $L - \lambda I$ vanishes only at the point $x = 0$ it follows that $-1 \notin \sigma(S L_0 R_\lambda^0)$. The operator $S L_0 R_\lambda^0$ is nuclear, so that $(I + S L_0 R_\lambda^0)^{-1} = I + K$, where K is a nuclear operator. But then we have $(L - \lambda I)^{-1} = (L_0 - \lambda I)^{-1}(I + K)$ so that $(L - \lambda I)^{-1}$ is a compact operator and L is a discrete operator.

Now, we introduce linear functionals φ_k :

$$\varphi_k(f) = \Delta^{-1} \sum_{j=1}^\infty (-1)^{j+k} \Delta_{jk} s_k^{-1/2} (L_0 R_\lambda^0 f, f'_j), \quad f \in H. \tag{7}$$

LEMMA 1. *The sequence of linear functionals φ_k ($k = 1, 2, 3, \dots$) on the space H is uniformly bounded.*

Proof. Since

$$\varphi_k(f) = \Delta^{-1} \sum_{j=1}^\infty (-1)^{j+k} \Delta_{jk} s_k^{-1/2} (L_0 R_\lambda^0 f, f'_j)$$

it follows that

$$\varphi_k(f) = \Delta^{-1} \Delta_{kk} (L_0 R_\lambda^0 f, f_k) + \Delta^{-1} \sum_{j \neq k} (-1)^{j+k} \Delta_{jk} s_j^{1/2} s_k^{-1/2} (L_0 R_\lambda^0 f, f_j)$$

and

$$|\varphi_k(f)| \leq |\Delta|^{-1} |\Delta_{kk}| \cdot \|L_0 R_\lambda^0\| \cdot \|f\| \\ + |\Delta|^{-1} \left(\sum_{j \neq k} |\Delta_{jk}|^2 s_j / s_k \cdot \sum_{j \neq k} |(L_0 R_\lambda^0 f, f_j)|^2 \right)^{1/2}.$$

This formulas imply, by using the Bessel inequality, that

$$|\varphi_k(f)| \leq \|f\| \left(|\Delta_{kk}| \cdot |\Delta|^{-1} \cdot \|L_0 R_\lambda^0\| + \|L_0 R_\lambda^0\| \cdot |\Delta|^{-1} \cdot \left(\sum_{j \neq k} |\Delta_{jk}|^2 s_j / s_k \right)^{1/2} \right).$$

We show that $\sup_{j,k} |\Delta_{jk}|/s_k < \infty$ and $\sup_k |\Delta_{kk}| < \infty$. By the Hadamard inequality

$$|\Delta_{jk}|^2 \leq \prod_{p \neq j, p \neq k} \left(\sum_{q \neq k} |\delta_{pq} - a_{pq}|^2 \right) \left(\sum_{q \neq k} |\delta_{kq} - a_{kq}|^2 \right)$$

and by using that $a_{kj} = -(s_k s_j)^{1/2} (L_0 R_\lambda^0 g_j, f_k)$ we get

$$\sum_{q \neq k} |\delta_{kq} - a_{kq}|^2 = \sum_{q \neq k} |a_{kq}|^2 \leq s_k \|L_0 R_\lambda^0\|^2 |S|_1$$

($|S|_1$ is the nuclear norm of operator S). The last formula implies that

$$|\Delta_{jk}|^2 / s_k \leq \|L_0 R_\lambda^0\|^2 \cdot |S|_1 \cdot \prod_{p \neq j, p \neq k} \left(\sum_{q \neq k} |\delta_{pq} - a_{pq}|^2 \right).$$

Because of

$$\prod_{p \neq j, p \neq k} \left(\sum_{q \neq k} |\delta_{pq} - a_{pq}|^2 \right) \leq \exp \left(2 \sum_{p=1}^{\infty} |a_{pp}| + \sum_{p,q=1}^{\infty} |a_{pq}|^2 \right) = M < \infty$$

we have $|\Delta_{jk}|^2 / s_k \leq M |S|_1 \cdot \|L_0 R_\lambda^0\|^2$.

We can show in a similar way that $|\Delta_{kk}| \leq \sqrt{M}$. This completes the proof of Lemma 1.

Using (7), the equality (6) can be written in the form

$$R_\lambda - R_\lambda^0 = - \sum_{k=1}^{\infty} s_k \varphi_k(\cdot) R_\lambda^0 g_k.$$

Since Lemma 1 implies $\sum_{k=1}^{\infty} s_k \|\varphi_k\| \|R_\lambda^0 g_k\| < \infty$, the operator

$$\sum_{k=1}^{\infty} s_k \varphi_k(\cdot) R_\lambda^0 g_k$$

is nuclear. So we have

$$\text{Sp}(R_\lambda - R_\lambda^0) = - \sum_{k=1}^{\infty} s_k \varphi_k(R_\lambda^0 g_k)$$

and

$$\text{Sp}(R_\lambda - R_\lambda^0) = -\Delta^{-1} \sum_{k,j=1}^{\infty} (-1)^{j+k} \Delta_{jk} (L_0(R_\lambda^0)^2 g'_k, f'_j). \quad (8)$$

Using that $a_{jk} = -(L_0 R_\lambda^0 g'_k, f'_j)$, we can get $a'_{jk}(\lambda) = -(L_0(R_\lambda^0)^2 g'_k, f'_j)$ and (8) can be written as

$$\text{Sp}(R_\lambda - R_\lambda^0) = \Delta^{-1} \sum_{k,j=1}^{\infty} (-1)^{j+k} \Delta_{jk} a'_{jk}(\lambda). \quad (9)$$

For the Koch matrix we have

$$\sum_{j,k=1}^{\infty} (-1)^{j+k} \Delta_{jk} a'_{jk}(\lambda) = \Delta'(\lambda),$$

so (9) implies

$$\text{Sp}(R_\lambda - R_\lambda^0) = \Delta'(\lambda)/\Delta(\lambda). \quad (10)$$

4. The proof of the theorem in the case when $c = 0$ in the formula for $N(\lambda)$

Let $\{\lambda_k\}$ be the eigenvalues of the operator L_0 and $\{\psi_k\}$ an orthonormal base in the space H , where ψ_k are eigenvectors (corresponding to λ_k) of the operator L_0 . We denote $c_n^j = (g_j, \psi_n)$, $b_n^j = (f_j, \psi_n)$; then $g_k = \sum_{n=1}^{\infty} c_n^k \psi_n$, $f_j = \sum_{n=1}^{\infty} b_n^j \psi_n$. The spectral theorem implies

$$(L_0 R_\lambda^0 g_k, f_j) = \sum_{n=1}^{\infty} \frac{\lambda_n c_n^k \overline{b_n^j}}{\lambda - \lambda_n}. \quad (11)$$

Since the assumptions are $g_k \in \mathcal{D}(L_0^{1-q})$, $f_j \in \mathcal{D}(L_0^q)$, the Bessel inequality implies

$$\sum_{n=1}^{\infty} \lambda_n |c_n^k b_n^j| \leq \|L_0^{1-q} g_k\| \|L_0^q f_j\|. \quad (12)$$

For the case $c = 0$ we have $N(\lambda) = O(\lambda^p)$ ($0 < p < 1$) and then there exists a sequence of real numbers r_n ($r_n \rightarrow \infty$) such that $d_n \rightarrow \infty$, where $d_n = d(\Gamma_n, \sigma(L_0))$ is the distance between the circle $\Gamma_n = \{\lambda : |\lambda| = r_n\}$ and the spectrum of the operator L [1].

LEMMA 2. *The operators L and L_0 have the same number of eigenvalues inside the circle Γ_n , for n large enough.*

Proof. We multiply equality (10) by $1/2\pi i$ and integrate it over Γ_n . Note that the well-known Riesz theorem implies

$$\frac{1}{2\pi i} \int_{\Gamma_n} \text{Sp}(R_\lambda - R_\lambda^0) d\lambda = N_1 - N_2,$$

where N_1 and N_2 denote the numbers of eigenvalues inside the contour Γ_n (counting their multiplicity) of the operators L and L_0 , respectively. So, this implies that

$$N_1 - N_2 = \frac{1}{2\pi i} \int_{\Gamma_n} \frac{\Delta'(\lambda)}{\Delta(\lambda)} d\lambda = \frac{1}{2\pi i} \int_{\Gamma_n} \frac{S'(\lambda)}{1 + S(\lambda)} d\lambda \quad (13)$$

where $S(\lambda) = \Delta(\lambda) - 1$. Let us estimate $S(\lambda)$. Since $a_{jk} = -(s_j s_k)^{1/2} (L_0 R_\lambda^0 g_k, f_j)$ by using (11) and (12) we have

$$|a_{jk}| \leq (s_j s_k)^{1/2} \|L_0^{1-q} g_k\| \cdot \|L_0^q f_j\| / d(\lambda) \quad (14)$$

where $d(\lambda)$ denotes the distance between a point λ and $\sigma(L_0)$. Note that for $\lambda \in \Gamma_n$ we have $d(\lambda) \geq d_n$. The definition of the function $S(\lambda)$ and (14) imply the following estimate

$$\left. \begin{aligned} |S(\lambda)| &\leq \sum_{k=1}^{\infty} M(\lambda)^k \\ \text{where } M(\lambda) &= \sum_{j=1}^{\infty} s_j \|L_0^q f_j\| \cdot \|L_0^{1-q} g_j\| / d(\lambda) \end{aligned} \right\} \quad (15)$$

From $\sum_{j=1}^{\infty} s_j \|L_0^q f_j\| \|L_0^{1-q} g_j\| < \infty$, and (15) it follows that $S(\lambda) \rightarrow 0$ when $n \rightarrow \infty$ and $\lambda \in \Gamma_n$; the convergence is uniform with respect to λ . Then (13) implies

$$N_1 - N_2 = \sum_{\nu=0}^{\infty} (-1)^\nu \frac{1}{2\pi i} \int_{\Gamma_n} S'(\lambda) S^\nu(\lambda) d\lambda. \quad (16)$$

From (15) we get

$$|S^\nu(\lambda)| \leq K_1^\nu / d(\lambda)^\nu \quad (17)$$

for $\lambda \in \Gamma_n$ and n is large enough (K_1 does not depend on λ and ν).

By differentiating (11) with respect to λ and using (12) we have

$$|a'_{jk}(\lambda)| \leq (s_j s_k)^{1/2} \|L_0^{1-q} g_k\| \cdot \|L_0^q f_j\| / d^2(\lambda). \quad (18)$$

From (14) and (18) and the definition of the function $S(\lambda)$ we can get the following estimate

$$|S'(\lambda)| \leq \frac{1}{d(\lambda)} \cdot \sum_{n=1}^{\infty} n M^n(\lambda) \leq \frac{K_2}{d^2(\lambda)} \quad (19)$$

where the constant K_2 does not depend on λ , $\lambda \in \Gamma_n$ and n is large enough. Then (17) and (19) imply

$$\left| \int_{\Gamma_n} S'(\lambda) S^\nu(\lambda) d\lambda \right| \leq K_2 K_1^\nu \int_{\Gamma_n} \frac{|d\lambda|}{d^{2+\nu}(\lambda)}.$$

Next we can easily check that

$$\int_{\Gamma_n} \frac{|d\lambda|}{d^2(\lambda)} \leq \frac{C_0}{d_n}$$

(the constant C_0 does not depend on n). Because of $d(\lambda) \geq d_n$ on Γ_n the inequality

$$\int_{\Gamma_n} \frac{|d\lambda|}{d^{2+\nu}(\lambda)} \leq \frac{1}{d_n^\nu} \int_{\Gamma_n} \frac{|d\lambda|}{d^2(\lambda)}$$

holds and consequently

$$\left| \int_{\Gamma_n} S'(\lambda) S^\nu(\lambda) d\lambda \right| \leq C_0 K_2 K_1^\nu / d_n^{\nu+1} \quad (20)$$

for n large enough. Formula (16) implies $|N_1 - N_2| \leq \sum_{\nu=0}^{\infty} C_0 K_2 K_1^\nu / d_n^{\nu+1}$ (n large enough) and $|N_1 - N_2| = O(1/d_n)$. Since $|N_1 - N_2|$ is a natural number or 0, this is possible if and only if $N_1 = N_2$. So the lemma is proved.

Let us now prove the Theorem in the case $c = 0$. If we multiply (10) by $\lambda/2\pi i$ and integrate it over Γ_n we have

$$\frac{1}{2\pi i} \int_{\Gamma_n} \lambda \operatorname{Sp}(R_\lambda - R_\lambda^0) d\lambda = \frac{1}{2\pi i} \int_{\Gamma_n} \lambda \frac{\Delta'(\lambda)}{\Delta(\lambda)} d\lambda. \quad (21)$$

The properties of Riesz projectors imply that

$$\frac{1}{2\pi i} \int_{\Gamma_n} \lambda \operatorname{Sp}(R_\lambda - R_\lambda^0) d\lambda = \sum_{\nu=1}^{k_n} (\mu_\nu - \lambda_\nu) \quad (22)$$

where k_n denotes the number of eigenvalues for operators L and L_0 inside the contour Γ_n (this number is the same for both operators by Lemma 2). From (21) and (22) it follows that

$$\sum_{\nu=1}^{k_n} (\mu_\nu - \lambda_\nu) = \frac{1}{2\pi i} \int_{\Gamma_n} \lambda \frac{\Delta'(\lambda)}{\Delta(\lambda)} d\lambda. \quad (23)$$

Since

$$\frac{1}{2\pi i} \int_{\Gamma_n} \lambda \frac{\Delta'(\lambda)}{\Delta(\lambda)} d\lambda = \sum_{\nu=0}^{\infty} (-1)^\nu \frac{1}{2\pi i} \int_{\Gamma_n} \lambda S'(\lambda) S^\nu(\lambda) d\lambda \quad (24)$$

for n large enough (because of $S(\lambda) \rightrightarrows 0$, as $n \rightarrow \infty$, $\lambda \in \Gamma_n$), it is enough to estimate each of the integrals $\int_{\Gamma_n} \lambda S'(\lambda) S^\nu(\lambda) d\lambda$.

By the definition of the function S (because of the structure of a_{jk}) we know that S is a holomorphic function on \mathbf{C} except at the points $\lambda_1, \lambda_2, \dots$. Inside the contour Γ_n , the function S has singularities at the points $\lambda_1, \lambda_2, \dots, \lambda_{k_n}$. We can easily check that

$$\int_{\Gamma_n} \lambda S'(\lambda) S^\nu(\lambda) d\lambda = -\frac{1}{\nu+1} \int_{\Gamma_n} S^{\nu+1}(\lambda) d\lambda. \quad (25)$$

Formulas (24) and (25) imply

$$\frac{1}{2\pi i} \int_{\Gamma_n} \lambda \frac{\Delta'(\lambda)}{\Delta(\lambda)} d\lambda = -\frac{1}{2\pi i} \int_{\Gamma_n} S(\lambda) d\lambda + \sum_{\nu=1}^{\infty} \frac{(-1)^{\nu+1}}{\nu+1} \frac{1}{2\pi i} \int_{\Gamma_n} S^{\nu+1}(\lambda) d\lambda. \quad (26)$$

From (17), for n large enough, it follows that

$$\left| \frac{1}{2\pi i} \int_{\Gamma_n} S^{\nu+1}(\lambda) d\lambda \right| \leq \frac{C_0 K_1}{2\pi} \left(\frac{K_1}{d_n} \right)^\nu,$$

therefore

$$\sum_{\nu=1}^{\infty} \frac{(-1)^{\nu+1}}{\nu+1} \frac{1}{2\pi i} \int_{\Gamma_n} S^{\nu+1}(\lambda) d\lambda = O(d_n^{-1}). \quad (27)$$

Next, (26) and (27) imply

$$\frac{1}{2\pi i} \int_{\Gamma_n} \lambda \frac{\Delta'(\lambda)}{\Delta(\lambda)} d\lambda = -\frac{1}{2\pi i} \int_{\Gamma_n} S(\lambda) d\lambda + O(d_n^{-1}). \quad (28)$$

Let us consider the function R given by

$$R(\lambda) = S(\lambda) + \sum_{j=1}^{\infty} a_{jj}(\lambda). \quad (29)$$

In a way similar to the proof of inequality (15), we can prove that $|R(\lambda)| \leq P \cdot (d(\lambda))^{-2}$ for $\lambda \in \Gamma_n$, for n large enough and a constant P which does not depend on λ . This inequality implies

$$-\frac{1}{2\pi i} \int_{\Gamma_n} R(\lambda) d\lambda = O(d_n^{-1}). \quad (30)$$

Then, from (28), (29) and (30) it follows that

$$\frac{1}{2\pi i} \int_{\Gamma_n} \lambda \frac{\Delta'(\lambda)}{\Delta(\lambda)} d\lambda = \frac{1}{2\pi i} \int_{\Gamma_n} \left(\sum_{j=1}^{\infty} a_{jj}(\lambda) \right) d\lambda + O(d_n^{-1}). \quad (31)$$

Since the series $\sum_{j=1}^{\infty} a_{jj}(\lambda)$ converges uniformly on Γ_n , we have

$$\frac{1}{2\pi i} \int_{\Gamma_n} \lambda \frac{\Delta'(\lambda)}{\Delta(\lambda)} d\lambda = \sum_{j=1}^{\infty} \frac{1}{2\pi i} \int_{\Gamma_n} a_{jj}(\lambda) d\lambda + O(d_n^{-1}). \quad (32)$$

Because of $a_{jj} = -s_j \sum_{\nu=1}^{\infty} \lambda_{\nu} a_{\nu}^j \overline{b_{\nu}^j} / (\lambda_{\nu} - \lambda)$, (32) gives

$$\frac{1}{2\pi i} \int_{\Gamma_n} \lambda \frac{\Delta'(\lambda)}{\Delta(\lambda)} d\lambda = \sum_{j=1}^{\infty} \left(s_j \sum_{\nu=1}^{k_n} \lambda_{\nu} a_{\nu}^j \overline{b_{\nu}^j} \right) + O(d_n^{-1}). \quad (33)$$

From (23) and (24) we can get

$$\sum_{\nu=1}^{k_n} (\mu_{\nu} - \lambda_{\nu}) = \sum_{j=1}^{\infty} \left(s_j \sum_{\nu=1}^{k_n} \lambda_{\nu} a_{\nu}^j \overline{b_{\nu}^j} \right) + O(d_n^{-1}). \quad (34)$$

Since $\sum_{\nu=1}^{\infty} \lambda_{\nu} |a_{\nu}^j b_{\nu}^j| \leq \|L_0^{1-q} g_j\| \cdot \|L_0^q f_j\|$ the series $\sum_{j=1}^{\infty} s_j \|L_0^{1-q} g_j\| \cdot \|L_0^q f_j\|$ is convergent and $\lim_{n \rightarrow \infty} \sum_{\nu=1}^{k_n} \lambda_{\nu} a_{\nu}^j \overline{b_{\nu}^j} = (L_0^{1-q} g_j, L_0^q f_j)$, then we have from (34), when $n \rightarrow \infty$,

$$\lim_{n \rightarrow \infty} \sum_{\nu=1}^{k_n} (\mu_{\nu} - \lambda_{\nu}) = \sum_{j=1}^{\infty} s_j (L_0^{1-q} g_j, L_0^q f_j).$$

So, the theorem is proved in the case $c = 0$.

**5. The proof of the theorem in the case
when $c \neq 0$ in formula for $N(\lambda)$**

If $g_k \in \mathcal{D}(L_0^{2-q})$ the spectral theorem implies $g_k \in \mathcal{D}(L_0^{1-q})$ and

$$\|L_0^{1-q} g_k\| \leq \lambda_1^{-1} \|L_0^{2-q} g_k\|. \quad (35)$$

Because of the assumption $\sum_{k=1}^{\infty} s_k \|L_0^{2-q} g_k\| \cdot \|L_0^q f_k\| < \infty$ (which holds in the case $c \neq 0$) (35) implies

$$\sum_{k=1}^{\infty} s_k \|L_0^{1-q} g_k\| \cdot \|L_0^q f_k\| < \infty.$$

If $g_k \in \mathcal{D}(L_0^{2-q})$, $f_j \in \mathcal{D}(L_0^q)$ using the Bessel inequality, we have

$$\sum_{n=1}^{\infty} \lambda_n^2 |c_n^k \overline{b_n^j}| \leq \|L_0^{2-q} g_k\| \cdot \|L_0^q f_j\|. \quad (36)$$

In the case $c \neq 0$, the formula for $N(\lambda)$ can be written in the form

$$\lambda_n = n(k + o(1)) \quad (37)$$

where $k = c^{-1}$.

The following lemma was proved in [1].

LEMMA 3. *Suppose that $\alpha > 0$, N_α is so large that $N_\alpha > 8/\alpha$ and for $o(1)$ in (37) the following holds: $|o(1)| \leq \alpha k/8(1+\alpha)$, for every $n > N_\alpha$. Then between λ_n and $\lambda_{n'}$, where $n > N_\alpha$ and $n' = [n(1+\alpha)] > N_\alpha$, there are eigenvalues λ_{n_ν} and $\lambda_{n_\nu+1}$ of the operator L_0 such that*

$$\lambda_{n_\nu+1} - \lambda_{n_\nu} \geq k/2. \quad (38)$$

Let us define $\Gamma_{n_\nu} = \{\lambda : |\lambda| = r_{n_\nu}\}$, where $r_{n_\nu} = (\lambda_{n_\nu} + \lambda_{n_\nu+1})/2$. Since

$$(L_0 R_\lambda^0 g_k, f_j) = \sum_{n=1}^{\infty} \frac{\lambda_n c_n^k \overline{b_n^j}}{\lambda_n - \lambda},$$

we have

$$(L_0 R_\lambda^0 g_k, f_j) + \frac{1}{\lambda} (L_0^{1-q} g_k, L_0^q f_j) = \sum_{n=1}^{\infty} \frac{\lambda_n^2 c_n^k \overline{b_n^j}}{\lambda(\lambda_n - \lambda)}. \quad (39)$$

From (35), (36), and (39) it follows that

$$|(L_0 R_\lambda^0 g_k, f_j)| \leq \frac{\|L_0^{2-q} g_k\| \cdot \|L_0^q f_j\|}{|\lambda| d(\lambda)} + \frac{1}{\lambda_1} \frac{\|L_0^{2-q} g_k\| \cdot \|L_0^q f_j\|}{|\lambda|}$$

and then, $d(\lambda) \geq k/4$ on Γ_{n_ν} , implies

$$|a_{jk}(\lambda)| \leq K (s_j s_k)^{1/2} \cdot \|L_0^{2-q} g_k\| \cdot \|L_0^q f_j\| \cdot |\lambda|^{-1} \quad (40)$$

on Γ_{n_ν} ($K = 4 \cdot k^{-1} + \lambda_1^{-1}$). Using (40), we can get an estimate for the function

$$|S(\lambda)| \leq \sum_{n=1}^{\infty} M_1^n(\lambda) \quad (41)$$

where $M_1(\lambda) = |\lambda|^{-1} K \sum_{j=1}^{\infty} s_j \|L_0^{2-q} g_j\| \cdot \|L_0^q f_j\|$.

Clearly, for $\lambda \in \Gamma_{n_\nu}$ and ν large enough, we have $|S(\lambda)| \leq Q/|\lambda|$ where the constant Q does not depend on λ ($\lambda \in \Gamma_{n_\nu}$). This implies that $S(\lambda) \rightarrow 0$, for $\lambda \in \Gamma_{n_\nu}$, $\nu \rightarrow \infty$, therefore the function $\Delta(\lambda)$ does not have zeros in region $\{\lambda : |\lambda| > r_{n_\nu}\}$ for ν large enough.

In a similar way, as it was done in Lemma 2, we can show that the operators L and L_0 have the same number of eigenvalues inside the contour Γ_{n_ν} , for ν large enough. This number is equal to n_ν .

As in the previous case, we have

$$\sum_{k=1}^{n_\nu} (\mu_k - \lambda_k) = -\frac{1}{2\pi i} \int_{\Gamma_{n_\nu}} S(\lambda) d\lambda + \sum_{r=1}^{\infty} \frac{(-1)^{r+1}}{r+1} \frac{1}{2\pi i} \int_{\Gamma_{n_\nu}} S(\lambda)^{r+1} d\lambda. \quad (42)$$

Using the estimate $|S(\lambda)| \leq Q/|\lambda|$, which holds for $\lambda \in \Gamma_{n_\nu}$ and ν large enough, we can get easily

$$\sum_{r=1}^{\infty} \frac{(-1)^{r+1}}{r+1} \frac{1}{2\pi i} \int_{\Gamma_{n_\nu}} S(\lambda)^{r+1} d\lambda = O(r_{n_\nu}^{-1}), \quad \nu \rightarrow \infty$$

and consequently (42) implies

$$\sum_{k=1}^{n_\nu} (\mu_k - \lambda_k) = -\frac{1}{2\pi i} \int_{\Gamma_{n_\nu}} S(\lambda) d\lambda + O(r_{n_\nu}^{-1}). \quad (43)$$

Since $|S(\lambda) + \sum_{j=1}^{\infty} a_{jj}(\lambda)| \leq \text{const}/|\lambda|^2$ on Γ_{n_ν} (ν large enough) we have

$$\frac{1}{2\pi i} \int_{\Gamma_{n_\nu}} \left(S(\lambda) + \sum_{j=1}^{\infty} a_{jj}(\lambda) \right) d\lambda = O(r_{n_\nu}^{-1}). \quad (44)$$

From (43) and (44) it follows that

$$\sum_{k=1}^{n_\nu} (\mu_k - \lambda_k) = \frac{1}{2\pi i} \int_{\Gamma_{n_\nu}} \sum_{j=1}^{\infty} a_{jj}(\lambda) d\lambda + O(r_{n_\nu}^{-1}). \quad (45)$$

Because of

$$\frac{1}{2\pi i} \int_{\Gamma_{n_\nu}} a_{jj}(\lambda) d\lambda = s_j \sum_{s=1}^{n_\nu} \lambda_s a_s^j \bar{b}_s^j$$

and the uniform convergence of the series $\sum a_{jj}(\lambda)$ on Γ_{n_ν} , (45) becomes

$$\sum_{k=1}^{n_\nu} (\mu_k - \lambda_k) = \sum_{j=1}^{\infty} \left(s_j \sum_{s=1}^{n_\nu} \lambda_s a_s^j \bar{b}_s^j \right) + O(r_{n_\nu}^{-1}). \quad (46)$$

Hence

$$\sum_{j=1}^{\infty} s_j \sum_{s=1}^{\infty} \lambda_s |a_s^j b_s^j| \leq \sum_{j=1}^{\infty} s_j \|L_0^{1-q} g_j\| \cdot \|L_0^q f_j\| < \infty$$

(46) and the fact that $\lim_{\nu \rightarrow \infty} \sum_{s=1}^{n_\nu} (\mu_s - \lambda_s) = (L_0^{1-q} g_j, L_0^q f_j)$ imply

$$\lim_{\nu \rightarrow \infty} \sum_{s=1}^{n_\nu} (\mu_s - \lambda_s) = \sum_{j=1}^{\infty} s_j (L_0^{1-q} g_j, L_0^q f_j).$$

So, the proof of the theorem in case $c \neq 0$ is completed.

Remark. Under the condition stated in the theorem can get directly that

$$\sum_{j=1}^{\infty} s_j (L_0^{1-q} g_j, L_0^q f_j) = \sum_{j=1}^{\infty} \lambda_j (S\varphi_j, \varphi_j) \quad (47)$$

so the statement of the theorem can be formulated in terms of the operator S by formula (47).

Example. Let L_0 be a semibounded ($L_0 > 0$), discrete, selfadjoint operator which acts on H . Suppose that $S > 0$ is a nuclear operator which has properties $R(S) \subset \mathcal{D}(L_0)$ and $\sum \lambda_n (S\varphi_n, \varphi_n) < \infty$ (λ_n denotes the eigenvalues and φ_n the corresponding eigenvectors of the operator L_0). Then

$$\lim_{n \rightarrow \infty} \sum_{\nu=1}^{k_n} (\mu_\nu - \lambda_\nu) = \sum_{\nu=1}^{\infty} \lambda_\nu (S\varphi_\nu, \varphi_\nu).$$

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