

## A GEOMETRIC CHARACTERIZATION OF HELICODIAL SURFACES OF CONSTANT MEAN CURVATURE

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**Abstract.** We prove that a helicodial surface has constant mean curvature if and only if its principal axes make an angle constant with the orbits. Moreover, the arguments used lead to a simple proof of the fact that all helicodial surfaces with constant mean curvature  $H$  can be isometrically deformed, through helicodial surfaces of the same  $H$ , into surfaces of revolution of the same  $H$  (Delaunay surfaces).

**1. Introduction.** A one parameter subgroup of the group of rigid motions of  $E^3$  is a differentiable mapping  $\gamma : \mathbf{R} \times E^3 \rightarrow E^3$  with the following properties:

a) The map  $\gamma_t : E^3 \rightarrow E^3$  given by  $x \rightarrow \gamma(t, x)$ ,  $x \in E^3$ ,  $t \in \mathbf{R}$  is a rigid motion,

b)  $\gamma_t \circ \gamma_s = \gamma_{t+s}$ ;      c)  $\gamma_0 =$  the identity.

Let  $x \in E^3$  have coordinates  $x = (x^1, x^2, x^3)$ . It may be shown that, possibly after a change of basis, any one parameter subgroup may be written either as

$$\gamma(t, x) = (x^1 \cos t + x^2 \sin t, -x^1 \sin t + x^2 \cos t, x^3 + bt),$$

where  $-\infty < b < +\infty$  is constant, or as

$$\gamma(t, x) = (x^1, x^2, x^3) + t(0, 0, 1) = (x^1, x^2, x^3 + t).$$

In the former case, if  $b \neq 0$ ,  $\gamma(t, x)$  is called a *helicodial motion with axis the  $x^3$ -axis and pitch  $b$* . The orbit  $t \in \mathbf{R} \rightarrow \gamma(t, x) \in E^3$  of a point  $x = (x^1, x^2, x^3)$  which does not lie on the  $x^3$ -axis is a helix. All such helices have the  $x^3$ -axis as common axis. If  $b = 0$ ,  $\gamma(t, x)$  is a *rotational motion about the  $x^3$ -axis*. The orbits of points not lying on the  $x^3$ -axis are circles having the  $x^3$ -axis as common axis. In the latter case,  $\gamma(t, x)$  is called a *parallel translational motion in the direction of the  $x^3$ -axis*. All orbits are straight lines parallel to the  $x^3$ -axis.

A helicoidal surface with axis the  $x^3$ -axis and pitch  $b \neq 0$  is surface that is invariant by  $\gamma(t, x)$  for all  $t$ . When  $b = 0$  the surface reduces to a *surface of revolution*. Finally, the translational motions generate the *cylinders*.

If we consider a curve  $c(s)$  on any of these surfaces which intersects all the orbits without touching them, the surface can be parametrized by  $(s, t)$  as  $\gamma(t, c(s))$ .

The main result states:

**THEOREM 4.** *A helicoidal surface has constant mean curvature if and only if its principal axes make an angle constant with the orbits.*

This is not true for the surfaces of revolution and the cylinders since, regardless of  $H$ , the orbits are principal curves.

In [3], an analytic parametrization of the helicoidal surfaces of constant mean curvature was exhibited. Also, in [3] was first shown that these surfaces can be isometrically deformed under preservation of the mean curvature and through helicoidal surfaces into Delaunay surfaces. Here, a simpler proof of this fact is given by making use of facts for general (not necessarily with  $H = \text{constant}$ ) helicoidal surfaces. The main tool for the proof of this part is the following:

**THEOREM (O. Bonnet)** (cf. [1, 2, 4]) *A surface of constant mean curvature in  $E^3$ , other than the plane and the sphere, can be isometrically deformed so that the mean curvature is preserved. During this deformation the principal directions rotate by a fixed angle, and for any fixed angle as rotation angle of the principal directions a surface of this isometric deformation is obtained.*

**2. Some local surface theory.** We consider a surface  $M^2$  in  $E^3$ , orientable and sufficiently smooth. We consider a well defined field of orthonormal frames  $(x, e'_1, e'_2, e'_3)$  over  $M^2$ , such that  $x \in M$ , and  $e_1, e_2$  comprise an orthonormal basis of the tangent space of  $M$  at  $x$ . We have then

$$\eta_i = dx \cdot e'_i, \quad \eta_{ij} = de'_i \cdot e'_j, \quad \eta_{ij} = -\eta_{ji} \text{ (so } \eta_{ii} = 0)$$

$d\eta_i = \sum_{j=1}^3 \eta_{ij} \wedge \eta_j$  (1-st structural equation),  $d\eta_{ij} = \sum_{k=1}^3 \eta_{ik} \wedge \eta_{kj}$  (2nd structural equation) where  $1 \leq i, j \leq 3$ . On  $M^2$ ,  $\eta_3 = 0$  so we have  $\eta_{13} \wedge \eta_1 + \eta_{23} \wedge \eta_2 = 0$ . So, by Cartan's Lemma we get

$$\eta_{13} = a\eta_1 + b\eta_2, \quad \eta_{23} = b\eta_1 + c\eta_2. \quad (1)$$

Then the mean and Gaussian curvatures of  $M^2$  are

$$H = (a + c)/2, \quad K = ac - b^2.$$

We also have

$$\left. \begin{aligned} d\eta_{12} &= -K\eta_1 \wedge \eta_2 && \text{(Gauss Equation) (GE).} \\ d\eta_{13} &= \eta_{12} \wedge \eta_3 = -bd\eta_2 + cd\eta_1 \\ d\eta_{23} &= \eta_{21} \wedge \eta_{13} = ad\eta_2 - bd\eta_1 \end{aligned} \right\} \text{(Codazzi-Mainardi Equations) (CME).}$$

A given Riemannian surface can be realized in  $E^3$  if the CME are satisfied.

We now let

$$e_1 = \cos \psi e_1 - \sin \psi e_2, \quad e_2 = \sin \psi e_1 + \cos \psi e_2$$

be the principal frame of  $M^2$ . For this frame the function  $b$  defined by (1) is zero and  $a, c$  are the principal curvatures.

In the sequel, we consider  $M^2$  with no umbilic points. We may then assume for the principal curvatures  $k_1, k_2$  of  $M^2$  that  $k_1 > k_2$  and we put  $J = k_1 - k_2 > 0$ . We will show that the CME are equivalent to:

$$dH = H_1 \eta_1 + H_2 \eta_2, \quad (\text{thus defining } H_1, H_2) \quad (2)$$

$$d\psi = -\cos 2\psi [H_2 J^{-1} \eta_1 + H_1 J^{-1} \eta_2] = \sin 2\psi [-H_1 J^{-1} \eta_1 + H_2 J^{-1} \eta_2] + \quad (3)$$

$$+ 1/2 * d \log J + \eta_{12},$$

where  $*$  is the Hodge operator whose action on the 1-forms is described by

$$*\eta_1 = \eta_2, \quad *\eta_2 = -\eta_1 \quad (*^2 = -1).$$

*Proof.* The principal coframe is

$$\omega_1 = \cos \psi \eta_1 - \sin \psi \eta_2, \quad \omega_2 = \sin \psi \eta_1 + \cos \psi \eta_2.$$

Exterior differentiation gives

$$d\omega_1 = (-d\psi + \eta_{12}) \wedge \omega_2, \quad d\omega_2 = \omega_1 \wedge (-d\psi \wedge \eta_{12}).$$

Thus, the connection form associated to the principal coframe is  $\omega_{12} = -d\psi + \eta_{12}$ , so that  $d\psi = -\omega_{12} + \eta_{12}$ . We then need to show that

$$\omega_{12} = \cos 2\psi [H_2 J^{-1} \eta_1 + H_1 J^{-1} \eta_2] - \sin 2\psi [-H_1 J^{-1} \eta_1 + H_2 J^{-1} \eta_2] - 1/2 * d \log J.$$

We work with  $\omega_1, \omega_2$  and the associated forms  $\omega_{12}, \omega_{13}, \omega_{23}$ . We have that  $\omega_{13} = k_1 \omega_1$  and  $\omega_{23} = k_2 \omega_2$ . We put  $\omega_{12} = p\omega_1 + q\omega_2$ . Then the equations  $d\omega_{13} = \omega_{12} \wedge \omega_{23}, d\omega_{23} = \omega_{21} \wedge \omega_{13}$  imply:

$$[dk_1 - (k_1 - k_2)p\omega_2] \wedge \omega_1 = 0, \quad [dk_2 - (k_1 - k_2)q\omega_1] \wedge \omega_2 = 0.$$

We have  $J = k_1 - k_2 > 0$  and we set  $dH = d(k_1 + k_2)/2 = u\omega_1 + v\omega_2$ . Then we get

$$dk_1 = (2u - Jq)\omega_1 + Jp\omega_2, \quad dk_2 = Jq\omega_1 + (2v - Jp)\omega_2.$$

By subtracting the second relation from the first and dividing by  $J$  we have

$$d \log J = 2uJ^{-1}\omega_1 - 2vJ^{-1}\omega_2 + 2(-q\omega_1 + p\omega_2).$$

But  $-q\omega_1 + p\omega_2 = *\omega_{12}$ , thus we obtain

$$\omega_{12} = vJ^{-1}\omega_1 + uJ^{-1}\omega_2 - 1/2 * d \log J.$$

Now by comparing  $dH$  given by (2) and  $dH$  given by  $dH = u\omega_1 + v\omega_2$ , with  $\omega_1, \omega_2$  expressed in terms of  $\eta_1, \eta_2$  we get

$$u = H_1 \cos \psi - H_2 \sin \psi, \quad v = H_1 \sin \psi + H_2 \cos \psi.$$

Now convert the last expression for  $\omega_{12}$  to the desired final form by inserting the above relations for  $u, v$  and by again rewriting  $\omega_1, \omega_2$  in terms of  $\eta_1, \eta_2$ .

The converse direction of the equivalence follows by reversing the process of this computation  $\square$

*Remark 1.* The Theorem of 0. Bonnet quoted in section 1 follows from (3) and the fundamental theorem of surfaces. With  $H$  constant on  $M^2$  we get from (3) that

$$d\psi = -1/2 * d \log J + \eta_{12}.$$

Thus if some  $\psi$  satisfies this equation, so does  $\psi + \text{constant}$ , and the fundamental theorem of surfaces applies for each new  $\psi$ .

**3. Some facts about Helicoidal Surfaces.** The geometry of a helicoidal surface allows us to parametrize it by  $(s, t)$ , where

$s$  =arc-length of curves orthogonal to orbits measured from a fixed orbit,

$t$  =time along orbits from a fixed  $t = t_0$ , (see also [3]). Then the curves  $t = \text{constant}$  are carried along the orbits by the motion. They remain orthogonal to the orbits and foliate the surface. So an orthonormal frame  $(e'_1, e'_2)$  is determined along these coordinate curves with  $e'_2$  tangent to the orbits. The corresponding coframe may be written as

$$\eta_1 = ds, \quad \eta_2 = q(s)dt \quad (q \text{ depends only on } s).$$

Thus,

$$\eta_{12} = \frac{q'(s)}{q(s)}\eta_2 = \mu(s)\eta_2.$$

Hence the  $\eta_1$ -curves are geodesics and  $\eta_2$ -curves (orbits) have geodesic curvature equal to

$$\mu(s) = \frac{d}{ds} \log(|q(s)|).$$

Along each orbit  $a, c, \mu, \psi$  are constant. So, in this case we get  $H_2 = 0$ . Also if we put  $dJ = J_1\eta_1 + J_2\eta_2$ , we get  $J_2 = 0$  and  $d \log J = J_1 J^{-1} \eta_1$ . Hence relation (3) becomes

$$d\psi = -\cos 2\psi[H_1 J^{-1} \eta_2] - \sin 2\psi[H_1 J^{-1} \eta_1] + 1/2 J_1 J^{-1} \eta_2 + \mu \eta_2.$$

Since  $\psi = \psi(s)$ , this implies

$$\frac{d\psi}{ds} = -\sin 2\psi \frac{dH/ds}{J}, \quad \mu = \cos 2\psi \frac{dH/ds}{J} - \frac{1}{2} \frac{dJ/ds}{J}. \quad (4)$$

By direct computation or by well known facts about curves on surfaces, we get the following:

The space curvature of orbits is

$$\sqrt{\mu^2 + (k_1 \cos^2 \psi + k_2 \sin^2 \psi)^2} \text{ or } \sqrt{\mu^2 + (k_1 \sin^2 \psi + k_2 \cos^2 \psi)^2},$$

and the space torsion of orbits is  $\pm(k_1 - k_2) \sin \psi \cos \psi$ , (5) for the respective cases when  $e_1$  is the major and minor principal direction.

Finally, we show that a helicoidal surface of constant mean curvature is free of umbilic points. This fact allows us to apply the previous local theory everywhere to such a surface.

*Proof.* Let  $M^2$  be an oriented connected surface in  $E^3$  with unit normal vector field  $v$ . We consider isothermal parameters  $(x_1, x_2)$  on  $M^2$ . If  $M^2$  has constant mean curvature then the (locally defined) function

$$f(z) = (b_{11} - b_{22} - 2ib_{12}(z)), \quad z = x_1 + ix_2$$

is known to be holomorphic in  $z$ ; here  $b_{ij} = -\langle \nabla_{\partial/\partial x_i} v, \partial/\partial x_j \rangle$ ,  $1 \leq i, j \leq 2$ , are the components of the second fundamental form of  $M^2$ . The zeros of  $f(z)$  are exactly the umbilic points of  $M^2$ . So, if an umbilic point is not isolated then  $f(z) \equiv 0$  and thus  $M^2$  is totally umbilical; i.e., a piece of a sphere or a plane.

Now if we assume that a point on a helicoidal surface is umbilic then all points belonging to the same orbit (helix) are umbilics and therefore that point is not isolated. So, when  $H$  is constant umbilic points cannot exist.  $\square$

**4. Proof of the results.** Looking at equation (4), we see that if  $\psi \not\equiv$  multiple of  $\pi/2$  then:

$$\psi \equiv \text{constant if and only if } H \equiv \text{constant}.$$

If  $\psi \equiv$  multiple of  $\pi/2$  then, by (5) we have that orbits are plane curves. This happens only if the surface is a surface of revolution or cylinder. This finishes the proof of the geometric characterization claimed.  $\square$

Next, let us consider a helicoidal surface  $M^2$  invariant under the motion  $\gamma$  with constant  $H$  and its deformation guaranteed by the Theorem of 0. Bonnet (section 1). Let  $N^2$  be a surface in this deformation, so that there is an isometry  $f : M^2 \rightarrow N^2$  which is onto, preserves  $H$  (and hence preserves the principal curvatures), and rotates the principal frame by a fixed angle. It is easy to see that

$$f \circ \gamma \circ f^{-1}(y, t) := f[\gamma(f^{-1}(y), t)], \quad y \in N^2, \quad t \in \mathbf{R},$$

forms a one-parameter group of isometries (one checks that (b) and (c) hold), under which  $N^2$  is invariant. Also these isometries preserve the principal curvatures and directions of  $N^2$ , hence they extend to rigid motions of  $E^3$  (and thus property (a) is valid as well). So,  $N^2$  is invariant under a one-parameter subgroup of rigid motions of  $E^3$ . The surface of revolution (Delaunay surface in this case) is obtained when, by rotating the principal frame,  $\psi$  becomes a multiple of  $\pi/2$ .  $\square$

*Remark 2.* The Delaunay surfaces of a given constant mean curvature form a one-parameter family of surfaces. Thus the helicoidal surfaces of a given constant mean curvature form a two-parameter family. The extra parameter is  $\psi$ .

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