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INTEGRABILITY OF TENSOR STRUCTURES OF ELECTROMAGNETIC TYPE

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Abstract. We study characterizations of the integrability of G-structures defined by tensor fields of elektromagnetic type.

1. Introduction. In [3] were considered the *G*-structures defined by a (1,1) tensor \tilde{J} on a differentiable manifold M^n such that

$$(\tilde{J}^2 - f^2)(\tilde{J}^2 + g^2) = 0,$$

where f, g are C^{∞} functions on M^n nowhere zero. This situation generalizes that of Hlavaty [4] and Mishra [7]. They consider the so called elektromagnetic tensor fields (of first class) on a 4-manifold which is the space-time of General Relativity. In [3] it was proved that the *G*-structure *P* defined by such a tensor field \tilde{J} is identical to the *G*-structure defined by a (1, 1) tensor field *J* that satisfies the same conditions as \tilde{J} but with f = g = 1, an so we have $J^4 = 1$.

On the other hand, that situation generalizes also the almost product and almost complex structures simultaneously. In [9] the family of linear connections that parallelize J (and an adapted metric also) is given. Connections partially adapted to such a structure are studied in [11]. In this note we study several characterizations and conditions of integrability of the (*G*-stucture defined by the) tenzor field \tilde{J} .

Thus, we consider the following situation:

Let M^n be a differentiable manifold and \tilde{J} a (1,1) tensor field such that:

a) $(\tilde{J}^2 - f^2)(\tilde{J}^2 + g^2) = 0$, where f, g are C^{∞} functions on M^n with f, g nowhere null;

b) The characteristic polynomial of \tilde{J} is $(x-f)^{r_1}(x+f)^{r_2}(x^2+g^2)^s$, where r_1, r_2, s are constants greater than or equal to 1 such that $r_1 + r_2 + 2s = n$. Then,

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since J which satisfies a) and b), but with f = g = 1, defines the same G-structure P as \tilde{J} (not an associated G-structure, but exactly the same P see [3]), we can characterize the integrability of P in terms of J.

2. Integrability in terms of the Nijenhuis tensor. We denote from now on by X, Y, \ldots , vectors fields on M^n . We consider the complementary projection operators $l = (J^2 + 1)/2$, $l_3 = (1 - J^2)/2$, which verify

$$Jl = lJ; J^2l = l, Jl_3 = l_3J, J^2l_3 = -l_3;$$

denote by L and L_3 the corresponding distributions, and put $L = L_1 \oplus L_2$, where L_1 and L_2 are distributions corresponding to the projectors l_1 and l_2 on L given by the eigenvalues +1 and -1 of $J \mid_L$. Let us decompose the Nijenhuis tensor of in the following manner:

$$N(X,Y) = lN(lX,lY) + l_3N(lX,lY) + N(lX,l_3Y) + N(l_3X,lY) + lN(l_3X,l_3Y) + l_3N(l_3X,l_3Y).$$

Then we have the following

PROPOSITION 2.1. a) L is integrable iff $(\forall X, Y)l_3N(lX, lY) = 0$;

b) L_3 is integrable iff $(\forall X, Y) lN(l_3X, l_3Y) = 0;$

c) If L is integrable, the almost product structure defined by $J \mid_L$ on each integral manifold of L is integrable iff $(\forall X, Y)N(lX, lY) = 0$;

d) If L_3 is integrable, the almost complex structure defined by $J \mid_{L_3}$ on each integral manifold of L_3 is integrable iff $(\forall X, Y)N(l_3X, l_3Y) = 0$.

Proof. a) $N(lX, lY) = [JlX, JlY] - J[JlX, lY] - J[lX, JlY] + J^2[lX, lY].$ Thus, if L is integrable, each bracket is an element of L and so $l_3N(lX, lY) = 0$. Conversely, suppose now that $l_3N(lX, lY) = 0$; then we obtain easily:

$$\begin{split} &l_3N(JlX, JlY) + Jl_3N(JlX, lY) + Jl_3N(lX, JlY) \\ &= 3l_3[lX, lY] + l_3(N(lX, lY) - J^2[lX, lY]) \\ &= 4l_3[lX, lY] + l_3N(lX, lY), \end{split}$$

and since by the hypothesis $l_3N(lX, lY) = 0$, L is integrable;

b) Analogous to a), if we consider now

$$lN(Jl_3X, Jl_3Y) + JlN(Jl_3X, l_3Y) + JlN(l_3X, Jl_3Y);$$

c) If L is integrable, then $J \mid_L$ induces on each integral manifold of L an almost product structure. As such a structure is integrable iff its Nijenhuis tensor is zero, that is, $N_{J|L}(lX, lY) = 0$, and since $N_{J|L}(lX, lY) = N(lX, lY)$, we obtain c);

d) Similar to c).

Definition 2.2. We say that J is partially integrable iff L and L_3 are integrable, and also the almost product and almost complex structure induced by J on the integral manifolds of L and L_3 , respectively.

Thus J is partially integrable iff $N(X, Y) = N(lX, l_3Y) + N(l_3X, lY)$.

So, we consider now the condition $N(lX, l_3Y) = 0$. Since the Lie derivative $L_Y^{\cdot}J$ verifies by definition $(L_Y^{\cdot}J)X = J[X,Y] - [JX,Y]$, we deduce:

a) $N(lX, l_3Y) = J(L_{l_3Y}J)lX - (L_{Jl_3Y}J)lX;$

b) $N(l_3X, lY) = J(L_{lY}J)l_3X - (L_{JlY}J)l_3X;$

and from these expressions it is immediate that:

PROPOSITION 2.3. $lN(lX, l_3Y) = 0$ (resp. $l_3N(lX, l_3Y) = 0$) for every X, Y iff $l(L_{l_3Z}J)l = 0$ (resp. $l_3(L_{lZ}J)l_3 = 0$ for every Z.

COROLLARY 2.4. $N(lX, l_3Y) = 0$ iff $l(L_{l_3Z}J)l = l_3(l_{lZ}J)l_3 = 0$, for every X, Y, Z.

Now, we have

THEOREM 2.5. J is integrable iff $N_J = 0$.

Proof. J is integrable iff for every $x \in M^n$, there exists a heighbourhood U of x and a coordinate system in $U, \{x^i\}$, such that the basis $\{\partial/\partial x^i\}, i = 1, ..., n$ is adapted in U. That is, J can be expressed as a linear combination of products $\partial/\partial x^i \otimes dx^j$ with constant coefficients, and so, trivially, N = 0.

Conversely, suppose N = 0. By a) and b) of Prop. 2.1, L and L_3 are integrable. Thus, for each $x \in M^n$ there exists a chart centered at $x, (U, \varphi)$, with coordinates $\{x^i, y^a\}, i = 1, \ldots, r_1 + r_2, a = 1, \ldots 2s$, such that

$$\partial/\partial x^i \in L, \ \partial/\partial y^a \in L_3.$$

So, in the local basis $\{\partial/\partial x^i, \partial/\partial y^a\}$, J has a matrix of the from

$$J = \begin{pmatrix} J_j^i & 0\\ 0 & J_b^a \end{pmatrix};$$

that is, $J = J_j^i \partial \partial x^i \otimes dx^j + J_b^a \partial \partial y^a \otimes dy^b$. Moreover, $l = \partial \partial x^i \otimes dx^i$, and $l_3 = \partial \partial y^a \otimes dy^a$. Thus,

$$L_{\partial/\partial y^c}J=rac{\partial J^i_j}{\partial y^c}rac{\partial}{\partial x^i}\otimes dx^i+rac{\partial J^a_b}{\partial y^c}rac{\partial}{\partial y^a}\otimes dy^b.$$

Hence

$$l(L_{\partial/\partial y^a}J)l = rac{\partial J^i_j}{\partial y^a}rac{\partial}{\partial x^i}\otimes dx^j.$$

So, Corollary 2.4 implies $\partial J_i^i / \partial y^a = 0$ (and analogously $\partial J_b^a / \partial x^i = 0$).

Consider now the integral manifold L_x of L, of coordinates $y^a = 0$. Since the almost product structure on L_x is integrable, there exist coordinate functions u^i in a neighbourhood of $x \in L_x$ in L_x such that

$$\partial/\partial u^i \in L_1, \ i = 1, \dots, r_1, \ \partial/\partial u^i \in L_2, \ i = r_1 + 1, \dots, r_1 + r_2$$

where these fields are considered in the regular submanifold $L_x \cap U$.

We define new coordinates in a neighbourhood W of x in $U\colon$ put, for $x'\in W,$

$$u^{i}(x') = u^{i}(\varphi^{-1}(x^{1}(x'), \dots, x^{r_{1}+r_{2}}(x')), 0, \dots, 0), \ \bar{y}^{a}(x') = y^{a}(x).$$

Then, for the new coordinate system $\{u^i, y^{-a}\}$ we have

$$\frac{\partial}{\partial x^i} = \frac{\partial u^j}{\partial x^i} \frac{\partial}{\partial u^j}, \ dx^i = \frac{\partial x^i}{\partial u^j} du^j, \ \text{ where } \ \frac{\partial}{\partial \bar{y}^a} \left(\frac{\partial u^i}{\partial x^i}\right) = \frac{\partial}{\partial \bar{y}^a} \left(\frac{\partial x^i}{\partial u^j}\right) = 0$$

This is a new coordinate system adapted to L and L_3 , and we have

$$J = \overline{J}^i_j \frac{\partial}{\partial u^i} \otimes du^j + \overline{J}^a_b \frac{\partial}{\partial \overline{y}^a} \otimes d\overline{y}^b, \text{ with } \frac{\partial \overline{J}^i_j}{\partial \overline{y}_a} = 0.$$

But for the points of coordinates $\bar{y}^a = 0$ we have by construction (see 2.1)

$$\overline{J}_{j}^{i} = \begin{pmatrix} I_{r_{1}} & 0\\ 0 & -I_{r_{2}} \end{pmatrix}.$$

Hence, in certain neighbourhood of x we also have the same matrix expression for \overline{J}_{I}^{i} .

Similarly, since the structure defined by J in L_{3x} is almost complex, a change of coordinates analogous to the previous one gives for \overline{J}_b^a the expression

$$\begin{pmatrix} 0 & -I_s \\ I_s & 0 \end{pmatrix}.$$

In other words, we have deduced that J is integrable.

3. Integrability in terms of a linear connection. Now, let ∇ be a linear connection without torsion on M^b and let Q be the (1, 2) tensor field on M^n defined by $Q(X, Y) = \{(\nabla_{JY}^{\cdot}J)X + J(\nabla_{Y}^{\cdot}J)X + 2J(\nabla_{X}^{\cdot}J)Y\}/4$.

We define a new connection D by means of the expression

 $D_X Y = l \nabla_X l X - l \nabla_{iY} l_3 X + l_3 \nabla_X l_3 Y - l_3 \nabla_{l_3 Y} l X + l Q(l X, l Y) - l_3 Q(l_3 X, l_3 Y).$

It is easily proved that:

$$i) D_{X}lY = l\nabla_{X}lY - l\nabla_{lX}l_{3}X + lQ(lX, lY);$$

$$ii) D_{X}l_{3}Y = l_{3}\nabla_{X}l_{3}Y - l_{3}\nabla_{l_{3}Y}lX - l_{3}Q(l_{3}X, l_{3}Y);$$

$$iii) D_{IX}lY = l\nabla_{IX}lY + lQ(lX, lY);$$

$$iv) D_{l_{3}X}lY = l\nabla_{l_{3}X}lY - l\nabla_{lY}l_{3}X;$$

$$v) D_{IX}l_{3}Y = l_{3}\nabla_{lX}l_{3}Y - l_{3}\nabla_{l_{3}Y}lX;$$

$$vi) D_{l_{3}X}l_{3}X = l_{3}\nabla_{l_{3}X}l_{3}Y - l_{3}Q(l_{3}X, l_{3}Y);$$

$$vii) D_{X}l = D_{X}l_{3} = 0.$$
(3.1)

So we have

PROPOSITION 3.1. The torsion T of D has the expression $T(X,Y) = 1/4\{lN(lY,lX) + l_3N(l_3N(l_3X,l_3Y))\} - l[l_3X,l_3Y] - l_3[lX.lY].$

Proof. Immediate from the expression for D and Q, applying that ∇ is torsionless, and proving that Q(X,Y) - Q(Y,X) = N(Y,X)/4

From that we obtain

COROLLARY 3.2. $(\forall X, Y)T(lX, l_3Y) = 0.$

We now prove

LEMMA 3.3. a) The distribution L is integrable iff $l_3T(lX, lY) = 0$;

b) The distribution L_3 is integrable iff $lT(l_3X, l_3Y) = 0$;

c) If L is integrable, then the almost product structure induced by $J \mid_L$ on each integral manifold of L is integrable iff lT(lX, lY) = 0;

d) If L_3 is integrable, then the almost complex structure induced by $J|_{L_3}$ on each integral manifold of L_3 is integrable iff $l_3T(l_3X, l_3Y) = 0$;

e) J is partially integrable iff T(X, Y) = 0;

f)
$$lN(JlX, lY) = lN(lX, JlY);$$

h) $l_3N(Jl_3X, l_3Y) = l_3N(l_3X, Jl_3Y);$
j) $(D_{l_3X}J)l_3Y = 0;$
Proof. a) If suffices to prove $l_3T(lX, lY) = -l_3[lX, lY];$
g) $lT(JX, lY) = lT(lX, JY);$
i) $l_3T(JX, l_3Y) = l_3T(l_3X, JY);$
k) $(D_{l_X}J)lY = 0.$
Proof. a) If suffices to prove $l_3T(lX, lY) = -l_3[lX, lY];$

b) analogous to a); c) it suffices to consider lT(lX, lY) = lN(lY, lX)/4 and a), c) of Prop. 2.1; d) it is deduced from $l_3T(l_3X, l_3Y) = l_3N(l_3X, l_3Y)/4$, and b), d) of Prop. 2.1.;

e) from Cor. 3.2 we obtain $T(X,Y) = T(lX,lY) + T(l_3X,l_3Y)$. If J is partially integrable, from a) and c) we deduce T(lX,lY) = 0 and from b) and d) that $T(l_3X,l_3Y) = 0$. Hence T(X,Y) = 0. Conversely, T(X,Y) = 0 implies $T(lX,lY) = T(l_3X,l_3Y) = 0$ and thus $lT(lX,lY) = l_3T(lX,lY) = lT(l_3X,l_3Y) = l_3T(l_3X,l_3Y) = 0$.

From these equalities and from a), b), c), d) we deduce that J is partially integrable; f) the proof is immediate and moreover, as a consequence, we obtain lN(JlX, JlY) = lN(lX, lY); g) lT(JX, lY) = lN(lY, JlX)/4 and lT(JX, lX) =lN(JlY, lX)/4, and from f) we obtain the result; h) the proof is analogous to that of f) and we deduce here $l_3N(Jl_3X, Jl_3Y) = -l_3N(l_3X, l_3Y)$;

i) we have $l_3T(JX, l_3Y) = l_3N(Jl_3X, l_3Y)/4$ and $l_3T(l_3X, JY) = l_3N(l_3X, Jl_3Y)/4$, and the conclusion follows from h);

j) from (3.1), vi) we have $D_{l_3X}Jl_3Y = l_3\nabla_{l_3X}Jl_3Y - l_3Q(l_3X, Jl_3Y)$ and $JD_{l_3X}l_3Y = Jl_3\nabla_{l_3X}l_3Y - l_3JQ(l_3X, l_3Y)$. Substracting we obtain

$$\begin{split} (D_{l_3X}J)l_3Y &= l_3(\nabla_{l_3X}J)l_3Y + l_3\{J(\nabla_{Jl_3Y}J)l_3X + J^2(\nabla_{l_3X}J)l_3X \\ &+ 2J^2(\nabla_{l_3X}J)l_3Y - (\nabla_{J^2l_3Y}J)l_3X - J(\nabla_{Jl_3X}J)l_3X \\ &- 2J(\nabla_{l_3X})Jl_3Y\}/4 = l_3(\nabla_{l_3X}J)l_3Y - l_3(\nabla_{l_3X}J)l_3Y = 0; \end{split}$$

k) analogous to that of j), by using (3.1), iii).

THEOREM 3.4. J is integrable iff there exists a linear connection without torsion that parallelzes J. If J is integrable, then D gives an explicit example of such a connection.

Proof. Suppose J integrable. Then for the earlier connection D we have: 1) D is torsionless, and 2) DJ = 0. Indeed, if J is integrable then it is partially integrable and, from e) of Lemma 3.3. we obtain 1). On the other hand, if J is integrable, N(X, Y) = 0, and since from (3.1), v) we have

$$(D_{JlX}J)l_3Y + (D_{lX}J)Jl_3Y = l_3N(lX, l_3Y) = 0,$$

we deduce

$$(D_{JlX}^{\cdot}J)l_{3}Y = -(D_{lX}J)Jl_{3}Y.$$
(3.2)

Substituting Y by JY we have

$$(D_{JlX}J)l_3JY = -(D_{lX}J)J^2l_3Y = (D_{lX}J)l_3Y,$$
(3.3)

and, if in (3.2) we substitute X by JX, we have

$$(D_{JlX}^{\cdot}J)Jl_{3}Y = -(D_{J^{2}lX}J)l_{3}Y = -(D_{l}^{\cdot}J)l_{3}Y.$$
(3.4)

From (3.3) and (3.4) we deduce

$$(D_{lX}J)l_{3}Y = 0. (3.5)$$

From (3.1), iv) we have $lN(l_3X, lY) = (D_{Jl_3X}J)lY + (D_{l_3X}J)JlY$; if J is integrable we obtain analogously

$$(D_{l_3X}J)lY = 0. (3.6)$$

But in j) and k) of Lemma 3.3. we have

$$(D_{l_3X}J)l_3Y = 0$$
 and $(D_{lX}J)lY = 0$ (3.7)

Hence, from (3.5) (3.6) and (3.7) follows 2).

Conversely, suppose now that there is a linear connection ∇ without torsion such that $\nabla J = 0$. If we consider Q from ∇ as before, we see that Q = 0. But in the proof of the Prop. 3.1 we have seen that Q(X,Y) - Q(Y,X) = N(Y,X)/4; that is, J is integrable.

Remark. As is well known, Lehmann-Lejeune [5] proves that, for 0-deformable (1,1) tensor fields, the integrability is equivalent to the existence of a torsionless structural *local* connection. In our case, we have a global connection and we also give its explicit expression when J is integrable.

4. Integrability in terms of the structure tensor. We have now at disposal two criteria of integrability of the G-structure P defined by \tilde{J} . The first

one in terms of the Nijenhuis tensor of the field J, the second one in terms of a linear connection. A third criterion is that which expresses the integrability in terms of the Guillemin stucture tensors [2].

The field \tilde{J} is not 0-deformable, but the associted field J, which defines the same structure P, is 0-deformable and so, we can anew characterize the integrability of P in terms of J; but from the results of Lehmann-Lejeune [5] it suffices to consider, in this case, the Chern-Ehresmann-Bernard tensor [1] and have the equivalence of the integrability with nullity of the 1-st structure tensor of P, as we express in the final theorem.

5. Integrability in terms of prolongations and complete lifts. Now, we consider of the one hand the complete lift J^c of J in the sense of Yano-Kobayashi [12], whic is a (1,1) tensor field on TM^n defined from J and, on the other hand, the canonical prolongation \hat{J} of J in the sense of Morimoto [8], which is also a (1,1) tensor field on TM^n . Firstly, we have the following:

PROPOSITION 5.1. The canonical prolongation to TN^n of the $(J^4 = 1)$ -structure J is a $(J^4 = 1)$ -structure \hat{J} which coincides with the complete lift J^c of J.

Proof. The structural group G corresponding to J (and \tilde{J}) is that of matrices of the form [3]

$$\begin{bmatrix} A & 0 \\ 0 & B \\ \hline 0 & C & -D \\ D & C \end{bmatrix}$$

where $A \in Gl(r_1, \mathbf{R}), B \in Gl(r_2, \mathbf{R}), C + iD \in Gl(s, \mathbf{C}), r_1 + r_2 + 2s = n$. If we denote \hat{G} the structural group of the canonical prolongation \hat{P} of P defined by J (and \tilde{J}), it has as elements the \hat{g} obtained by means of

$$\hat{g} = j_n(\{g, X\}), \quad g \in G, \quad X \in \mathbf{g},$$

where **g** denotes the Lie algebra of G, X the translated $R_{g^{-1}*}Y$, for a certain $Y \in T_g G$, and j_n the imbedding $j_n : TGl(n, \mathbf{R}) \to Gl(2n, \mathbf{R})$.

More precisely,

$$j_{n}(\{g, X\}) = \begin{bmatrix} g & 0 \\ Xg & g \end{bmatrix} = \begin{bmatrix} A & 0 \\ 0 & B \\ 0 & C \\ 0 & D & C \\ \hline \alpha & 0 & 0 \\ 0 & \beta & 0 \\ \hline 0 & \beta & 0 \\ \hline 0 & \beta & \gamma & 0 \\ \hline 0 & B & 0 \\ \hline 0 & C & -D \\ 0 & C \end{bmatrix}$$
(5.1)

since

$$Xg = \begin{bmatrix} M & 0 \\ 0 & N \\ 0 \\ 0 \\ 0 \\ 0 \\ Q \\ P \end{bmatrix} \begin{bmatrix} A & 0 \\ 0 & B \\ 0 \\ 0 \\ D \\ C \end{bmatrix} = \begin{bmatrix} MA & 0 \\ 0 \\ NB \\ 0 \\ 0 \\ PC - QD - PD - QC \\ 0 \\ QC + PD - QD + PC \end{bmatrix}$$

where $M \in \mathcal{M}(r_1; \mathbf{R}), N \in \mathcal{M}(r_2; \mathbf{R}), P + iQ \in \mathcal{M}(s, \mathbf{C}).$

It is immediate that $j_n(\{g, X\})$, belongs to the matrix group $Gl(2r_1, \mathbf{R}) \times Gl(2r_2, \mathbf{R}) \times Gl(2s, \mathbf{C})$, by means of a convenient rearrangement of the boxes of the matrix (5.1). Hence, we have the structural group od a $(J^4 = 1)$ -structure on TM^n .

On the other hand, for a given local coordinate system $\{U, x^1, \ldots, x^n\}$ on M^n , and a section σ of the principal bundle of frames FM^n on U, expressed as

$$\sigma(x) = \left(\dots, \sum_{i=1}^{n} \sigma_{j}^{i}(x) \frac{\partial}{\partial x^{i}} \Big|_{x}, \dots\right), \ x \in U,$$

Morimoto [8] proves that $\tilde{\sigma} = j_{M^n} \cdot T\sigma$ (where j_{M^n} is the canonical embedding $j_{M^n} : TFM^n \to FTM^n$), is a section of FTM^n on TU, which can be expressed as

$$\tilde{\sigma}\left(\sum_{i=1}^{n} y^{i} \frac{\partial}{\partial x^{i}} \Big|_{x}\right) = \left(\dots, \sum_{i=1}^{n} \sigma_{j}^{i}(x) \frac{\partial}{\partial x^{i}} \Big|_{x} + \sum_{i,k=1}^{n} \frac{\partial \sigma^{i}(x)}{\partial x^{k}} y^{k} \frac{\partial}{\partial y^{i}} \Big|_{X} \dots, \sum_{i=1}^{n} \sigma_{j}^{i}(x) \frac{\partial}{\partial y^{i}} \Big|_{X}, \dots\right)$$

where $\{x^1, \ldots, x^n, y^1, \ldots, y^n\}$ is the local coordinate system induced in TU, and $X = \sum_{i=1}^n y^i \frac{\partial}{\partial x^i} \Big|_x \in TU.$

We now consider the diagram



with the explained notations. Let again $\{U, x^1, \ldots, x^n\}$ a local coordinate system on M^n , and σ a section of the *G*-structure *P* on *U*; then $\tilde{\sigma} = j_{M^n} \cdot T\sigma$ is a section of the canonical prolongation \hat{P} of *P*, since

$$\hat{\sigma}(TU) = j_{M^n} \cdot T\sigma(TU) \subset j_{M^n}(T(\sigma(U)) \subset j_{M^n}(TP) = \hat{P}.$$

Now, let $J_0: \mathbf{R}^n \to \mathbf{R}^n$ be an automorphism such that $J_0^4 = 1$. From the diagram

$$\begin{array}{ccc} T_x M^n & \xrightarrow{J_x} & T_x M^n \\ \sigma(x) & \uparrow & \uparrow \sigma(x) \\ \mathbf{R}^n & \xrightarrow{J_0} & \mathbf{R}^n \end{array}$$

we define J_x as $J_x = \sigma(x) \cdot J_0 \cdot \sigma(x)^{-1}$.

Then $J: x \sim \to J_x$ is the $(J^4 = 1)$ -structure associated to P globally defined. Indeed:

a) $J^4 = 1$. Immediate from $J_0^4 = 1$;

b) globaly defined: If $x \in U \cap U'$, where U, U' are coordinate neghbourhoods and σ' is a section of P on U', then $J'_x = \sigma'(x) \cdot J_0 \cdot \sigma'(x)^{-1}$, but since

$$\sigma'(x) = g(x) \cdot \sigma(x), \quad g(x) \in G.$$

we deduce

$$J'_x = g(x) \cdot \sigma(x) \cdot J_0 \cdot \sigma(x)^{-1} \cdot g(x)^{-1} = g(x) \cdot J_x \cdot g(x)^{-1}$$

We now consider $TJ_0: T\mathbf{R}^n \to T\mathbf{R}^n$. Since

$$TJ_0 = \begin{bmatrix} J_0 & 0\\ 0 & J_0 \end{bmatrix}$$

we have $(TJ_0)^4 = 1$, and we define

$$\hat{J}(X) = \hat{\sigma}(X) \cdot TJ_0 \cdot \hat{\sigma}(X)^{-1}$$
, for every $X \in TU$.

It is clear that $\hat{J}^4 = 1$, and \hat{J} is the $(J^4 = 1)$ -structure on TM^n canonical prolongation of J, since $\hat{\sigma} = j_{M^n} \cdot T\sigma$.

On the other hand, we can choose as a basis of $T_X T M^n$ the set

$$\bigg\{\frac{\partial}{\partial x^1}\Big|_X, \dots, \frac{\partial}{\partial x^n}\Big|_X, \frac{\partial}{\partial y^1}\Big|_X, \dots, \frac{\partial}{\partial y^n}\Big|_X\bigg\}.$$

Then, using the earlier expressions for $\hat{\sigma}$ and TJ_0 , we obtain

$$\hat{J}(x) = \hat{\sigma}(x) \cdot TJ_0 \cdot \hat{\sigma}(x)^{-1} = \begin{bmatrix} \sigma(x) & 0\\ \partial \sigma(x) & \sigma(x) \end{bmatrix} \begin{bmatrix} J_0 & 0\\ 0 & J_0 \end{bmatrix} \begin{bmatrix} \sigma(x) & 0\\ \partial \sigma(x) & \sigma(x) \end{bmatrix}^{-1} = \begin{bmatrix} J_x & 0\\ \partial J_x & J_x \end{bmatrix},$$

which is precisely the formula of the complete lift J^c of J (see [12]), being

$$\partial J_x = \left(\sum_{k=1}^n y^k \frac{\partial J_j^i}{\partial x^k}\right)$$

But Morimoto [8] proves that a G-structure P is integrable if and only if the canonical prolongation \hat{P} is integrable; hence we obtain.

PROPOSITION 5.2. Let (M^n, J) be a $(J^4 = 1)$ -manifold. Then the following statements are equivalent:

a) The G-structure P defined by J is integrable;

b) The Nijenhuis tensor of the tensor \hat{J} corresponding to the canonical prolongation \hat{P} of P is zero;

c) The Nijenhuits tensor of the complete lift J^c of J is zero.

6. J-Lie groups. Now, we consider a sufficient condition in order to a $(j^4 = 1)$ -structure be integrable.

Let (M_1, J_1) and (M_2, J_2) be two $(J^4 = 1)$ -manifolds. We say that a differentiable map $f : M_1 \to M_2$ is a *J*-map if and anly if the following diagram is commutative

$$\begin{array}{cccc} T_x M_1 & \xrightarrow{J_*} & T_{f(x)} M_2 \\ \\ J_1 \downarrow & & \downarrow J_2 & \text{for every } x \in M_1. \\ \\ T_x M_1 & \xrightarrow{f_*} & T_{f(x)} M_2 \end{array}$$

Definition 6.1. We call J-Lie group a Lie Group G with a $(J^4 = 1)$ -structure J such that the usual translations L_g and R_g are J-maps, for every $g \in G$.

Thus, we have

PROPOSITION 6.2. If G is a J-Lie group, then J is integrable.

Proof. Since L_g and R_g are J-maps, we have ad $g \cdot J = J \cdot ad g$.

In particular for $g = \exp tX$, $X \in \mathbf{g}$, $t \in \mathbf{R}$, we have

$$\exp(\operatorname{Ad} tX) \cdot JY = J(\exp(\operatorname{Ad} tX)Y), \text{ for every } Y \in g.$$
(6.1)

Moreover, we obtain

$$\exp(\operatorname{Ad} tX)JY = JY + t[X, JY] + t^{2}[X, [X, JY]]/2 + \cdots$$
$$J \exp(\operatorname{Ad} tX)Y = JY + tJ[X, Y] + t^{2}J[X, [X, Y]]/2 + \cdots$$

Hence, from (6.1) and taking the limit for $t \to 0$ we deduce [X, JY] = J[X, Y], and also J[X, Y] = -J[Y, X] = -[Y, JX] = [JX, Y].

Thus, it is immediate $N(X, Y) = 0, X, Y \in \mathbf{g}$.

7. Characterizations of the integrability. Finally, according to the earlier results we can give the following:

THEOREM 7.1. Let M^n be a differentiable manifold with a tensor field of electromagnetic type and class \tilde{J} . Then the following statements are equivalent:

a) The G-structure P defined by \tilde{J} is integrable;

b) The Nijenhuis tensor of the associated tensor field J is zero;

c) There exists a linear torsionless connection which parallelizes J;

d) The structure tensor of P is zero;

e) The Nijenhuis tensor of the tensor field \hat{J} corresponding to the canonical prolongation \hat{P} of P is zero;

f) The Nijenhuis tensor of the complete lift J^c of J is zero; moreover,

g) If G is a J-Lie group, then J is integrable.

When M^n is J-Kaehlerian [10], other conditions can be given.

We note that any linear connection which parallelizes \tilde{J} does not exist.

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