

THE LAPLACE TRANSFORM OF THE MODIFIED
BESSEL FUNCTION OF THE SECOND KIND*

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Dedicated to the memory of Professor Arthur Erdélyi

1. Denoting the Mellin transform of $f \in L(0, \infty)$ by

$$(1) \quad \mathfrak{M}\{f(x):s\} = \int_0^{\infty} x^{s-1} f(x) dx,$$

we know that [2, Vol. I, p. 331, Eq. (26)]

$$(2) \quad \mathfrak{M}\{K_\nu(\alpha x):s\} = 2s^{-2} \alpha^{-s} \Gamma\left(\frac{1}{2}s - \frac{1}{2}\nu\right) \Gamma\left(\frac{1}{2}s + \frac{1}{2}\nu\right),$$

provided that $\operatorname{Re}(\alpha) > 0$ and $\operatorname{Re}(s) > |\operatorname{Re}(\nu)|$.

If in the known result (2) we set $x = t^\rho$, $dx = \rho t^{\rho-1} dt$, ρ being a real number, and (for simplicity) replace s by s/ρ , $\rho \neq 0$, we shall at once have

$$(3) \quad \int_0^{\infty} t^{s-1} K_\nu(\alpha t^\rho) dt = \frac{2^{(s-2\rho)/\rho}}{|\rho|^{s/\rho}} \Gamma\left(\frac{s}{2\rho} - \frac{1}{2}\nu\right) \Gamma\left(\frac{s}{2\rho} + \frac{1}{2}\nu\right),$$

where, for convergence, $\operatorname{Re}(\alpha) > 0$, and either

$$(4) \quad \rho > 0, \operatorname{Re}(s/\rho) > |\operatorname{Re}(\nu)|$$

or

$$(5) \quad \rho < 0, \operatorname{Re}(\nu) < \operatorname{Re}(s/\rho) < \frac{1}{2}.$$

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In order to compute the Laplace transform of the modified Bessel function $K_\nu(\alpha t)$, $\rho > 0$, we put

$$(6) \quad I_\rho^\mu(s) \equiv \int_0^\infty e^{-st} t^{\mu-1} K_\nu(\alpha t^\rho) dt,$$

expand the exponential function e^{-st} in powers of t , and integrate term-by-term using (3). We thus find that

$$(7) \quad I_\rho^\mu(s) = \frac{2^{(\mu-2\rho)/\rho}}{\rho \alpha^{\mu/\rho}} \sum_{n=0}^{\infty} \frac{(-s)^n}{n!} \left(\frac{2}{\alpha}\right)^{n/\rho} \Gamma\left(\frac{\mu+n}{2\rho} - \frac{1}{2}\nu\right) \Gamma\left(\frac{\mu+n}{2\rho} + \frac{1}{2}\nu\right),$$

which holds true under the following sufficient) conditions:

$$(8) \quad \rho > 0, \operatorname{Re}(s) > 0, \operatorname{Re}(\alpha) > 0, \operatorname{Re}(\mu/\rho) > |\operatorname{Re}(\nu)|,$$

it being assumed that the series on the right makes sense.

Obviously, owing to the *second* inequality in (5), the integral (6) cannot be evaluated in this manner using (3) when $\rho < 0$. To overcome this difficulty, however, we can appeal to the Eulerian integral for the Γ -function in conjunction with the (Mellin-Barnes) contour integral representation

$$(9) \quad K_\nu(x) = \frac{1}{2\pi i} \int_{\sigma-i\infty}^{\sigma+i\infty} 2^{s-2} \Gamma\left(\frac{1}{2}s - \frac{1}{2}\nu\right) \Gamma\left(\frac{1}{2}s + \frac{1}{2}\nu\right) x^{-s} ds,$$

$$x > 0, \sigma > |\operatorname{Re}(\nu)|,$$

which indeed follows rather immediately from (2) and the Mellin inversion theorem. We thus obtain

$$(10) \quad I_\rho^\mu(s) = \frac{s^{-\mu}}{4} H_{0,3}^{3,0} \left[\frac{\alpha}{2s^\rho} \left| (\mu, -\rho), \left(\frac{1}{2}\nu, \frac{1}{2}\right), \left(-\frac{1}{2}\nu, \frac{1}{2}\right) \right. \right],$$

where $\rho < 0$, $H_{p,q}^{m,n}[z|\dots]$ denotes the familiar H -function of C. Fox ([3]), p. 408; see also [5], p. 214, Eq. (2.4) *et seq.*, and for convergence,

$$(11) \quad \min \left\{ \operatorname{Re}(s), \operatorname{Re}(\alpha), \operatorname{Re}\left(\mu - \frac{1}{2}\rho\right) \right\} > 0.$$

Evidently, for $\mu=1$, our formulas (7) and (10) would provide the Laplace transforms of $K_\nu(\alpha t^\rho)$ for positive or negative values of ρ .

2. Special cases of $I_\nu^\mu(s)$

There are a number of interesting special situations in which the right-hand sides of our formulas (7) and (10) can be written in relatively more familiar forms. We, therefore, find it worthwhile to record the following special cases:

Case 1. When $\rho = m$, where m is a positive integer, we can make use of the elementary identity

$$(12) \quad \sum_{n=0}^{\infty} \Phi(n) = \sum_{j=0}^{N-1} \sum_{n=0}^{\infty} \Phi(nN+j), \quad N=1, 2, 3, \dots,$$

which is easy to verify, and (7) would thus yield

$$(13) \quad I_m^\mu(s) = \frac{2^{(\mu-2m)/m}}{m \alpha^{\mu/m}} \sum_{j=0}^{2m-1} \sum_{n=0}^{\infty} \frac{(-s)^{2mn+j}}{(2mn+j)!} \left(\frac{2}{\alpha}\right)^{2n+j/m} \cdot \Gamma\left(\frac{\mu+j}{2m} - \frac{1}{2} \nu + n\right) \Gamma\left(\frac{\mu+j}{2m} + \frac{1}{2} \nu + n\right),$$

provided that the conditions in (8) hold true with $\rho = m \in \{1, 2, 3, \dots\}$.

By Gauss's multiplication theorem (cf. [1], Vol. I, p. 4, Eq. (11)), we have

$$(14) \quad (2mn+j)! = j! (2m)^{2mn} \prod_{k=1}^{2m} \left(\frac{j+k}{2m}\right)_n,$$

where, for convenience, $(\lambda)_n = \Gamma(\lambda+n)/\Gamma(\lambda)$, and (13) will readily reduce to the elegant form:

$$(15) \quad I_m^\mu(s) = \frac{2^{(\mu-2m)/m}}{m \alpha^{\mu/m}} \sum_{j=0}^{2m-1} \frac{(-s)^j}{j!} \left(\frac{2}{\alpha}\right)^{j/m} \Gamma\left(\frac{\mu+j}{2m} - \frac{1}{2} \nu\right) \Gamma\left(\frac{\mu+j}{2m} + \frac{1}{2} \nu\right) \cdot {}_2F_{2m-1} \left[\begin{matrix} (\mu-m\nu+j)/2m, (\mu+m\nu+j)/2m; \\ (j+1)/2m, \dots * \dots, (j-2m)/2m; \end{matrix} \frac{4}{\alpha^2} \left(\frac{s}{2m}\right)^{2m} \right],$$

where m is a positive integer, $\text{Re}(s) > 0$, $\text{Re}(\alpha) > 0$, $\text{Re}(\mu) > m|\text{Re}(\nu)|$, and the asterisk in the generalized hypergeometric functions represents the fact that the denominator parameter $(2m)/2m$ is always omitted.

Remark 1. For $m=1$, the conditions $\text{Re}(s) > 0$ and $\text{Re}(\alpha) > 0$ can obviously be replaced by a single condition: $\text{Re}(s+\alpha) > 0$. Indeed, in this case our formula (15) would simplify considerably, and we first have

$$(16) \quad I_1^\mu(s) = \frac{2^{\mu-2}}{\alpha^\mu} \left\{ \Gamma\left(\frac{1}{2}\mu - \frac{1}{2}\nu\right) \Gamma\left(\frac{1}{2}\mu + \frac{1}{2}\nu\right) {}_2F_1 \left[\begin{matrix} \frac{1}{2}(\mu-\nu), \frac{1}{2}(\mu+\nu); \\ \frac{1}{2}; \end{matrix} \frac{s^2}{\alpha^2} \right] \right. \\ \left. - \frac{2s}{\alpha} \Gamma\left(\frac{1}{2}\mu - \frac{1}{2}\nu + \frac{1}{2}\right) \Gamma\left(\frac{1}{2}\mu + \frac{1}{2}\nu + \frac{1}{2}\right) {}_2F_1 \left[\begin{matrix} \frac{1}{2}(\mu-\nu+1), \frac{1}{2}(\mu+\nu+1); \\ \frac{3}{2}; \end{matrix} \frac{s^2}{\alpha^2} \right] \right\},$$

provided that $\operatorname{Re}(s+\alpha) > 0$ and $\operatorname{Re}(\mu) > |\operatorname{Re}(\nu)|$.

Making use of Gauss's quadratic transformations (*cf.*, *e.g.*, [1], Vol. I, p. 111, Eq. (7) and Eq. (11)), the familiar (Euler's) formula [*op. cit.*, p. 64, (22)], and Legendre's duplication theorem [*op. cit.*, p. 5, Eq. (15)], we shall thus arrive at the well-known result [1, Vol. II, p. 50, Eq. (26)]:

$$(17) \quad I_1^\mu(s) = \frac{(2\alpha)^\nu \sqrt{\pi} \Gamma(\mu-\nu) \Gamma(\mu+\nu)}{(s+\alpha)^{\mu+\nu} \Gamma\left(\mu + \frac{1}{2}\right)} \left[\begin{matrix} \mu+\nu, \nu + \frac{1}{2}; \\ \mu + \frac{1}{2}; \end{matrix} \frac{s-\alpha}{s+\alpha} \right],$$

where, for convergence, $\operatorname{Re}(s+\alpha) > 0$ and $\operatorname{Re}(\mu) > |\operatorname{Re}(\nu)|$.

Case 2. When $\rho = -m$, where m is a positive integer, we can use Gauss's multiplication theorem once again to get the following special case of our formula (11):

$$(18) \quad I_{-m}^\mu(s) = \frac{2^{\mu-m-1} m^{\mu-\frac{1}{2}}}{s^\mu \pi^{m-\frac{1}{2}}} G_{0,2m+2}^{2m+2,0} \left(\frac{\alpha^2}{4} \left(\frac{s}{2m} \right)^{2m} \middle| \Delta(2m; \mu), \frac{1}{2}\nu, -\frac{1}{2}\nu \right),$$

where, as before, m is a positive integer, $\operatorname{Re}(s) > 0$, $\operatorname{Re}(\alpha) > 0$, $\operatorname{Re}(\mu) > -\frac{1}{2}m$, $G_{p,q}^{m,n}(z|\dots)$ denotes the familiar (Meijer's) G -function (*cf.*, *e.g.*, [1], Vol. I, p. 207 *et seq.*), and for convenience, $\Delta(m; \lambda)$ denotes the set of m parameters $\lambda/m, (\lambda+1)/m, \dots, (\lambda+m-1)/m, m \geq 1$.

Remark 2. For $\mu=1$, our formulas (15) and (18) would obviously yield the Laplace transforms of the modified Bessel functions $K_\nu(\alpha t^{\pm m})$, where $m=1, 2, 3, \dots$; in fact, these consequences of our general results (7) and (11), and their *trivial* cases when $s=1$, can be shown to lead us fairly easily to the

corrected (or modified) versions of the known results [4, p. 325, Eq. (1) and Eq. (2); p. 326, Eq. (3) and Eq. (4)], which were derived erroneously from certain integrals involving MacRobert's E -function (cf. [1], Vol. I, p. 203, § 5.2).

Case 3. Yet another interesting special case of our formula (7) would occur when $\rho = \frac{1}{2}$. Indeed, in terms of the Whittaker functions [1, Vol. I, p. 264, Eq. (5)], we readily obtain the elegant formula (cf. [2], Vol. I, p. 199, Eq. (37)):

$$(19) \quad I_{\frac{1}{2}}^{\mu}(s) = \frac{\Gamma\left(\mu - \frac{1}{2}\nu\right) \Gamma\left(\mu + \frac{1}{2}\nu\right)}{\alpha s^{\mu - \frac{1}{2}}} e^{\alpha^2/8s} W_{\frac{1}{2} - \mu, \frac{1}{2}\nu}\left(\frac{\alpha^2}{4s}\right),$$

where, for convergence, $\text{Re}(s) > 0$, $\text{Re}(\alpha) > 0$, and $\text{Re}(\mu) > \frac{1}{2}|\text{Re}(\nu)|$.

Since [1, Vol. I, p. 265, Eq. (14)]

$$(20) \quad W_{0,\nu}(z) = \sqrt{(z/\pi)} K_{\nu}\left(\frac{1}{2}z\right),$$

this last formula (19), with $\mu = \frac{1}{2}$, yields another known result (cf. [2], Vol. I, p. 199, Eq. (35)):

$$(21) \quad I_{\frac{1}{2}}^{\frac{1}{2}}(s) = \frac{\sec\left(\frac{1}{2}\nu\pi\right)}{2\sqrt{(s/\pi)}} e^{\alpha^2/8s} K_{\frac{1}{2}\nu}\left(\frac{\alpha^2}{8s}\right),$$

provided that $\min\{\text{Re}(s), \text{Re}(\alpha)\} > 0$ and $|\text{Re}(\nu)| < 1$.

Case 4. For $\rho = -\frac{1}{2}$, our formula (10) evidently gives us the result

$$(22) \quad I_{-\frac{1}{2}}^{\mu}(s) = \frac{s^{-\mu}}{2} G_{0,3}^{3,0}\left(\frac{\alpha^2 s}{4} \left| \begin{matrix} - \\ \mu, \frac{1}{2}\nu, -\frac{1}{2}\nu \end{matrix} \right. \right),$$

$$\min\left\{\text{Re}(s), \text{Re}(\alpha), \text{Re}\left(\mu + \frac{1}{4}\right)\right\} > 0,$$

in terms of the G -function occurring in (18).

Case 5. Our formula (7) simplifies also in the special case $\rho = m + \frac{1}{2}$, where m is a positive integer, the case $m=0$ being already covered by case 3. Applying the identities (12) and (14) appropriately, we thus find that

$$(23) \quad I_{m+\frac{1}{2}}^{\mu}(s) = \frac{4^{(\mu-2m-1)/(2m+1)}}{\left(m+\frac{1}{2}\right)\alpha^{(2\mu)/(2m+1)}} \sum_{j=0}^{2m} \frac{(-s)^j}{j!} \left(\frac{2}{\alpha}\right)^{2j/(2m+1)} \\ \cdot \Gamma\left(\frac{\mu+j}{2m+1} - \frac{1}{2}\nu\right) \Gamma\left(\frac{\mu+j}{2m+1} + \frac{1}{2}\nu\right) \\ \cdot {}_2F_{2m} \left[\begin{matrix} (j+1)/(2m+1) - \frac{1}{2}\nu, (\mu+j)/(2m+1) + \frac{1}{2}\nu; \\ (j+1)/(2m+1), \dots, (j+2m+1)/(2m+1); \end{matrix} \frac{4}{\alpha^2} \left(\frac{s}{2m+1}\right)^{2m+1} \right],$$

where m is a positive integer, $\operatorname{Re}(s) > 0$, $\operatorname{Re}(\alpha) > 0$, $\operatorname{Re}(\mu) > \left(m + \frac{1}{2}\right) |\operatorname{Re}(\nu)|$, and, as before, the asterisk in the generalized hypergeometric functions represents the fact that the denominator parameter $(2m+1)/(2m+1)$ is always omitted.

Case 6. For $\rho = -m - \frac{1}{2}$, where m is a positive integer, our formula (10) would reduce to the form:

$$(24) \quad I_{-m-\frac{1}{2}}^{\mu}(s) = \frac{(2m+1)^{\frac{\mu-1}{2}} s^{-\mu}}{2^{m+1} \pi^m} \\ \cdot G_{0,2m+3}^{2m+3,0} \left(\frac{\alpha^2}{4} \left(\frac{s}{2m+1}\right)^{2m+1} \left| \frac{1}{\Delta(2m+1; \mu, \frac{1}{2}\nu, -\frac{1}{2}\nu)} \right. \right),$$

$$\min \left\{ \operatorname{Re}(s), \operatorname{Re}(\alpha), \operatorname{Re} \left(2\mu + m + \frac{1}{2} \right) \right\} > 0,$$

which evidently would also incorporate Case 4 if we let $m=0$. {See also Equation (18) above.}

We remark in passing that such simplifications of our general results (7) and (10) as in Cases 5 and 6 above can be made whenever ρ is a rational number.

3. Further generalizations

In view of (9), if we define

$$(25) \quad \Xi(x) = \frac{1}{2\pi i} \int_{\sigma-i\infty}^{\sigma+i\infty} \frac{\prod_{j=1}^q \Gamma(\beta_j + B_j s)}{\prod_{j=1}^r \Gamma(\gamma_j - C_j s) \prod_{j=1}^p \Gamma(\alpha_j + A_j s)} x^{-s} ds,$$

where $x > 0$, the A 's, B 's, and C 's are positive numbers such that

$$(26) \quad \Omega \equiv \sum_{j=1}^p A_j - \sum_{j=1}^q B_j - \sum_{j=1}^r C_j \leq 0,$$

in which the equality holds provided that

$$(27) \quad 0 < x < R \equiv \prod_{j=1}^p (A_j)^{A_j} \prod_{j=1}^q (B_j)^{-B_j} \prod_{j=1}^r (C_j)^{-C_j},$$

and σ is any real number satisfying the inequality

$$(28) \quad \sigma > \omega \equiv - \min_{1 \leq j \leq q} \left\{ \operatorname{Re} \left(\frac{\beta_j}{B_j} \right) \right\},$$

then it is easily verified that $\Xi(x)$ can be expressed as the sum of q Wright's generalized hypergeometric functions (cf. [6]). Furthermore, in terms of the familiar (Fox's) H -function [3, p. 408], we have the following relationship (see also [5], p. 214, Eq. (2.4) *et seq.*)

$$(29) \quad \Xi(z) = H_{p, q+r}^{q, 0} \left[z \left| \begin{matrix} (\alpha_1, A_1), \dots, (\alpha_p, A_p) \\ (\beta_1, B_1), \dots, (\beta_q, B_q), (1 - \gamma_1, C_1), \dots, (1 - \gamma_r, C_r) \end{matrix} \right. \right],$$

provided, for example, that

$$(30) \quad |\arg(z)| < \frac{1}{2} \Lambda \pi, \quad \Lambda \equiv \sum_{j=1}^q B_j - \sum_{j=1}^r C_j - \sum_{j=1}^p A_j > 0.$$

Equation (25), in conjunction with (1) and the Mellin inversion theorem, can be shown to imply

$$(31) \quad \mathfrak{M} \{ \Xi(x) : s \} = \frac{\prod_{j=1}^q \Gamma(\beta_j + B_j s)}{\prod_{j=1}^r \Gamma(\gamma_j - C_j s) \prod_{j=1}^p \Gamma(\alpha_j + A_j s)},$$

where, for convergence, $\operatorname{Re}(s) > \omega$, ω being given by (28).

Evidently, in view of the relationship (29), this last formula (31) can also be deduced as a special case of our earlier result [5, p. 214, Eq. (2.14)].

Since the function $\Xi(z)$, defined by (25) under the conditions in (30), vanishes exponentially when $|z| \rightarrow \infty$ and $\Omega < 0$, we can appeal to (31) in order to derive a generalization of (7) given by

$$(32) \quad \int_0^{\infty} e^{-st} t^{\mu-1} \Xi(\alpha t^{\rho}) dt = \rho^{-1} \alpha^{-\mu/\rho} \sum_{n=0}^{\infty} \frac{(-s \alpha^{-1/\rho})^n}{n!} \\ \frac{\prod_{j=1}^q \Gamma[\beta_j + B_j(\mu+n)/\rho]}{\prod_{j=1}^r \Gamma[\gamma_j - C_j(\mu+n)/\rho] \prod_{j=1}^p \Gamma[\alpha_j + A_j(\mu+n)/\rho]},$$

which holds true whenever the series on the right has a meaning, $\rho > 0$, $\text{Re}(s) > 0$, $\text{Re}(\alpha) > 0$, and $\text{Re}(\mu) > \rho\omega$, ω being given by (28).

A similar generalization of our formula (10) can be obtained from the definition (25) by applying the method of derivation of (10) *mutatis mutandis*, and we omit the details involved. Instead, however, we find it worthwhile to record an interesting unification of (7), (10) and (32), involving the general H -function, in the following form:

$$(33) \quad \int_0^{\infty} e^{-ts} t^{\mu-1} H_{p,q}^{m,n} \left[\alpha t^{\rho} \left| \begin{matrix} (\alpha_1, A_1), \dots, (\alpha_p, A_p) \\ (\beta_1, B_1), \dots, (\beta_q, B_q) \end{matrix} \right. \right] dt \\ = s^{-\mu} H_{p+1,q}^{m,n+1} \left[\alpha s^{-\rho} \left| \begin{matrix} (1-\mu, \rho), (\alpha_1, A_1), \dots, (\alpha_p, A_p) \\ (\beta_1, B_1), \dots, (\beta_q, B_q) \end{matrix} \right. \right], \text{ if } \rho > 0, \\ = s^{-\mu} H_{p,q+1}^{m+1,n} \left[\alpha s^{-\rho} \left| \begin{matrix} (\alpha_1, A_1), \dots, (\alpha_p, A_p) \\ (\mu, -\rho), (\beta_1, B_1), \dots, (\beta_q, B_q) \end{matrix} \right. \right], \text{ if } \rho < 0,$$

provided for example, that $\text{Re}(s) > 0$, $|\arg(\alpha)| < \frac{1}{2} \Theta\pi$, where

$$(34) \quad \Theta \equiv \sum_{j=1}^n A_j - \sum_{j=n+1}^p A_j + \sum_{j=1}^m B_j - \sum_{j=m+1}^q B_j > 0,$$

and either

$$(35) \quad \text{Re}(\mu) > -\rho \min_{1 \leq j \leq m} \left\{ \text{Re} \left(\frac{\beta_j}{B_j} \right) \right\}, \text{ if } \rho > 0,$$

or

$$(36) \quad \operatorname{Re}(\mu) > -\rho \max_{1 \leq j \leq n} \left\{ \operatorname{Re} \left(\frac{\alpha_j - 1}{A_j} \right) \right\}, \text{ if } \rho < 0.$$

Remark 3. Evidently, for $\mu = 1$, our last result (33) would provide us with the Laplace transform of the general $H_{p,q}^{m,n}[\alpha t^\rho | \dots]$ function for real exponents $\rho \neq 0$, and under the appropriately modified conditions stated above. Furthermore, (33) can be rewritten in terms of the G -function whenever ρ , the A 's and the B 's are *rational* numbers.

We should like to conclude by remarking that a fairly wide variety of special functions, which occur rather frequently in problems of applied mathematics and mathematical analysis, can be expressed in terms of the H -function which evidently reduces, when $A_j = B_k = 1$, $j = 1, \dots, p$; $k = 1, \dots, q$, to the relatively more familiar G -function of Meijer (or the E -function of Mac Robert) and which includes, as a special case, the generalized hypergeometric ${}_p\psi_q[z]$ function of E. M. Wright (cf. [6]). Thus the Laplace transform pair, given by Equation (33), would apply not only to the simpler special functions of mathematical physics such as the Bessel functions $J_\nu(z)$ and $I_\nu(z)$, the Legendre functions $P_\nu^\mu(z)$ and $Q_\nu^\mu(z)$, the Whittaker function $M_{k,\mu}(z)$, the Bessel-Wright function $J_\nu^\mu(z)$, the classical orthogonal polynomials of Hermite, Jacobi (and, of course, Gegenbauer, Legendre, and Tchebycheff), Laguerre, *et cetera*, which are particular cases of the generalized hypergeometric function ${}_pF_q[z]$ or ${}_p\psi_q[z]$, but also to the more involved Bessel functions $K_\nu(z)$ and $Y_\nu(z)$, the Whittaker function $W_{k,\mu}(z)$, their various combinations and several other related functions. Indeed, by making use of extensive tables of functions (and their various products and other combinations) expressible in terms of the E - or G -function or the generalized hypergeometric functions (see, for instance, [1], Vol. I, pp. 215—222; [2], Vol. I, pp. 373—384; [2], Vol. II, pp. 433—444), the interested reader can easily construct from the results of this paper a considerably large tableau of higher transcendental functions and their respective images under the Laplace transform.

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