ON CERTAIN TRANSFORMATIONS IN OPERATIONAL CALCULUS

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1. Introduction. Mainra [3] has defined

(1.1)
$$\widetilde{\omega}_{\mu,\lambda}^{\nu}(x) = \int_{0}^{\infty} \widetilde{\omega}_{\mu,\lambda}(xy) J_{\nu}(y) \sqrt{y} dy, R(\mu,\lambda,\nu) \geqslant -\frac{1}{2};$$

and has established that $\tilde{\omega}_{\mu,\lambda}^{\nu}(x)$ is a Fourier kernel and plays the role of a transform.

Bhatnagar [1] has defined

(1.2)
$$\tilde{\omega}_{\mu, \lambda, \nu}(x) = \int_{0}^{\infty} \tilde{\omega}_{\mu, \lambda}(x/t) J_{\nu}(t) \frac{dt}{\sqrt{t}}, (\mu, \lambda, \nu) \geqslant -\frac{1}{2};$$

and proved that this is also a Fourier kernel and plays the role of a transform.

Further the integral equation

(1.3)
$$\psi(p) = p \int_{0}^{\infty} e^{-pt} f(t) dt, \ R(p) > 0;$$

is symbolically denoted as

$$\psi(p) \stackrel{.}{\rightleftharpoons} f(t)$$
 or $f(t) \stackrel{.}{\rightleftharpoons} \psi(p)$,

where $\psi(p)$ is known as the classical Laplace transform of f(t).

In this paper, we shall establish that if $f(t) = \psi(p)$, $g(t) = \phi(p)$ and the two originals f(t), g(t) are related by any of the Fourier kernels defined above, then their images $\psi(p)$, $\phi(p)$ are also related by one of the Fourier kernels different from the first one and vice—versa. We may add that converse theorems can be proved either under similar conditions or under modified conditions. As $\tilde{\omega}_{\mu,\lambda}^{\nu}(x)$ and $\tilde{\omega}_{\mu,\lambda,\nu}(x)$ involve Bessel functions, it is well-known that $R(\mu,\lambda,\nu)>-1$. So at some places in the conditions we have not mentioned it as this is implied. For self-reciprocity we take

$$R(\mu, \lambda, \nu) \geqslant -\frac{1}{2}$$
.

(2.1) 2. Let
$$f(t) = \psi(p)$$

(2.2) and
$$F(at) = x(p/a)$$
.

Applying Goldstein's theorem [2] to (2.1) and (2.2), we get

(2.3)
$$\int_0^\infty f(x) \times (x/a) \frac{dx}{x} = \int_0^\infty \psi(x) F(ax) \frac{dx}{x},$$

provided the necessary change in the order of integrations involved are permissible and the integrals converge.

Putting $a = \frac{1}{p}$ and interpreting with the help of (2.2), we get

(2.4)
$$\int_0^\infty F(y/x) f(x) \frac{dx}{x} \stackrel{\sim}{=} \int_0^\infty F(x/p) \psi(x) \frac{dx}{x}.$$

Let us put $F(t) = t^m \tilde{\omega}_{\mu, \lambda} (1/t)$ in (2.4). We get

$$(2.5) \quad y^m \int_0^\infty \tilde{\omega}_{\mu, \lambda}(x/y) f(x) x^{-1-m} dx \stackrel{:}{\rightleftharpoons} p^{-m} \int_0^\infty \tilde{\omega}_{\mu, \lambda}(p/x) \psi(x) x^{m-1} dx.$$

Now Mitra and Bose [4] have proved that if $f_1(t) = \psi_1(p)$,

(2.6) then
$$t^{\nu+1} \int_0^\infty J_{\nu}(tz) f_1(z) z^{-\nu} dz = p^{1-\nu} \int_0^\infty J_{\nu+1}(pz) \psi_1(z) z^{\nu} dz$$
,
$$R(\nu) \geqslant -\frac{1}{2}.$$

Making use of (2.5) in (2.6), we get

$$t^{\nu+1} \int_0^\infty J_{\nu}(tz) z^{m-\nu} dz \int_0^\infty \tilde{\omega}_{\mu, \lambda}(x/z) f(x) x^{-1-m} dx$$

$$\stackrel{:}{=} p^{1-\nu} \int_0^\infty J_{\nu+1}(pz) z^{\nu-m} dz \int_0^\infty \tilde{\omega}_{\mu, \lambda}(z/x) \psi(x) x^{m-1} dx.$$

Changing the order of integrations on both the sides and putting

$$m = \nu - \frac{1}{2}, \text{ we get}$$

$$t^{\nu+1} \int_{0}^{\infty} f(x) x^{-\frac{1}{2} - \nu} dx \int_{0}^{\infty} \tilde{\omega}_{\mu, \lambda} (x/z) J_{\nu} (tz) z^{-\frac{1}{2}} dz$$

$$\stackrel{:}{=} p^{1-\nu} \int_{0}^{\infty} \psi(x) x^{\nu-\frac{3}{2}} dx \int_{0}^{\infty} \tilde{\omega}_{\mu, \lambda} (z/x) J_{\nu+1} (pz) \sqrt{z} dz,$$

provided $x^{\nu-\frac{1}{2}}\psi(x)$ and $x^{-\nu}f(x)$ are bounded and absolutely integrable in $(0, \infty)$ and $R\left(\mu, \nu, \lambda, +\frac{1}{2}\right) > 0$.

On writing z/t for z on the l. h. s. and z/p for z on the r. h. s. and making use of the definitions of $\tilde{\omega}_{\mu, \lambda, \nu}(x)$ and $\tilde{\omega}_{\mu, \lambda}^{\nu}(x)$, the above can be written after a little change as

$$(2.7) t^{\nu+\frac{1}{2}} \int_{0}^{\infty} \tilde{\omega}_{\mu, \lambda, \nu}(tx) f(x) x^{-\frac{1}{2}-\nu} dx p^{-\nu-\frac{1}{2}} \int_{0}^{\infty} \tilde{\omega}_{\mu, \lambda}^{\nu+\frac{1}{2}}(x/p) \psi(1/x) x^{-\frac{1}{2}-\nu} dx.$$

Let us put
$$\phi(p) = p^{-\nu - \frac{1}{2}} \int_{0}^{\infty} \tilde{\omega}_{\mu, \lambda}^{\nu + 1}(x/p) \psi(1/x) x^{-\frac{1}{2} - \nu} dx = \text{r. h. s. of } (2.7)$$

i. e. $p^{-\nu - \frac{1}{2}} \phi(1/p)$ is the $\tilde{\omega}_{\mu,\lambda}^{\nu+1}(x)$ transform of $p^{-\nu - \frac{1}{2}} \psi(1/p)$.

Taking $g(t) = \phi(p)$, we get

$$g(t) = t^{\nu + \frac{1}{2}} \int_{0}^{\infty} \tilde{\omega}_{\mu, \lambda, \nu} (tx) f(x) x^{-\frac{1}{2} - \nu} dx = \text{l.h.s. of (2.7)},$$

which shows that $t^{-\nu-\frac{1}{2}}g(t)$ is the $\tilde{\omega}_{\mu,\lambda,\nu}(x)$ transform of $t^{-\nu-\frac{1}{2}}f(t)$. Thus:

Theorem 1. Let $f(t) \stackrel{\cdot}{=} \psi(p)$, $g(t) \stackrel{\cdot}{=} \phi(p)$ and $p^{-\nu - \frac{1}{2}} \phi(1/p)$ be the $\tilde{\omega}_{u,\lambda}^{\nu+1}(x)$ transform of $p^{-\nu - \frac{1}{2}} \psi(1/p)$.

Then $t^{-\nu-\frac{1}{2}}g(t)$ will be the $\tilde{\omega}_{\mu, \lambda, \nu}(x)$ transform of $t^{-\nu-\frac{1}{2}}f(t)$, provided $x^{\nu-\frac{1}{2}}\psi(x)$ and $x^{-\nu}f(x)$ are bounded and absolutely integrable in $(0, \infty)$, g(t) and $t^{\nu+\frac{1}{2}}\int_{0}^{\infty}\tilde{\omega}_{\mu, \lambda, \nu}(tx)f(x) x^{-\frac{1}{2}-\nu}dx$ are continuous functions of t in (0, t) and $R(\mu, \lambda, \nu, +\frac{1}{2}) > 0$.

Further let $x^{-\nu-\frac{1}{2}}\psi(1/x)$ be $R^{\nu+1}_{\mu,\lambda}$; then from (2.7), we have $t^{\nu+\frac{1}{2}}\int_{x}^{\infty} \tilde{\omega}_{\mu,\lambda,\nu}(tx) f(x) x^{-\frac{1}{2}-\nu} dx = \psi(p). \text{ But } \psi(p) = f(t).$

(2.8)
$$f(t) = t^{\nu + \frac{1}{2}} \int_{0}^{\infty} \tilde{\omega}_{\mu, \lambda, \nu}(tx) f(x) x^{-\frac{1}{2} - \nu} dx,$$

which shows that $x^{-\nu-\frac{1}{2}}f(x)$ is $R_{\mu,\lambda,\nu}$, provided both the sides of (2.8) are continuous functions of t.

Thus

Cor. 1. Let $f(t) = \psi(p)$ and $x^{-\nu - \frac{1}{2}} \psi(1/x)$ be $R_{\mu, \lambda}^{\nu + 1}$. Then $x^{-\nu - \frac{1}{2}} f(x)$ will be $R_{\mu, \lambda, \nu}$, provided the conditions of the theorem are satisfied and $\left[f(t) - t^{\nu + \frac{1}{2}} \int_{0}^{\infty} \tilde{\omega}_{\mu, \lambda, \nu}(tx) f(x) x^{-\frac{1}{2} - \nu} dx \right]$ is a continuous function of t.

Again putting $\lambda = v + 1$ in (2.7), we can write it after changing the sides as

$$(2.9) p^{-\nu - \frac{1}{2}} \int_{0}^{\infty} J_{\mu}(x/p) \sqrt{x/p} \psi(1/x) x^{-\frac{1}{2} - \nu} dx$$

$$= p \int_{0}^{\infty} e^{-pt} t^{\nu + \frac{1}{2}} dt \int_{0}^{\infty} \tilde{\omega}_{\mu, \nu + 1, \nu}(tx) f(x) x^{-\frac{1}{2} - \nu} dx$$

$$= \int_{0}^{\infty} \tilde{\omega}_{\mu, \nu + 1, \nu}(x) x^{-\nu - \frac{1}{2}} \left[p \int_{0}^{\infty} e^{-pt} f(x/t) t^{2\nu} dt \right] dx$$

$$= \int_{0}^{\infty} \tilde{\omega}_{\mu, \nu + 1, \nu}(x) x^{\nu - \frac{1}{2}} \left[px \int_{0}^{\infty} e^{-pxt} f(1/t) t^{2\nu} dt \right] dx$$

$$= \int_0^\infty \widetilde{\omega}_{\mu, \nu+1, \nu}(x) F(px) x^{\nu-\frac{1}{2}} dx, \text{ [assuming } F(p) \stackrel{...}{=} t^{2\nu} f(1/t)].$$

Let $\xi(p)$ be the Hankel-transform of order μ of $x^{-\nu-\frac{1}{2}}\psi\left(\frac{1}{x}\right)$.

Then from (2.9), we get

$$p^{-\nu - \frac{1}{2}} \xi (1/p) = \int_{0}^{\infty} \tilde{\omega}_{\mu, \nu+1, \nu}(x) F(px) x^{\nu - \frac{1}{2}} dx$$
$$\xi (1/p) = \int_{0}^{\infty} \tilde{\omega}_{\mu, \nu+1, \nu}(x/p) F(x) x^{\nu - \frac{1}{2}} dx,$$

or

which shows that $\xi(p)$ is the $\tilde{\omega}_{\mu, \nu+1, \nu}(x)$ transform of $p^{\nu-\frac{1}{2}}F(p)$. Thus:

Cor. 2. Let $f(t) = \psi(p)$, $t^{2\nu} f(1/t) = F(p)$ and $\xi(p)$ be the Hankel-transform of order μ of $x^{-\nu - \frac{1}{2}} \psi(1/x)$.

Then $\xi(p)$ will be the $\tilde{\omega}_{\mu, \nu+1, \nu}(x)$ transform of $p^{\nu-\frac{1}{2}} F(p)$, provided the conditions of the theorem hold good.

Further, if we put $t^{2\nu} f(1/t) = f(t)$ in the above corollary, we get $\psi(p) = F(p)$. Hence we have

Cor. 3. Let $f(t) = t^{2\nu} f(1/t)$, $f(t) = \psi(p)$ and $\xi(p)$ be the Hankel-transform of order μ of $x^{-\nu - \frac{1}{2}} \psi(1/x)$.

Then $\xi(p)$ will be the $\tilde{\omega}_{\mu, \nu+1, \nu}(x)$ transform of $p^{\nu-\frac{1}{2}} \psi(p)$ provided the conditions of the theorem hold good.

Again, we know that if $f(t) = \psi(p)$,

(2.10) then
$$\left(t\frac{d}{dt}\right)^n f(t) \stackrel{...}{=} (-1)^n \left(p\frac{d}{dp}\right)^n \psi(p).$$

Making use of the above relation in (2.7), we get

$$(2.11) t^{\nu+\frac{1}{2}} \int_{0}^{\infty} \tilde{\omega}_{\mu, \lambda, \nu}(tx) \left[\left(x \frac{d}{dx} \right)^{n} f(x) \right] x^{-\nu-\frac{1}{2}} dx \stackrel{\cdot}{=} p^{-\nu-\frac{1}{2}} \int_{0}^{\infty} \tilde{\omega}_{\mu, \lambda}^{\nu+\frac{1}{2}}(x/p) \times \left[(-1)^{n} \left(\frac{1}{x} \frac{d}{d\frac{1}{x}} \right)^{n} \psi(1/x) \right] x^{-\nu-\frac{1}{2}} dx,$$

provided $x^{-\nu - \frac{1}{2}} \left(x \frac{d}{dx} \right)^n f(x)$ and $x^{-\nu - \frac{1}{2}} \left(\frac{1}{x} \frac{d}{dx} \right)^n \psi(1/x)$ are bounded and

absolutely integrable in $(0, \infty)$.

Now proceeding in the same manner as in theorem 1, we have the following theorem:

Theorem 2. Let $f(t) = \psi(p)$, $g(t) = \phi(p)$ and $p^{-\nu - \frac{1}{2}} \phi(1/p)$ be the $\tilde{\omega}_{\mu, \lambda}^{\nu+1}(x)$ transform of $(-1)^n p^{-\nu - \frac{1}{2}} \left(\frac{1}{p} \frac{d}{d \cdot \frac{1}{p}}\right)^n \psi(1/p)$.

Then $t^{-\nu-\frac{1}{2}}g(t)$ will be the $\tilde{\omega}_{\mu, \lambda, \nu}(x)$ transform of

$$t^{-\nu-\frac{1}{2}}\left(t\frac{d}{dt}\right)^n f(t)$$
, provided $x^{-\nu-\frac{1}{2}}\left(x\frac{d}{dx}\right)^n f(x)$ or $x^{-\nu}\left(x\frac{d}{dx}\right)^n f(x)$

and $x^{-\nu-\frac{1}{2}} \left(\frac{1}{x} \frac{d}{d\frac{1}{x}}\right)^n \psi(1/x)$ are bounded and absolutely integrable in $(0, \infty)$.

Cor. Let
$$f(t) = \psi(p)$$
 and $(-1)^n x^{-\nu - \frac{1}{2}} \left(\frac{1}{x} \frac{d}{d \frac{1}{x}} \right)^n \psi(1/x)$ be $R_{\mu, \lambda}^{\nu+1}$.

Then $t^{-\nu-\frac{1}{2}}\left(t\frac{d}{dt}\right)^n f(t)$ will be $R_{\mu, \lambda, \nu}$, provided the conditions of the above theorem hold good.

Similarly,

Theorem 3. Let $f(t) = \psi(p)$, $g(t) = \phi(p)$ and $p^{-\nu - \frac{1}{2}} \phi(1/p)$ be the $\tilde{\omega}_{\mu, \lambda}^{\nu+1}(x)$ transform of $p^{-\nu - n - \frac{1}{2}} \psi(1/p)$.

Then $t^{-\nu-\frac{1}{2}}g(t)$ will be the $\tilde{\omega}_{\mu, \lambda, \nu}(x)$ transform of $t^{-\nu-\frac{1}{2}}f^n(t)$, provided $f^r(0)=0, r=0,1,\ldots, n-1; x^{-\nu}f^n(x)$ and $x^{-\nu-n}\psi(1/x)$ are bounded and absolutely integrable in $(0, \infty)$.

Cor. Let $f(t) = \psi(p)$, f'(0) = 0, $r = 0,1,2,\ldots, n-1$, and $x^{-\nu-n-\frac{1}{2}}\psi(1/x)$ be $R_{\mu,\lambda}^{\nu+1}$. Then $t^{-\nu-\frac{1}{2}}f^n(t)$ will be $R_{\mu,\lambda,\nu}$, provided the conditions of the above theorem hold good.

Here we make use of the following relation in place of (2.10), i.e. If $f(t) = \psi(p)$, then $f^n(t) = p^n \psi(p)$, provided $f^r(0) = 0$, $r = 0,1,2,\ldots, n-1$, and proceed in the same way as in theorem 2.

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REFERENCES

- [1] Bhatnagar, K. P., Bull. Cal. Math. Soc., 46, (1954), 179.
- [2] Goldstein, S., Proc. Lond. Math. Soc. (2), 34. (1932), 103.
- [3] Mainra, V. P., Bull. Cal. Math. Soc. vol. 50, No. 3. (1958), 123.
- [4] Mitra, S. C., and Bose, B. N., Acta Mathematica, 88. (1952), 227-240.
- [5] Titchmarsh, E. C., (1948), Theory of Fourier Integrals, (Oxford).
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