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HOLOMORPHICALLY SEMI-SYMMETRIC CONNEXIONS

§1. Object of the paper. — Let us consider a n -dimensional differentiable manifold M . The connexion of M is termed semi-symmetric if its torsion tensor S_{ij}^k has the form

$$(1.1) \quad S_{ij}^k = \delta_i^k S_j - \delta_j^k S_i,$$

where S_i are the components of a vector field. The geometrical meaning of the semi-symmetric connexion was given by E. Bartolotti [1] and it consists in the following. Let U^i and V^j be two vectors. The vectors $S_{ij}^k U^i V^j$, U^k and V^k are, in the general case, linearly independent. But if

$$(1.2) \quad S_{ij}^k U^i V^j = a U^k + b V^k$$

for every U and V , where a and b are scalars, then S_{ij}^k has the form (1.1) and conversely.

Let us now suppose that in the manifold M a tensor field $F_j^i \neq \delta_j^i$ satisfying

$$(1.3) \quad F_j^i F_k^j = \omega \delta_k^i, \quad \omega = -1 \quad \text{or} \quad \omega = +1$$

is given. We suppose that the structure F_j^i is integrable. Then, if $\omega = -1$, the dimension of M is even, and M is a complex space [2]. If $\omega = +1$, M is a locally product space.

A semi-symmetric connexion is generalized on a complex space in [3] and on a locally product space in [4]. In [3], two connexions of a Kähler space are investigated, each of which can be considered a generalization of a semi-symmetric connexion. Furthermore, each of these connexions is a metric F -connexion.

The purpose of the present paper is the generalization of the semi-symmetric connexion, too. We generalize Bartolotti's geometrical property in the case that the manifold M has the integrable structure (1.3). Namely, we consider the skew-symmetric tensor S_{ij}^k satisfying condition

$$(1.4) \quad S_{ij}^k U^i F_a^j U^a = a U^k + b F_a^k U^a$$

instead of (1.2). We shall show, in § 2, that if S_i^k satisfies (1.4) for every U , then

$$(1.5) \quad S_{ij}^k = \delta_j^k S_i - \delta_i^k S_j - \omega S_a F_i^a F_j^k + \omega S_a F_j^a F_i^k + \\ + \frac{1}{2} (\delta_i^a \delta_j^b + \omega F_i^a F_j^b) W_{ab}^k$$

where S_i is a vector field and W_{ij}^k is an arbitrary skew-symmetric tensor. The conversely being obvious, we have:

Theorem 1. — *The skew-symmetric tensor S_i^k satisfies condition (1.4) for every U^i if and only if it has the form (1.5).*

Definition. — We term an affine connexion having a torsion tensor of the form (1.5) a *holomorphically semi-symmetric connexion*.

The simplest connexion of the Riemannian space whose torsion tensor has the form (1.5), is

$$(1.6) \quad \Gamma_{ij}^k = \left\{ \begin{matrix} k \\ ij \end{matrix} \right\} + S_i \delta_j^k - \omega S_a F_i^a F_j^k$$

This connexion is an F -connexion (*i. e.* $\nabla F=0$), but it is not a metric one.

Let R_{rkj}^i and K_{rkj}^i be the components of the curvature tensors of the connexions (1.6) and $\left\{ \begin{matrix} k \\ ij \end{matrix} \right\}$ respectively and $R_{rk} = R_{rka}^a$, $K_{rk} = K_{rka}^a$. We shall prove, in §3,

Theorem 2. — *We have*

$$(1.7) \quad R_{rkj}^i + \frac{1}{2} (R_{kj} - R_{jk}) \delta_r^i + \frac{1}{2} (R_{ka} F_j^a - R_{ja} F_k^a) F_r^i = K_{rkj}^i + K_{ka} F_j^a F_r^i$$

in the Kähler space, and

$$(1.8) \quad R_{rkj}^i + \frac{1}{2} (R_{kj} - R_{jk}) \delta_r^i - \frac{1}{2} (R_{ka} F_j^a - R_{ja} F_k^a) F_r^i = K_{rkj}^i$$

in the locally decomposable Riemannian space.

Finally, in the case of a locally decomposable Riemannian space, we consider the connexion

$$(1.9) \quad \Gamma_{ij}^k = \left\{ \begin{matrix} k \\ ij \end{matrix} \right\} + S_i \delta_j^k + g_{ij} S^i - F_i^a F_j^k S_a + F_{ij} F_a^k S^a$$

where g_{ij} are the components of the metric tensor and $F_{rj} = F_r^a g_{aj}$. The torsion tensor of this connexion has the form (1.5), too. It is also an F -connexion. We shall prove, in §4,

Theorem 3. — The curvature tensor R_{rkj}^i of the connexion (1.9) satisfies the relation

$$(1.10) \left\{ \begin{aligned} &R_{rkj}^i - \frac{1}{2} \delta_r^i (R_{ab} - R_{ba}) O_{rk}^{ab} + \frac{1}{2} F_r^i R_{pq} (O_{ak}^{pq} F_j^a - O_{aj}^{pq} F_k^a) \\ &+ \frac{1}{2} R_{pq} (g_{kr} g^{ia} O_{aj}^{pq} - g_{jr} g^{ia} O_{ak}^{pq} - F_{kr} F^{ia} O_{aj}^{pq} + F_{jr} F^{ia} O_{ak}^{pq}) \\ &+ \frac{1}{2} (\alpha R + \beta R^*) (g_{jr} \delta_k^i - g_{kr} \delta_j^i + F_{jr} F_k^i - F_{kr} F_j^i) \\ &+ \frac{1}{2} (\beta R + \alpha R^*) (g_{jr} \delta_k^a - g_{kr} \delta_j^a + F_{jr} F_k^a - F_{kr} F_j^a) F_a^i = \\ &= K_{rkj}^i - \frac{1}{2} (g_{jr} K_k^i - g_{kr} K_j^i + F_{jr} F_a^i K_k^a - F_{kr} F_a^i K_j^a) \\ &+ \frac{1}{2} (\alpha K + \beta K^*) (g_{jr} \delta_k^i - g_{kr} \delta_j^i + F_{jr} F_k^i - F_{kr} F_j^i) \\ &+ \frac{1}{2} (\beta K + \alpha K^*) (g_{jr} \delta_k^a - g_{kr} \delta_j^a + F_{jr} F_k^a - F_{kr} F_j^a) F_a^i, \end{aligned} \right.$$

where we use the notations (2.3), (4.3) and

$$\alpha = \frac{n-2}{(n-2)^2 - \varphi^2}, \quad \beta = -\frac{\varphi}{(n-2)^2 - \varphi^2}, \quad \varphi = F_a^a.$$

§2. Proof of Theorem 1. — To eliminate a and b from (1.4), we multiply (1.4) by $\delta_{[l}^r \delta_i^s \delta_{k]}^m U^l F_a^i U^a$. We get

$$\delta_{[l}^r \delta_i^s \delta_{k]}^m S_{ij}^k U^l U^i F_a^j F_b^k U^a U^b = 0$$

This condition being satisfied for every U^i , we have

$$\delta_{[l}^r \delta_i^s \delta_{k]}^m S_{ij}^k F_a^j F_b^k + (l, i, a, b) = 0,$$

where the above expression means that it should sum up all the terms obtained by the cyclic permutation of indices l, i, a, b .

Multiplying this equation by $F_p^a F_q^b$, and then contracting with respect to r and l and with respect to s and q , we obtain

$$\begin{aligned} &[n(n+2)\omega - \varphi^2] (\omega S_p^m + S_{ab}^m F_p^a F_b^b) + 2\varphi (\omega S_p^r + S_{ab}^r F_p^a F_b^b) F_r^m \\ &+ \delta_p^m [\omega(n+2)\Theta_a F_i^a + \omega\varphi\Theta_i - \omega^2(n+2)\psi_i - \omega\varphi\psi_a F_i^a] \\ &- \delta_i^m [\omega(n+2)\Theta_a F_p^a + \omega\varphi\Theta_p - \omega^2(n+2)\psi_p - \omega\varphi\psi_a F_p^a] \\ &- F_p^m [\omega(n+2)\Theta_i + \varphi\Theta_a F_i^a - \omega(n+2)\psi_a F_i^a - \omega\varphi\psi_i] \\ &+ F_i^m [\omega(n+2)\Theta_p + \varphi\Theta_a F_p^a - \omega(n+2)\psi_a F_p^a - \omega\varphi\psi_p] = 0, \end{aligned}$$

where

$$\varphi = F_a^a, \quad S_{ir}^r = \psi_i, \quad S_{ib}^a F_a^b = \Theta_i.$$

Putting

$$S_i = \omega (n+2) \Theta_a F_i^a + \omega \varphi \Theta_i - (n+2) \psi_i - \omega \varphi \psi_a F_i^a,$$

we reduce the preceding equation to the form:

$$(2.1) \quad [n(n+2) \omega - \varphi^2] (\omega S_{ip}^m + S_{ab}^m F_p^a F_i^b) + 2\varphi (\omega S_{ip}^r + S_{ab}^r F_p^a F_i^b) F_r^m = \\ = \delta_i^m S_p - \delta_p^m S_i - \omega S_a F_p^a F_i^m + \omega S_a F_i^a F_p^m.$$

In complex spaces $\varphi=0$. Also, in some locally product spaces, $\varphi=0$. Namely, in every locally product space the structure tensor F_j^i determines two globally complementary distributions. If the dimensions of the distributions are p and q , then $\varphi=p-q$. We mean here by the distribution of dimension p , p -dimensional plane in the n -dimensional tangent plane at each point of the manifold.

Let us consider firstly, the case $\varphi=0$. The vector field S_i being arbitrary we can put S_i instead of $\bar{S}_i = \frac{1}{2n(n+2)} S_i$. Taking into account the skew-symmetry of S_{jk}^i , we transcribe (2.1) in the following form:

$$(2.2) \quad *O_{pi}^{ab} S_{ab}^m = S_i \delta_p^m - S_p \delta_i^m + \omega S_a F_p^a F_i^m - \omega S_a F_i^a F_p^m,$$

where

$$(2.3) \quad O_{ij}^{ab} = \frac{1}{2} (\delta_i^a \delta_j^b + \omega F_i^a F_j^b), \quad *O_{ij}^{ab} = \frac{1}{2} (\delta_i^a \delta_j^b - \omega F_i^a F_j^b).$$

The condition

$$O_{ij}^{pi} (S_i \delta_p^m - S_p \delta_i^m + \omega S_a F_p^a F_i^m - \omega S_a F_i^a F_p^m) = 0$$

being satisfied, we can apply the well known lemma ([2], p. 133). Thus, (2.2) permits a solution, and the general solution is given by (1.5).

Now suppose that $\varphi \neq 0$. This can be the case only if $\omega=1$. Taking this into account, we put

$$(2.4) \quad S_{ip}^m + S_{ab}^m F_p^a F_i^b = T_{ip}^m$$

$$(2.5) \quad \delta_i^m S_p - \delta_p^m S_i - S_a F_p^a F_i^m + S_a F_i^a F_p^m = V_{ip}^m.$$

Then (2.1) obtains the form:

$$(2.6) \quad [n(n+2) - \varphi^2] T_{ip}^m + 2\varphi T_{ip}^r F_r^m = V_{ip}^m.$$

Multiplying (2.6) by F_m^s , we get

$$(2.7) \quad 2\varphi T_{ip}^m + [n(n+2) - \varphi^2] T_{ip}^r F_r^m = V_{ip}^r F_r^m.$$

Multiplying (2.6) by $[n(n+2)-\varphi^2]$ and (2.7) by 2φ , and subtracting the second from the first, we obtain

$$\{[n(n+2)-\varphi^2]^2-4\varphi^2\} T_{ip}^m = [n(n+2)-\varphi^2] V_{ip}^m - 2\varphi V_{ip}^a F_a^m.$$

This can be written in the form:

$$(2.8) \quad T_{ip}^m = AV_{ip}^m + BV_{ip}^a F_a^m,$$

where

$$A = \frac{n(n+2)-\varphi^2}{[n(n+2)-\varphi^2]^2-4\varphi^2}, \quad B = -\frac{2\varphi}{[n(n+2)-\varphi^2]^2-4\varphi^2}.$$

It can be easily verified that

$$[n(n+2)-\varphi^2]^2-4\varphi^2 \neq 0.$$

The tensor V_{ip}^m having the form (2.5), we find

$$\begin{aligned} AV_{ip}^m + BV_{ip}^a F_a^m &= (AS_p - BS_a F_p^a) \delta_i^m - (AS_i - BS_a F_i^a) \delta_p^m - \\ &- (AS_a F_p^a - BS_p) F_i^m + (AS_a F_i^a - BS_i) F_p^m. \end{aligned}$$

Consequently, putting

$$V_i = AS_i - BS_a F_i^a,$$

and taking into account (2.4) and the skew-symmetry of S_{ij}^k , we express (2.8) in the form:

$$*O_{ip}^{ba} S_{ba}^m = \frac{1}{2} (V_p \delta_i^m - V_i \delta_p^m - V_a F_p^a F_i^m + V_a F_i^a F_p^m),$$

i. e. in the form (2.2). This completes the proof of Theorem 1.

§ 3. Proof of Theorem 2. — Suppose that the Riemannian metric g_{ij} satisfying

$$g_{is} F_j^s F_i^s = g_{ij}, \quad \overset{\circ}{\nabla}_k F_j^i = 0,$$

where $\overset{\circ}{\nabla}$ denotes the covariant differentiation with respect to the Christoffel symbols $\{j_k^i\}$, is given. Then M is, in the case $\omega = -1$, a Kähler space and, in the case $\omega = +1$, a locally decomposable Riemannian space.

Let us consider the connexion

$$(3.1) \quad \Gamma_{jk}^i = \{j_k^i\} + S_j \delta_i^k + P_j F_k^i.$$

If $P_i = -\omega S_a F_i^a$, the connexion (3.1) reduces to (1.6), i. e. to the holomorphically semi-symmetric connexion.

Let ∇ be the operator of the covariant derivative with respect to Γ_{jk}^i . Then $\nabla_k F_j^i = 0$, i. e. (3.1), is an F -connexion.

The curvature tensor R_{rkj}^i of the connexion (3.1) can be expressed in the form:

$$(3.2) \quad R_{rkj}^i = K_{rkj}^i + (\overset{\circ}{\nabla}_k S_j - \overset{\circ}{\nabla}_j S_k) \delta_r^i + (\overset{\circ}{\nabla}_k P_j - \overset{\circ}{\nabla}_j P_k) F_r^i,$$

where K_{rkj}^i is the curvature tensor of the Riemannian connexion $\{j_k\}$.

On the other hand, from the Ricci identity

$$\nabla_s \nabla_r F_j^i - \nabla_r \nabla_s F_j^i = F_j^a R_{asr}^i - F_a^i R_{jrs}^a - \nabla_a F_j^i (\Gamma_{sr}^a - \Gamma_{rs}^a)$$

and the condition $\nabla_k F_j^i = 0$, we get

$$F_j^a R_{asr}^i - F_a^i R_{jrs}^a = 0$$

from which

$$R_{ksr}^i - \omega R_{bsr}^a F_a^i F_k^b = 0.$$

Putting $R_{ksa}^a = R_{ks}$, we find

$$R_{ks} = \omega R_{bst}^a F_a^t F_k^b, \quad R_{pq} F_k^p F_s^q = R_{kqt}^a F_a^t F_s^q.$$

and

$$(3.3) \quad R_{ks} + \omega R_{pq} F_k^p F_s^q = \omega (R_{bst}^a F_a^t F_k^b + R_{kqt}^a F_a^t F_s^q).$$

Similarly, if $K_{ksa}^a = K_{ks}$,

$$(3.4) \quad K_{ks} + \omega K_{pq} F_k^p F_s^q = \omega (K_{bst}^a F_a^t F_k^b + K_{kqr}^a F_a^r F_s^q).$$

On the other hand, we have from (3.2)

$$R_{bst}^a F_a^t F_k^b = K_{bst}^a F_a^t F_k^b + \omega (\overset{\circ}{\nabla}_s S_k - \overset{\circ}{\nabla}_k S_s) + \omega (\overset{\circ}{\nabla}_s P_a - \overset{\circ}{\nabla}_a P_s) F_k^a$$

and

$$R_{kqt}^a F_a^t F_s^q = K_{kqt}^a F_a^t F_s^q + (\overset{\circ}{\nabla}_q S_a - \overset{\circ}{\nabla}_a S_q) F_k^a F_s^q + \omega (\overset{\circ}{\nabla}_q P_k - \overset{\circ}{\nabla}_k P_q) F_s^q.$$

Adding these two equations and taking into account that $\overset{\circ}{\nabla}_k F_j^i = 0$, both for the Kähler and for the locally decomposable Riemannian space, we get

$$\begin{aligned} R_{bst}^a F_a^t F_k^b + R_{kqt}^a F_a^t F_s^q &= K_{bst}^a F_a^t F_k^b + K_{kqt}^a F_a^t F_s^q \\ &+ \omega \overset{\circ}{\nabla}_s S_k - \omega \overset{\circ}{\nabla}_k S_s + \omega \overset{\circ}{\nabla}_s (P_a F_k^a) - \omega (\overset{\circ}{\nabla}_a P_s) F_k^a \\ &+ \overset{\circ}{\nabla}_q (S_a F_k^a) F_s^q - \overset{\circ}{\nabla}_a (S_q F_s^q) F_k^a + \omega (\overset{\circ}{\nabla}_a P_k) F_s^a - \omega \overset{\circ}{\nabla}_k (P_a F_s^a). \end{aligned}$$

In the following, we shall limit ourselves to the holomorphically semi-symmetric connexion (1.6). Then

$$P_i = -\omega S_a F_i^a, \quad S_i = -P_a F_i^a,$$

and the preceding equation reduces to

$$R_{bst}^a F_a^t F_k^b + R_{kbt}^a F_s^b F_a^t = K_{bst}^a F_a^t F_k^b + K_{kqt}^a F_a^t F_s^q.$$

From this, substituting (3.3) and (3.4), we find

$$(3.5) \quad O_{ij}^{ab} K_{ab} = O_{ij}^{ab} R_{ab}.$$

On the other hand, for the connexion (1.6), (3.2) becomes

$$(3.6) \quad R_{rkj}^i = K_{rkj}^i + (\overset{\circ}{\nabla}_k S_j - \overset{\circ}{\nabla}_j S_k) \delta_r^i - \omega (F_j^a \overset{\circ}{\nabla}_k S_a - F_k^a \overset{\circ}{\nabla}_j S_a) F_r^i.$$

This can be written in the form

$$R_{rkj}^i = K_{rkj}^i + \overset{\circ}{\nabla}_k S_a (\delta_j^a \delta_r^i - \omega F_j^a F_r^i) - \overset{\circ}{\nabla}_j S_a (\delta_k^a \delta_r^i - \omega F_k^a F_r^i).$$

From this we find, by contraction with respect to *i* and *j*,

$$*O_{rk}^{ab} \overset{\circ}{\nabla}_b S_a = \frac{1}{2} (K_{rk} - R_{rk}).$$

The condition (3.5) being satisfied, we can apply the lemma [2], p. 133. Thus the preceding equation permits a solution, and the general solution has the form

$$\overset{\circ}{\nabla}_k S_r = \frac{1}{2} (K_{rk} - R_{rk}) + O_{rk}^{ab} W_{ab},$$

where W_{ab} is an arbitrary tensor.

Hence, we can put

$$(3.7) \quad \overset{\circ}{\nabla}_k S_r = \frac{1}{2} (K_{rk} - R_{rk}).$$

Substituting (3.7) into (3.6), we have

$$(3.8) \quad \begin{aligned} R_{rkj}^i + \frac{1}{2} (R_{kj} - R_{jk}) \delta_r^i - \frac{1}{2} \omega (R_{ka} F_j^a - R_{ja} F_k^a) F_r^i = \\ = K_{rkj}^i - \frac{1}{2} \omega (K_{ka} F_j^a - K_{ja} F_k^a) F_r^i. \end{aligned}$$

In the case of a Kähler space

$$K_{ij} + \omega K_{ab} F_i^a F_j^b = 0, \quad \text{that is } K_{aj} F_i^a = -K_{ia} F_j^a,$$

and (3.8) reduces, taking into account that K_{ij} is a symmetric tensor, to

$$(3.9) \quad R_{rkj}^i + \frac{1}{2} (R_{kj} - R_{jk}) \delta_r^i + \frac{1}{2} (R_{ka} F_j^a - R_{ja} F_k^a) F_r^i = K_{rkj}^i + K_{ka} F_j^a F_r^i.$$

In the case of a locally decomposable Riemannian space, $\omega = +1$ and

$$K_{ij} - K_{ab}F_i^a F_j^b = 0, \quad \text{that is } K_{aj}F_i^a - K_{ia}F_j^a = 0,$$

and (3.8) reduces to

$$(3.10) \quad R_{rkj}^i + \frac{1}{2}(R_{kj} - R_{jk})\delta_r^i - \frac{1}{2}(R_{ka}F_j^a - R_{ja}F_k^a)F_r^i = K_{rkj}^i.$$

(3.9) and (3.10) complete the proof of Theorem 2.

§4. Proof of Theorem 3. — In this section we shall consider the connexion (1.9). R_{rkj}^i is the curvature tensor of this connexion, K_{rkj}^i is the curvature tensor of the Riemannian connexion and $\overset{\circ}{\nabla}$ denotes the covariant differentiation with respect to the Riemannian connexion. Then it is easy to see that

$$\begin{aligned} R_{rkj}^i &= K_{rkj}^i + \delta_r^i \overset{\circ}{\nabla}_k S_j + g_{jr} \overset{\circ}{\nabla}_k S^i - F_r^i F_j^a \overset{\circ}{\nabla}_k S_a + F_{jr} F_a^i \overset{\circ}{\nabla}_k S^a \\ &\quad - \delta_r^i \overset{\circ}{\nabla}_j S_k - g_{kr} \overset{\circ}{\nabla}_j S^i + F_r^i F_k^a \overset{\circ}{\nabla}_j S_a - F_{kr} F_a^i \overset{\circ}{\nabla}_j S^a \\ &\quad + g_{jr} S^i S_k - F_{rk} F_j^a S_a S^i - g_{kr} F_a^i F_j^b S^a S_b + F_{jr} F_a^i S^a S_k \\ &\quad - g_{kr} S^i S_j + F_{jr} F_k^a S_a S^i + g_{jr} F_a^i F_k^b S^a S_b - F_{kr} F_a^i S^a S_j. \end{aligned}$$

If we put

$$\psi_{jk} = \overset{\circ}{\nabla}_k S_j + S_j S_k + F_j^a F_k^b S_a S_b, \quad \psi_k^i = g^{ij} \psi_{jk},$$

the tensor R_{rkj}^i may be written in the form

$$\begin{aligned} R_{rkj}^i &= K_{rkj}^i + \delta_r^i \psi_{jk} - \delta_r^i \psi_{kj} - F_r^i F_j^a \psi_{ak} + F_r^i F_k^a \psi_{aj} \\ &\quad + g_{jr} \psi_k^i - g_{kr} \psi_j^i + F_{jr} F_a^i \psi_k^a - F_{kr} F_a^i \psi_j^a. \end{aligned} \quad (4.1)$$

Contracting (4.1) with respect to i and j , we obtain

$$(4.2) \quad R_{rk} = K_{rk} + 2\psi_{rk} - \psi_{kr} + F_k^a F_r^b \psi_{ab} - g_{kr} \psi_a^a - F_{kr} F_a^b \psi_b^a.$$

We now introduce the notations

$$(4.3) \quad R_{rk} g^{rk} = R, \quad K_{rk} g^{rk} = K, \quad R_{rk} F^{rk} = R^* \\ K_{rk} F^{rk} = K^*, \quad F^{rk} = F_a^r g^{ak}, \quad F_{rk} g^{rk} = F_a^a = \varphi = p - q$$

and suppose $p \geq 2, q \geq 2$.

Contracting (4.2) by g^{rk} and F^{rk} , we find that

$$(4.4) \quad \left\{ \begin{aligned} \psi_a^a &= \frac{n-2}{\varphi^2 - (n-2)^2} (R - K) - \frac{\varphi}{\varphi^2 - (n-2)^2} (R^* - k^*), \\ F_b^a \psi_a^b &= -\frac{\varphi}{\varphi^2 - (n-2)^2} (R - K) + \frac{n-2}{\varphi^2 - (n-2)^2} (R^* - k^*). \end{aligned} \right.$$

We note that according to the supposition $p \geq 2$, $q \geq 2$, $\varphi^2 - (n-2)^2 \neq 0$.

On the other hand, we have from (4.1)

$$R_{rk} + R_{ab}F_r^a F_k^b = K_{rk} + K_{ab}F_r^a F_k^b + 2\psi_{ab}(\delta_r^a \delta_k^b + F_r^a F_k^b) - 2g_{kr}\psi_a^a - 2F_{kr}F_b^a \psi_a^b,$$

or

$$(4.5) \quad \psi_{ab}O_{rk}^{ab} = \frac{1}{2} (R_{ab}O_{rk}^{ab} - K_{ab}O_{rk}^{ab} + g_{kr}\psi_a^a + F_{kr}F_b^a \psi_a^b).$$

It is easy to see that

$$(R_{ab}O_{pq}^{ab} - K_{ab}O_{pq}^{ab} + g_{pq}\psi_a^a + F_{pq}F_b^a \psi_a^b) * O_{rk}^{pq} = 0.$$

Consequently, the linear equation (4.5) with the unknown tensor ψ_{rk} permits a solution

$$\psi_{rk} = \frac{1}{2} (R_{ab}O_{rk}^{ab} - K_{ab}O_{rk}^{ab} + g_{rk}\psi_a^a + F_{rk}F_b^a \psi_a^b).$$

Taking into account that for the locally decomposable Riemannian space $K_{ab}O_{rk}^{ab} = K_{rk}$, this solution reduces to

$$(4.6) \quad \psi_{rk} = \frac{1}{2} (R_{ab}O_{rk}^{ab} - K_{rk} + g_{rk}\psi_a^a + F_{rk}F_b^a \psi_a^b).$$

Substituting (4.6) and (4.4) into (4.1), we find (1.10).

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HOLOMORFNO SEMI-SIMETRIČNE KONEKSIJE

Rezime

Holomorfno semi-simetrična koneksija mnogostrukosti je ona koneksija čiji tenzor torzije zadovoljava uslov (1.5). Do tog smo uslova došli tako što smo na prostore sa kompleksnom odnosno lokalnom produkt strukturom proširili geometrijsku interpretaciju semi-simetrične koneksije koju je dao E. Bartolotti [1]. Zatim smo posmatrali dve specijalne holomorfno semi-simetrične koneksije: (1.6) i (1.9). Pokazali smo da tenzori krivina tih koneksija zadovoljavaju uslove (1.7) i (1.8), odnosno (1.10).