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STRONGLY DISTINGUISHED CONNECTIONS IN A RECURRENT K-HAMILTON SPACE

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Abstract

In the n+mK dimensional differentiable manifold E^* (K-Hamilton space) special coordinate transformations are allowed. In $T^*(E^*) \otimes T^*(E^*)$ the metric tensor is given, and using the nonlinear connection $N, T(E^*)$ may be decomposed in K+1 orthogonal subspaces (with respect to G): $T_H(E^*)$ and $T_{K}(E^*) = T_{K}(E^*)$. In $T(E^*) = T_{K}(E^*)$ a strongly distinguished connection is introduced in such a way that Y and $\nabla_X Y$ belong to the same subspace of $T(E^*)$, $\forall X, Y \in T(E^*)$. The law of transformation of connection coefficients is given. For the metrical and recurrent case the connection coefficients and the torsion tensor are determined.

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1. Coordinate bases in $T(E^*)$ and $T^*(E^*)$

Let E^* be an n + mK dimensional differentiable manifold. If u is one point of E^* , then in some local chart u has coordinates

$$u = ((x^i), (p_a^1), (p_a^2), \dots, (p_a^K)) = ((x^i), (p_a^{\alpha})) = (x, p),$$

where
$$(x^i) = (x^1, x^2, \dots, x^n) = (x), (p_a^{\alpha}) = ((p_1^{\alpha}), \dots, (p_m^{\alpha})) = (p^{\alpha})$$
 and $a, b, c, d, e, f = \overline{1, m}, \quad i, j, h, k, l, m = \overline{1, n}, \quad \alpha, \beta, \ \gamma, \delta = \overline{1, K}.$

We shall consider the following transformation of a coordinate system. If $((x^{i'}),(p^{\alpha}_{a'}))=(x',p')$ are the coordinates of the same point u in the new coordinate system, then

(1.1)
$$\begin{aligned} & (\mathbf{a}) \ \ x^{i'} = x^{i'}(x^1, \dots x^n) & \operatorname{rank} \left[\partial x^{i'} / \partial x^i \right] = n \\ & (\mathbf{b}) \ \ p^{\alpha}_{a'} = M^{\alpha}_{a'} \left(x^1, \dots, x^n \right) p^{\alpha}_a & \operatorname{rank} \left[\partial p^{\alpha}_{a'} / \partial p^{\alpha}_a \right] = m. \end{aligned}$$

The Einstein summation convention will be used for all three kinds of indices, except when the index is in brackets. If (1.1) is valid, then an inverse transformation exists

(1.2) (a)
$$x^i = x^i(x^{1'}, \dots x^{n'})$$
 (b) $p_a^{\alpha} = M_a^{\alpha'}(x^{1'}, \dots, x^{n'})p_{\alpha'}^{\alpha}$.

The natural basis $\bar{B} = \{(\partial_i), (\partial_1^a), \dots, (\partial_K^a)\}$ of $T(E^*)$ is formed by n vectors of type $\partial_i = \partial/\partial x^i$ and $m \cdot K$ vectors of type $\partial_\alpha^a = \partial/\partial p_\alpha^\alpha$. Any vector field $X \in T(E^*)$ may be written in the form

(1.3)
$$X = \bar{X}^i \partial_i + \bar{X}_a^{\alpha} \partial_{\alpha}^a.$$

With respect to the coordinate transformations (1.1) and (1.2) the basic vectors of \overline{B} have the following law of transformation

$$\begin{bmatrix}
\partial_{i} \\
\partial_{1}^{a} \\
\vdots \\
\partial_{K}^{a}
\end{bmatrix} = \begin{bmatrix}
\frac{\partial x^{i'}}{\partial x^{i}} & (\partial_{i} M_{a'}^{a}) p_{a}^{1} & \dots & (\partial_{i} M_{a'}^{a}) p_{a}^{K} \\
0 & M_{a'}^{a} & \dots & 0 \\
\vdots & \vdots & \vdots & \vdots \\
0 & 0 & \dots & M_{a'}^{a}
\end{bmatrix} \begin{bmatrix}
\partial_{i'} \\
\partial_{1}^{a'} \\
\vdots \\
\partial_{K}^{a'}
\end{bmatrix}$$

$$(1.5) \qquad \begin{bmatrix} \partial_{i'} \\ \partial_{1}^{a'} \\ \vdots \\ \partial_{K}^{a'} \end{bmatrix} = \begin{bmatrix} \frac{\partial x^{j}}{\partial x^{i'}} & (\partial_{i'} & M_{b}^{b'}) p_{b'}^{1} & \dots & (\partial_{i'} & M_{b}^{b'}) p_{b'}^{K} \\ 0 & M_{b}^{a'} & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & M_{a'}^{a'} \end{bmatrix} \begin{bmatrix} \partial_{j} \\ \partial_{1}^{b} \\ \vdots \\ \partial_{K}^{b} \end{bmatrix}.$$

Substituting (1.5) into (1.4), we obtain

(1.6)
$$(a) \frac{\partial x^{i'}}{\partial x^{i}} \frac{\partial x^{j}}{\partial x^{i'}} = \delta_{i}^{j} \qquad (b) \stackrel{(\alpha)}{M_{a'}^{a}} \stackrel{(\alpha)}{M_{b}^{a'}} = \delta_{b}^{a}$$

$$(c) \frac{\partial x^{i'}}{\partial x^{i}} \frac{\partial x^{j'}}{\partial x^{i'}} + (\partial_{i} \stackrel{(\alpha)}{M_{a'}^{a}}) p_{a}^{\alpha} \stackrel{(\alpha)}{M_{b'}^{a'}} = 0.$$

(1.6c) is the consequence of (1.4) and (1.6b).

2. Adapted bases in $T(E^*)$ and $T^*(E^*)$

From (1.4) and (1.5) it is obvious that ∂_i and $\partial_{i'}$ are not transformed as tensors, so we introduce a new, the so-called adapted basis $B = \{(\delta_i), (\partial_1^a), \ldots, (\partial_K^a)\}$ of $T(E^*)$, where by definition

(2.1)
$$\delta_i = \partial_i - N_{ai}^{\alpha}(x, p) \partial_{\alpha}^a$$

and $N_{ai}^{\alpha}(x,p)$ are the coefficients of the nonlinear connection. Under coordinate transformation (1.1) and (1.2), they transform in the following way:

$$(2.2) \qquad (a) \quad N_{a'i'}^{\alpha}(x',p') = \stackrel{(\alpha)}{M_{a'}^{a}} \frac{\partial x^{i}}{\partial x^{i'}} N_{ai}^{\alpha}(x,p) - p_{a}^{\alpha} \frac{\partial \stackrel{(\alpha)}{M_{a'}^{a}}}{\partial x^{i}} \frac{\partial x^{i}}{\partial x^{i'}}$$

$$(b) \quad N_{bj}^{\alpha}(x,p) = \stackrel{(\alpha)}{M_{b}^{b'}} \frac{\partial x^{j'}}{\partial x^{j}} N_{b'j'}^{\alpha}(x',p') + p_{a}^{\alpha} \stackrel{(\alpha)}{M_{b}^{a'}} \frac{\partial \stackrel{(\alpha)}{M_{a'}^{a'}}}{\partial x^{j'}}$$

Any vector field $X \in T(E^*)$ in the adapted basis B is given by

$$(2.3) X = X^i \delta_i + X_a^{\alpha} \partial_{\alpha}^a.$$

The coordinates of the vector X given by (2.3) and the elements of basis B are transformed as tensor in the following way:

(2.4)
$$(a) \delta_{i} = \frac{\partial x^{i'}}{\partial x^{i}} \delta_{i'} \qquad (b) \partial_{\alpha}^{a} = \stackrel{(\alpha)}{M_{a'}^{a}} (x) \partial_{\alpha}^{a'}$$

$$(c) X^{i} = \frac{\partial x^{i}}{\partial x^{i'}} X^{i'} \qquad (d) X_{a}^{\alpha} = \stackrel{(\alpha)}{M_{a'}^{a'}} (x') X_{a'}^{\alpha}.$$

From (1.3) and (2.3) we obtain the relation between the coordinates of the field X in the bases \bar{B} and B. They are connected by the formula

$$X^i = \bar{X}^i, \quad X^{\alpha}_a = \bar{X}^{\alpha}_a + N^{\alpha}_{ai} \bar{X}^i$$

The subspace of $T(E^*)$ spanned by $\{\delta_i\}$ shall be denoted by $T_H(E^*)$ (horizontal part) and the subspace spanned by $\{\partial_{\alpha}^a\}$ by $_{(\alpha)}T_V(E^*)$ (the vertical α part). So we have

$$T(E^*) = T_H(E^*) \oplus T_V(E^*),$$

where

$$T_V(E^*) = \sum_{\alpha=1}^K \alpha_{\alpha} T_V(E^*),$$

$$\dim T_H(E^*) = n, \quad \dim_{(\alpha)} T_V(E^*) = m.$$

 $X^i\delta_i$ is the horizontal and $X^{\alpha}_a\partial^a_{\alpha}$ the vertical part of the field X.

Now (2.3) may be written in the form

$$X = X_H + X_V, \quad X_H = X^i \delta_i, \quad X_V = X_a^{\alpha} \partial_{\alpha}^a.$$

Let us consider the dual tangent space of E^* , the space $T^*(E^*)$. The natural basis in $T^*(E^*)$ is

$$\begin{array}{lcl} \overline{B}^* & = & \{dx^1, \dots, dx^n, dp_1^1, \dots, dp_m^1, \dots, dp_1^K, \dots dp_m^K\} \\ & = & \{dx^i, dp_a^1, \dots, dp_a^K\}. \end{array}$$

From (1.1) we obtain

(2.5) (a)
$$dx^{i'} = \frac{\partial x^{i'}}{\partial x^i} dx^i$$

(b)
$$dp_{a'}^{lpha}=rac{\partial\stackrel{(lpha)}{M_{a'}^{lpha}}(x)}{\partial x^i}p_a^{lpha}dx^i+\stackrel{(lpha)}{M_{a'}^{lpha}}(x)dp_a^{lpha}.$$

From (2.5b) it is obvious that dp_a^{α} are not transformed as tensors, so we introduce a new basis $B^* = \{(dx^1), (\delta p_a^1), \dots, (\delta p_a^K)\}$, where

(2.6)
$$\delta p_a^{\alpha} = dp_a^{\alpha} + N_{ai}^{\alpha}(x, p)dx^i.$$

By the coordinate transformation (1.1) the bases \overline{B}^* and B^* are related by (2.5a) and

$$(2.7) \quad \text{(a)} \quad \delta p_a^{\alpha} = \stackrel{(\alpha)^{a'}}{M_a} (x') \delta p_{a'}^{\alpha} \qquad \qquad \text{(b)} \quad \delta p_{a'}^{\alpha} = \stackrel{(\alpha)^a}{M_{a'}} (x) \delta p_a^{\alpha}.$$

The proof of (2.7) is obtained using (2.6) and (2.2). Any field $w \in T^*(E^*)$ can be written in the bases and B^* and \overline{B}^* in the following way

$$(2.8) w = \bar{w}_i dx^i + \bar{w}^a_\alpha dp^\alpha_a = w_i dx^i + w^a_\alpha \delta p^\alpha_a,$$

where

$$(2.9) w_i = \bar{w}_i - N_{a_i}^{\alpha} \bar{w}_{\alpha}^a, \quad \bar{w}_{\alpha}^a = w_{\alpha}^a.$$

The subspace of $T^*(E^*)$ spanned by $\{(dx^i)\}$ shall be denoted by $T^*_H(E^*)$ and the subspace spanned by $\{(\delta p_a^{\alpha})\}$ by ${}_{\alpha}T^*_V(E^*)$.

So we have

$$T^*(E^*) = T_H^*(E^*) \oplus T_V^*(E^*),$$

where

$$T_V^*(E^*) = \sum_{\alpha=1}^K {}_{\alpha} T_V^*(E^*).$$

Now (2.8) may be written in the form

$$w=w_H+w_V, \ w_H=w_i dx^i, \ w_V=w^a_{\alpha}\delta p^{\alpha}_a.$$

If $\{(dx^i), (\delta p_a^1), ..., (\delta p_{a'}^K)\}$ and $\{(dx^{i'}), (\delta p_{a'}^1), ..., (\delta p_{a'}^K)\}$ are two bases in $T^*(E^*)$ related by (2.5a) and (2.7) then any $w \in T^*(E^*)$ satisfies the relation

$$(2.10) w = w_i dx^i + w_\alpha^a \delta p_a^\alpha = w_{i'} dx^{i'} + w_\alpha^{a'} \delta p_{a'}^\alpha.$$

Substitutingg $dx^{i'}$ from (2.5a) and $\delta p_{a'}^{\alpha}$ from (2.7b) into (2.10) and commparing the coefficients besides basis vectors, we obtain

(2.11)
$$w_i = w_{i'} \frac{\partial x^{i'}}{\partial x^i}, \quad w_\alpha^a = \stackrel{(\alpha)^a}{M_{a'}} w_\alpha^{a'}.$$

By a straightforward calculation we can prove

Proposition 2.1. The adapted bases $\{(\delta_i), (\partial_1^a), ..., (\partial_K^a)\}$ and $\{(dx^i), (\delta p_a^1), ..., (\delta p_a^K)\}$ are dual to each other, i.e.

$$<\delta_i, dx^j>=\delta_i^j <\delta_i, \delta p_a^{\alpha}>=0$$

 $<\partial_{\alpha}^a, dx^j>=0 <\partial_{\alpha}^a, \delta p_b^{\beta}>=\delta_b^a\delta_{\alpha}^{\beta}.$

3. Tensor Fields on E^*

a) A horizontal tensor fiels t_H has the local representation:

$$t_H = t^{i_1...i_p}_{j_1...j_q}(x,p)\delta_{i_1}\otimes...\otimes\delta_{i_p}\otimes dx^{j_1}\otimes...\otimes dx^{j_q}.$$

It is defined on

$$\underbrace{T_H(E^*) \otimes ... \otimes T_H(E^*)}_{p \ times} \otimes \underbrace{T_H^*(E^*) \otimes ... \otimes T_H^*(E^*)}_{q \ times}.$$

By changing the coordinates given by (1.1) and (1.2), the coordinates of the field t_H have the following transformation law

$$t^{i'_{1}...i'_{p}}_{j'_{1}...j'_{q}} = \frac{\partial x^{i'_{1}}}{\partial x^{i_{1}}}...\frac{\partial x^{i'_{p}}\partial x^{j_{1}}}{\partial x^{i_{p}}\partial x^{j'_{1}}}...\frac{\partial x^{j_{q}}}{\partial x^{j'_{q}}}, \ t^{i_{1}...i_{p}}_{j_{1}...j_{q}}.$$

b) The α vertical tensor field $(\alpha)^{t_V}$ has the local representation

$$_{(\alpha)}t_{V}=_{(\alpha)}t_{b_{1}...b_{s}\alpha...\alpha}^{\alpha...\alpha a_{1}...a_{r}}\partial_{\alpha}^{b_{1}}\otimes...\partial_{\alpha}^{b_{s}}\otimes\delta p_{a_{1}}^{\alpha}\otimes...\otimes\delta p_{a_{r}}^{\alpha}$$

(not summing over α).

 $(\alpha)t_V$ is defined on

$$\underbrace{(\alpha)T_V(E^*)\otimes ...\otimes_{(\alpha)}T_V(E^*)}_{s\ times}\otimes \underbrace{(\alpha)T_V^*(E^*)\otimes ...\otimes_{(\alpha)}T_V^*(E^*)}_{r\ times}.$$

By changing the coordinates of type (1.1) and (1.2) the coordinates of the field $_{(\alpha)}t_V$ given above have the following transformation law

$$(_{\alpha})t_{\alpha...\alpha b'_{1}...b'_{s}}^{a'_{1}...a'_{r}\alpha...\alpha} =$$

$$(_{\alpha})t_{\alpha...\alpha b_{1}...b_{s}}^{a_{1}...a_{r}\alpha...\alpha} (_{\alpha})t_{a'_{1}}^{a_{1}...a_{r}\alpha...\alpha} M_{a'_{1}}^{a'_{1}}...M_{a'_{r}}^{a'_{r}} M_{b'_{1}}^{b_{1}}...M_{b'_{s}}^{b_{s}}.$$

c) A vertical tensor field t_V on $T_V(E^*) \otimes T_V^*(E^*)$ has the form

$$t_V = t_{a_1\beta}^{\alpha b_1} \partial_{\alpha}^{a_1} \otimes \delta p_{b_1}^{\beta}$$

(summing over α and β).

The coordinate transformation of tensor t_V is given by

$$t_{a_{1}'\beta}^{\alpha\ b_{1}'}=t_{a_{1}\beta}^{\alpha\ b_{1}}\ M_{a_{1}'}^{(\alpha)}M_{b_{1}}^{(\beta)}\ .$$

d) A tensor field t on

$$\underbrace{T_{H}(E^{*}) \otimes ... \otimes T_{H}(E^{*})}_{p \ times} \otimes \underbrace{T_{H}^{*}(E^{*}) \otimes ... \otimes T_{H}^{*}(E^{*})}_{q \ times} \otimes \underbrace{T_{V}(E^{*}) \otimes ... \otimes T_{V}(E^{*})}_{q \ times} \otimes \underbrace{T_{V}^{*}(E^{*}) \otimes ... \otimes T_{V}^{*}(E^{*})}_{r \ times}$$

is given by

$$\begin{split} t &= t^{i_1 \dots i_p} \underset{j_1 \dots j_q}{\overset{\beta_1 \dots \beta_s a_1 \dots a_r}{b_1 \dots b_s \alpha_1 \dots \alpha_r}} (x,p) \delta_{i_1} \otimes \dots \otimes \delta_{i_p} \otimes \\ dx^{j_1} &\otimes \dots \otimes dx^{j_q} \otimes \partial_{\beta_1}^{b_1} \otimes \dots \otimes \partial_{\beta_s}^{b_s} \otimes \delta p_{a_1}^{\alpha_1} \otimes \dots \otimes \delta p_{a_r}^{\alpha_r}. \end{split}$$

The summation goes over all the indices.

The coordinate transformation of the above tensor is given by

$$\begin{split} t &= t^{i'_{1}...i'_{p}} \int_{j'_{1}...j'_{q}}^{\beta_{1}...\beta_{s}a'_{1}...a'_{r}} = t^{i_{1}...i_{p}} \int_{j_{1}...j_{q}}^{\beta_{1}...\beta_{s}a_{1}...a_{r}} \int_{j_{1}...j_{q}}^{\beta_{1}...\beta_{s}a_{1}...a_{r}} \\ &\frac{\partial x^{i'_{1}}}{\partial x^{i'_{1}}}...\frac{\partial x^{i'_{p}}\partial x^{j_{1}}}{\partial x^{j_{1}}}...\frac{\partial x^{j_{q}}}{\partial x^{j'_{q}}} \stackrel{(\beta_{1})}{M^{b_{1}}_{b'_{1}}}... \stackrel{(\beta_{s})}{M^{b_{s}}_{a'_{1}}} \stackrel{(\alpha_{1})}{\dots} \stackrel{(\alpha_{r})}{M^{a'_{r}}_{a'_{r}}} \end{split}$$

The order of spaces $T_H(E^*)$, $T_H^*(E)$, $T_V(E^*)$ and $T_V^*(E^*)$ can be taken arbitrary. It has an influence on the order of indices of tensor t which is defined on their tensor product.

4. Metric tensor in the K-Hamilton space

In the space $T^*(E^*)\otimes T^*(E^*)$ the metric tensor G with respect to the basis $\{(dx^i),(\delta p_a^1,...,(\delta p_a^K))\}$ has the form

$$(4.1) G = [(dx^{i})(\delta p_{a}^{1})...(\delta p_{a}^{K})].$$

$$\begin{bmatrix} [g_{ij}] & [g_{i1}^{b}] & \dots & [g_{iK}^{b}] \\ [g_{1j}^{a}] & [g_{11}^{ab}] & \dots & [g_{1K}^{ab}] \\ \vdots & \vdots & \vdots & \vdots \\ [g_{Ki}^{a}] & [g_{K1}^{ab}] & \dots & [g_{KK}^{ab}] \end{bmatrix} \otimes \begin{bmatrix} dx^{j} \\ \delta p_{b}^{1} \\ \vdots \\ \delta p_{b}^{K} \end{bmatrix}.$$

The matrices $[g_{ij}]$, $[g_{i\beta}^{\ b}]$, $[g_{\alpha j}^{\ a}]$ and $[g_{\alpha \beta}^{\ ab}]$ have the format $n \times n$, $n \times m$, $m \times n$ and $m \times m$ respectively. As G is a tensor, its coordinates in the new coordinate system (x', p') are transformed in the following way:

$$(4.2) a) g_{i'j'} = g_{ij} \frac{\partial x^{i}}{\partial x^{i'}} \frac{\partial x^{j}}{\partial x^{j'}} b) g_{\alpha j'}^{a'} = g_{\alpha j}^{a} M_{a}^{a'} \frac{\partial x^{j}}{\partial x^{j'}}$$

$$c) g_{i'\beta}^{b'} = g_{i\beta}^{\ b} \frac{\partial x^{i}}{\partial x^{i'}} M_{b}^{b'} d) g_{\alpha\beta}^{a'b'} = g_{\alpha\beta}^{ab} M_{a}^{a'} M_{b}^{b'}$$

We shall suppose that G is a symmetric, positive definite tensor field of rank n + mK. From the symmetry it follows that

$$g_{ji} = g_{ij}$$
 $g_{i\beta}^{\ \ b} = g_{\beta i}^{\ \ b}$ $g_{\alpha\beta}^{ab} = g_{\beta\alpha}^{ba}$.

The "covariant" coordinates of the field $X = X^i \delta_i + X_a^{\alpha} \partial_{\alpha}^a$ are given by

(4.3)
$$X_i = g_{ij} X^j + g^a_{i\alpha} X^a_{\alpha}, \quad X^a_{\alpha} = g^a_{\alpha i} X^i + g^{ab}_{\alpha \beta} X^{\beta}_b.$$

The inverse matrix of G (appearing in (4.1)) is given by

$$\left[egin{array}{cccc} [g^{jk}] & [g^{jl}_{c}] & \dots & [g^{jK}_{c}] \ [g^{1k}_{b}] & [g^{11}_{bc}] & \dots & [g^{1K}_{bc}] \ dots & dots & dots & dots \ [g^{Kk}_{bc}] & [g^{K1}_{bc}] & \dots & [g^{KK}_{bc}] \end{array}
ight]$$

The matrices $[g^{jk}]$, $[g^{j\gamma}_c]$, $[g^{\gamma k}_c]$ and $[g^{\beta\gamma}_{bc}]$ have the format $n \times n$, $n \times m$, $m \times m$ and $m \times m$ respectively. Now we have

(4.4) a)
$$g_{ij}g^{jk} + g_{i\beta}^{\ b}g_b^{\beta k} = \delta_i^k$$
 b) $a_{\alpha j}^a g^{j\gamma}_c + g_{\alpha\beta}^{ab}g_{bc}^{\beta\gamma} = \delta_c^a \delta_{\alpha}^{\gamma}$
c) $g_{ij}g^{j\gamma}_c + g_{i\beta}^{\ b}g_{bc}^{\beta\gamma} = 0$ d) $g_{\alpha j}^a g^{jk} + g_{\alpha\beta}^{ab}g_b^{\beta k} = 0$.

The contravariant coordinates of $w = w_i dx^i + w^a_\alpha \delta p^\alpha_a$ are given by

$$(4.5) w^i = g^{ij}w_j + g^{i\alpha}_{\ a}w^a_{\alpha} w^{\alpha}_{a} = g^{j\alpha}_{\ a}w_j + g^{\alpha\beta}_{ab}w^b_{\beta}.$$

Using (2.4), (2.11) and (4.2) it can be shown that following transformation laws are valid:

$$X_{i'} = X_i \frac{\partial x^i}{\partial x^{i'}} \quad X_{\alpha}^{a'} = X_{\alpha}^a \stackrel{(\alpha)}{M_a^{a'}}$$

$$w_{i'} = w_i rac{\partial x^i}{\partial x^{i'}} \quad w_{a'}^{lpha} = w_a^{lpha} \stackrel{(lpha)}{M_{a'}^a}.$$

If the k-Hamilton function H(X, p) is given in the space E^* , then the metric tensor G can be defined in the following way:

$$g_{ij}(x,p) = g_{ij}(x)$$
 $g_{i\beta}^{\ b} = 0$ $g_{\alpha j}^{\ a} = 0$

$$g_{\alpha\beta}^{ab} = 2^{-1} \partial_{\alpha}^{a} \partial_{\beta}^{b} H^{2}(x, p) \quad \forall \ \alpha, \beta = \overline{1, K},$$

where $g_{ij}(x)$ is some metric tensor defined on M and M is the π^* projection of E^*

$$\pi^*(E^*) = M, \quad \pi^*((x^i), (p_a^1), ..., (p_a^K)) = (x^i).$$

We can not define

$$g_{i\beta}{}^b(x,p) = 2^{-1}\delta_i\partial_\beta^b H^2(x,p), \quad g_{ij}(x,p) = 2^{-1}\delta_i\delta_j H^2(x,p),$$

because the above quantities are not transformed as tensor.

Using the metric G determined by (4.1) we define the scalar product (X,Y) of fields $X,Y\in T(E^*)$ by

$$(4.6) (X,Y) = g_{ij}X^{i}Y^{j} + g_{i\beta}^{\ b}X^{i}Y_{b}^{\beta} + g_{\alpha j}^{a}X_{\alpha}^{\alpha}Y^{j} + g_{\alpha\beta}^{ab}X_{\alpha}^{\alpha}Y_{b}^{\beta}.$$

Then length of X, |X| is defined by $|X|^2 = (X, X)$ and $\cos \theta$, where θ is the angle between X and Y by

(4.7)
$$\cos \theta = \frac{(X,Y)}{|X||Y|}.$$

When $\cos \theta = 0$, we say that the fields X and Y are orthogonal to each other. For the horizontal field X_H we have

$$X_H = X^i \delta_i, |X_H|^2 = g_{ij} X^i X^j$$

and for vertical vector X_V we have

$$X_V = X_a^{\alpha} \partial_{\alpha}^a |X_V|^2 = g_{\alpha\beta}^{ab} X_a^{\alpha} X_b^{\beta}.$$

For the field $(\alpha)X_V \in (\alpha) T_V(E^*)$ we have

$$_{(\alpha)}X_V = X_a^{\alpha}\partial_{\alpha}^a, \ |_{(\alpha)}X_V|^2 = g_{\alpha\alpha}^{ab}X_a^{\alpha}X_b^{\alpha} (\text{not summing over } \alpha).$$

Theorem 4.1. The necessary and sufficient conditions that the subspaces $T_H(E^*)$, $T_{V}(E^*)$, $T_{V}(E^*)$ of $T(E^*)$ should be orthogonal to each other with respect to the metric tensor G, are

$$[g_{i\beta}^{\ \ b}] = 0, \ [g_{\alpha i}^{\ a}] = 0, \ \forall \ \alpha, \beta = \overline{1, K} \ \alpha \neq \beta.$$

Definition 4.1. The differentiable manifold E^* in which the coordinate transformations of type (1.1) and (1.2) are allowed, supplied with the nonlinear connection N (see (2.2)) and the metric tensor G (given by (4.1)) is called the K- Hamilton space.

5. Strongly distinguished connection in $T(E^*)$

The distinguished connection ∇ or d— connection in the K— Hamilton space in [15], [9], [16] and others is defined as a function $\nabla:(X,Y)\to \nabla_X Y; X,Y, \nabla_X Y\in T(E^*)$ for which, besides the usual conditions for the linear connection, the following restrictions hold

(5.1) (a)
$$\nabla_X Y_H \in T_H(E^*)$$
 (b) $\nabla_X Y_V \in T_V(E^*)$
$$\forall X \in T(E^*), \forall Y_H \in T_H(E^*) \text{ and } \forall Y_V \in T_V(E^*).$$

The strongly distinguished or s.d. - connection is the linear connection for which (5.1 a) and (5.2)

(5.2)
$$\nabla_{X(\alpha)} Y_V \in {}_{(\alpha)} T_V(E^*)$$

hold.

Definition 5.1. The strongly distinguished connection ∇ in $T(E^*)$ is the linear connection defined by

(5.3)
$$(a) \quad \nabla_{\delta_{i}} \delta_{j} = F_{j i}^{k} \delta_{k}, \quad b) \quad \nabla_{\delta_{i}} \partial_{\alpha}^{a} = F_{c i}^{(\alpha)} \partial_{\alpha}^{c}$$

$$(c) \quad \nabla_{\partial_{\alpha}^{a}} \delta_{j} = C_{j \alpha}^{k a} \delta_{k} \quad d) \quad \nabla_{\partial_{\alpha}^{a}} \partial_{\beta}^{b} = C_{c \alpha}^{b a} \partial_{\beta}^{c}.$$

In (5.3b) and (5.3d) there is no summation over α and β respectively.

Proposition 5.1. If $X, Y \in T(E^*)$, where X is given by (2.3) and $Y = Y^j \delta_j + Y_b^{\beta} \partial_{\beta}^b$, then

(5.4)
$$\nabla_X Y = (Y_{li}^j X^i + Y^j |_{\alpha}^a X_a^{\alpha}) \delta_j + (Y_{bli}^{\beta} X^i + Y_b^{\beta} |_{\alpha}^a X_a^{\alpha}) \partial_{\beta}^b,$$

where

$$(a) Y_{\mathbf{i}i}^{j} = \delta_{i} Y^{j} + F_{k i}^{j} Y^{k}$$

$$(b) Y^{j}|_{\alpha}^{a} = \partial_{\alpha}^{a} Y^{j} + C_{k\alpha}^{ja} Y^{k}$$

$$(5.5) (c) Y_{b|i}^{\beta} = \delta_i Y_b^{\beta} + F_{bi}^d Y_d^{\beta}$$

$$(d) \qquad Y_b^\beta|_\alpha^a = \partial_\alpha^a Y_b^\beta + C^{(\beta)}_{b\alpha}{}^a Y_d^\beta.$$

Proposition 5.2. If $((x^i),(p_a^{\alpha}))$ and $((x^{i'}),(p_{a'}^{\alpha}))$ are two coordinate systems connected by (1.1) and (1.2) then

$$(5.6) \nabla_{X'}Y' = \nabla_X Y$$

iff $Y^j_{li},~Y^j|^a_{\alpha},Y^{\beta}_{b|i}$ and $Y^{\beta}_b|^a_{\alpha}$ are transformed as tensors, i.e.

$$Y_{|i'}^{j'} = Y_{|i}^{j} (\partial_{i'} x^{i})(\partial_{j} x^{j'}) \qquad Y_{|\alpha}^{j'}|_{\alpha}^{\alpha'} = Y_{|\alpha}^{j}|_{\alpha}^{\alpha}(\partial_{j} x^{j'}) \stackrel{(\alpha)}{M_{\alpha}^{\alpha'}}$$

$$Y_{b'li'}^{\beta} = Y_{bli}^{\beta} \stackrel{(\beta)}{M_{b'}^{b}} (\partial_{i'} x^{i}) \qquad Y_{b'l\alpha}^{\beta|\alpha'} = Y_{b}^{\beta|\alpha} \stackrel{(\beta)}{M_{b'}^{b}} M_{a'}^{\alpha'}$$

or equivalently iff the s.d.-connection coefficients have the following law of transformation:

(a)
$$F_{ji}^{k} = F_{j'i'}^{k'}(\partial_{j}x^{j'})(\partial_{k'}x^{k})(\partial_{i}x^{i'}) + (\partial_{i}\partial_{j}x^{k'})(\partial_{k'}x^{k})$$
(b)
$$F_{ci}^{b} = F_{c'i'}^{b'}M_{b'}^{b}M_{c}^{c'}(\partial_{i}x^{i'}) + (\partial_{i}M_{b'}^{b})M_{c}^{d'}$$
(5.7)

$$(c) \quad C_{ja}^{k\alpha} = C_{j'a'}^{k'\alpha'}(\partial_j x^{j'})(\partial_{k'} x^k) \stackrel{(\alpha)}{M_{a'}^a}$$

$$(d) \quad {\overset{(\beta)}{C}}_{c\alpha}^{b\ a} = {\overset{(\beta)}{C}}_{b'}^{b'\ a'} = {\overset{(\beta)}{M}}_{b'}^{b'} M_c^{c'} {\overset{(\alpha)}{M}}_{a'}^a \; .$$

Proof. The proof is obtained by direct calculation using (5.4) - (5.6), (2.1) and (2.4).

It follows from (5.7) that the F's which appear in (5.7a) and (5.7b) are transformed as connection coefficients, and the two C's in (5.7c) and (5.7b) as tensors.

The torsion tensor T(X,Y) for the s.d.-connection, is as usual, given by

$$(5.8) T(X,Y) = \nabla_X Y - \nabla_Y X - [X,Y]. \quad \Box$$

Theorem 5.1. In the k-Hamilton space the torsion tensor for the s.d. -connection has the form

$$(5.9) T(X,Y) = (F_{ki}^{j} - F_{ik}^{j})X^{i}Y^{k}\delta_{j} + C_{k\alpha}^{ja}X^{\alpha}Y^{k}\delta_{j} - C_{i\beta}^{jb}X^{i}Y_{b}^{\beta}\delta_{j} + C_{k\alpha}^{(\beta)^{da}}X^{\alpha}X_{a}^{\alpha}Y_{b}^{\beta}\partial_{\alpha}^{c} + C_{i\alpha}^{(\beta)^{da}}X_{a}^{\alpha}Y_{b}^{\beta}\partial_{\beta}^{b}X_{ai}^{\alpha} - N_{bi}^{\beta}\partial_{\beta}^{b}N_{aj}^{\alpha})X^{i}Y^{j}\partial_{\alpha}^{a} + (\partial_{i}N_{aj}^{\alpha} - \partial_{j}N_{ai}^{\alpha} + N_{bj}^{\beta}\partial_{\beta}^{b}N_{ai}^{\alpha} - N_{bi}^{\beta}\partial_{\beta}^{b}N_{aj}^{\alpha})X^{i}Y^{j}\partial_{\alpha}^{a} + (G_{i\alpha}^{\beta})Y_{bi}^{\beta}Y_{a}^{\alpha} - (G_{i\alpha}^{\alpha}N_{bi}^{\beta})Y_{a}^{\alpha}X_{a}^{\beta}Y^{j}\partial_{\beta}^{b} + ((\partial_{\alpha}^{a}N_{bj}^{\beta})X_{a}^{\alpha} - F_{bj}^{d}X_{a}^{\beta}Y^{j}\partial_{\beta}^{b}.$$

Proof. The proof is obtained by direct calculation using (5.3), (5.4), (5.5) and (5.8).

If we suppose that the nonlinear connection N is such that N_{bi}^{β} is the function only of (x) and $(p^{\beta}) = (p_1^{\beta}, ..., p_m^{\beta})$ for $\forall \beta = \overline{1, K}$, then the three

last lines in (5.9) have the form

$$(\partial_i N^\alpha_{aj} - \partial_j N^\alpha_{ai} + N^{(\alpha)}_{bj} \partial^b_{(\alpha)} N^\alpha_{bi} - N^{(\alpha)}_{b} \partial^b_{(\alpha)} N^\alpha_{aj}) X^i Y^j \partial^a_\alpha +$$

$$(5.10) + (F^{d}_{bi} - \partial^{d}_{(\beta)} N^{(\beta)}_{bi}) X^{i} Y^{\beta}_{d} \partial^{b}_{\beta} + (\partial^{d}_{(\beta)} N^{(\beta)}_{bj} - F^{d}_{bj}) X^{\beta}_{d} Y^{j} \partial^{b}_{\beta}.$$

The Einstein summation convection is not meant on (α) and (β) .

Theorem 5.2. In the K- Hamilton space the torsion tensor for the s.d. - connection is identically equal to zero iff

(a)
$$\partial_{\alpha}^{a} N_{\beta i}^{b} = 0$$
 for $\forall \alpha \neq \beta$

(b)
$$F_{ki}^j - F_{ik}^j = 0$$
,

(c)
$$C_k^{j\alpha} = 0$$

(5.11)

(d)
$$C^{(\alpha)}_{b\alpha}^{(\alpha)} - C^{(\alpha)}_{b\alpha}^{(\alpha)} = 0$$

(e)
$$C_{b\gamma}^{(\alpha)} = 0$$
 for $\forall \alpha \neq \gamma$

$$(f) \quad \partial^{i} N_{aj}^{\alpha} - \partial_{j} N_{ai}^{\alpha} - N_{bi}^{(\alpha)} \partial_{(\alpha)}^{b} N_{aj}^{\alpha} + N_{bj}^{(\alpha)} \partial_{(\alpha)}^{b} N_{ai}^{\alpha} = 0$$

$$(g) \quad F^d_{bi} = \partial^d_{(\beta)} N^\beta_{bi}.$$

Proof. The proof follows from (5.9) and (5.10).

Theorem 5.3. In the K-Hamilton space the torsion free s.d. - connection has the following properties

(a)
$$\nabla_{\delta_i}\delta_j = \nabla_{\delta_i}\delta_i$$
, $\nabla_{X_H}Y_H \in T_H(E^*)$

(b)
$$\nabla_{(\alpha)X_V}(\beta)Y_V = 0 \ \forall \alpha \neq \beta$$

(5.12)

(c)
$$\nabla_{\partial_{\alpha}^{a}} \partial_{\alpha}^{b} = \nabla_{\partial_{\alpha}^{b}} \partial_{\alpha}^{a} \qquad \nabla_{(\alpha)X_{V}(\alpha)} Y_{V} \in {}_{(\alpha)} T_{V}(E^{*})$$

(d)
$$\nabla_{(\alpha)} X_{\prime\prime} Y_H = 0$$
.

If instead of (5.11 a) the stronger condition

(5.13)
$$\partial_{\alpha}^{a} N_{\beta i}^{b} = 0 \qquad \forall \alpha \neq \beta = \overline{1, m}$$

is satisfied, then besides (5.12) the torsion free s.d. - connection has the property $\nabla_{X_H}(\alpha)Y_V=0$.

From Theorems 5.2 and 5.3 follows the following

Theorem 5.4. In the torsion free K- Hamilton space for which condition (5.13) holds the strongly distinguished connection reduces to the following form:

$$abla_{\delta_i}\delta_j=F_{ji}^k\delta_k,\ F_{ij}^k=F_{ji}^k,\
abla_{\delta_i}\partial_{lpha}^a=0,\
abla_{\partial_{oldsymbol{lpha}}^a}\delta_j=0,$$

(5.14)
$$\nabla_{\partial_{\alpha}^{a}} \partial_{\beta}^{b} = 0 \quad \text{for } \alpha \neq \beta, \quad \nabla_{\partial_{\alpha}^{a}} \partial_{\alpha}^{b} = \overset{(\alpha)}{C^{b}}_{c}^{a} \partial_{\alpha}^{c}, \quad \overset{(\alpha)}{C^{b}}_{c}^{a} = \overset{(\alpha)}{C^{a}}_{c}^{b},$$

6. Strongly distinguished connection in $T^*(E^*)$

The connection ∇ defined on $T(E^*)$ by (5.3) induces a connection ∇^* on $T^*(E^*)$ which will also be denoted by ∇ . For the field X defined by (2.9) and w defined by (2.14) we have

(6.1)
$$\nabla_X w = \nabla_{X^i \delta_i + X^{\alpha}_{\alpha} \partial_{\alpha}^{\alpha}} (w_j dx^j + w^b_{\beta} \delta p^{\beta}_b).$$

Definition 6.1. The connection ∇ on $T^*(E^*)$ is defined by

(a)
$$\nabla_{\delta_i} dx^j = \tilde{F}_k^{\ j}_{i} dx^k$$
 (b) $\nabla_{\delta_i} \delta p_b^{\beta} = \tilde{F}^c_{\ bi} \delta p_c^{\beta}$ (6.2)

$$(c) \ \nabla_{\partial_{\alpha}^{a}} dx^{j} = \tilde{C}_{k}^{\ j}{}^{a}_{\alpha} dx^{k} \ (d) \nabla_{\partial_{\alpha}^{a}} \delta p_{b}^{\beta} = \tilde{C}^{c}{}^{a}_{b\alpha} \ \delta p_{c}^{\beta}.$$

Proposition 6.1. If $X \in T(E^*)$ and $w \in T^*(E^*)$, then

(6.3)
$$\nabla_X w = (w_{j|i} X^i + w_j|_\alpha^a X_a^\alpha) dx^j + (w_{\beta|i}^b X^i + w_{\beta}^b|_\alpha^a w_a^\alpha) \delta p_b^\beta,$$

where

(a)
$$w_{j|i} = \delta_i w_j + \tilde{F}_j^{\ k} w_k$$

(b)
$$w_j|_{\alpha}^a = \partial_{\alpha}^a w_j + \tilde{C}_{j\alpha}^{ka} w_k$$

(6.4)

$$(c) \ w_{\beta|i}^b = \delta_i w_{\beta}^b + \tilde{F}^b_{\ ci} \ w_{\beta}^c$$

$$(d) \ w_{\beta}^b|_{\alpha}^a = \partial_{\alpha}^a w_{\beta}^b + \tilde{C}^b{}_{c\alpha}^a w_{\beta}^c.$$

Proof. Substituting (6.2) into (6.1) and using the linarity of connection ∇ we obtain (6.3) and (6.4). \square

Using the properties of the linear connection ∇ and relations (4.1) and (6.2), the following relations are fulfilled:

$$\nabla_X G = (g_{ij|k} X^k + g_{ij} \mid_{\gamma}^c X_c^{\gamma}) dx^i \otimes dx^j +$$

$$(g_{i\beta|k}^b X^k + g_{i\beta}^b \mid_{\gamma}^c X_c^{\gamma}) dx^i \otimes \delta p_b^{\beta} +$$

$$(g_{\alpha j|k}^a X^k + g_{\alpha j}^a \mid_{\gamma}^c X_c^{\gamma}) \delta p_a^{\alpha} \otimes dx^j +$$

$$(g_{\alpha \beta|k}^{ab} X^k + g_{\alpha \beta}^{ab} \mid_{\gamma}^c X_c^{\gamma}) \delta p_a^{\alpha} \otimes \delta p_b^{\beta},$$

where

(a)
$$g_{ij|k} = \delta_k g_{ij} + g_{hj} \tilde{F}_{ik}^{h} + g_{ih} \tilde{F}_{ik}^{h}$$

(b)
$$g_{ij}|_{\gamma}^{c} = \partial_{\gamma}^{c} g_{ij} + g_{hj} \tilde{C}_{i}|_{\gamma}^{hc} + g_{ih} \tilde{C}_{j}|_{\gamma}^{hc}$$

(c)
$$g_{i\beta|k}^{\ b} = \delta_k g_{i\beta}^{\ b} + g_{h\beta}^{\ b} \tilde{F}_{ik}^h + g_{i\beta}^{\ d} \, \tilde{F}_{\ dk}^b$$

(6.6)

(d)
$$g_{i\beta}^{\ b}|_{\gamma}^{c} = \partial_{\gamma}^{c} g_{i\beta}^{\ b} + g_{h\beta}^{\ b} \tilde{C}_{i\gamma}^{hc} + g_{i\beta}^{\ d} \tilde{C}_{\ d\gamma}^{\ d\gamma}$$

(e)
$$g_{\alpha\beta|k}^{ab} = \delta_k g_{\alpha\beta}^{ab} + g_{\alpha\beta}^{db} \tilde{F}_{dk}^a + g_{\alpha\beta}^{ad} \tilde{F}_{dk}^b$$

$$(\mathrm{f}) \ g^{ab}_{\alpha\beta}|^{c}_{\gamma} = \partial^{c}_{\gamma}g^{ab}_{\alpha\beta} + g^{db}_{\alpha\beta} \, \tilde{C}^{acc}_{d\gamma}^{c} + g^{ab}_{\alpha\beta} \, \tilde{C}^{bcc}_{\phantom{d\gamma}}^{\phantom{(\beta)}} \, .$$

Definition 6.2. The K-Hamilton space will be called recurrent if there exists a field $\lambda(x,p) = \lambda_k(x,p)dx^k + \lambda_{\gamma}^c(x,p)\delta p_c^{\gamma}$, such that

$$g_{ij|k} = \lambda_k g_{ij} \qquad g_{ij}|_{\gamma}^c = \lambda_{\gamma}^c g_{ij}$$

$$g_{i\beta|k}^b = \lambda_k g_{i\beta}^b \qquad g_{i\beta}^b|_{\gamma}^c = \lambda_{\gamma}^c g_{i\beta}^b$$

$$g_{\alpha j|k}^a = \lambda_k g_{\alpha j}^a \qquad g_{\alpha j}^a|_{\gamma}^c = \lambda_{\gamma}^c g_{\alpha j}^a$$

$$g_{\alpha \beta|k}^{ab} \lambda_k g_{\alpha \beta}^{ab} \qquad g_{\alpha \beta}^{ab}|_{\gamma}^c = \lambda_{\gamma}^c g_{\alpha \beta}^{ab}.$$

The K-Hamilton space will be called a metric space if

(6.8)
$$g_{\alpha j|k}^a = 0$$
 $g_{\alpha j}^a|_{\gamma}^c = 0$ $g_{\alpha \beta|k}^{ab} = 0$ $g_{\alpha \beta}^{ab}|_{\gamma}^c = 0$.

Theorem 6.1. In the recurrent and in the metric K-Hamilton spaces the strongly distinguished connection ∇ defined on $T(E^*)$ by (5.3) and on $T^*(E^*)$ by (6.2) is compatible with the raising and lowering of the indices by the metric tensor G iff following conditions are satisfied

(6.9)
$$\tilde{F}_{k}^{\ j}_{i} = -F_{k}^{\ j}_{i} \quad \tilde{F}_{bi}^{c} = -F_{bi}^{(\beta)},$$

$$\tilde{C}_{k}^{\ ja}_{\alpha} = -C_{k}^{\ ja}_{\alpha}, \quad \tilde{C}^{c}_{b\alpha}^{\ a} = -C_{b\alpha}^{(\beta)}.$$

Proof. The proof is obtained from (4.3), (6.6), (6.7), (6.4) and (5.5). It is similar to the proof of Theorem 6.1 in [8].

Theorem 6.2. In the recurrent K-Hamilton space the coordinates of the inverse metric tensor satisfy the following relations

$$(6.10) g_{|k}^{ij} = -\lambda_k g^{ij}, g_{b|k}^{i\beta} = -\lambda_k g_b^{i\beta}, g_{db|k}^{\delta\beta} = -\lambda_k g_{db}^{\delta\beta}$$
$$g^{ij}|_{\gamma}^c = -\lambda_{\gamma}^c g^{ij}, g_b^{i\beta}|_{\gamma}^c = -\lambda_{\gamma}^c g_b^{i\beta}, g_{db}^{\delta\beta}|_{\gamma}^c = -\lambda_{\gamma}^c g_{db}^{\delta\beta}.$$

Proof. The proof follows from (4.4) and is the same as the proof of Theorem 6.2 in [8.]

To avoid confusion in the application of the Einstein summation convention in the case when more than two indices are equal, we introduce the following notation:

$$F_{\alpha dk}^{a\delta} = \left\{ \begin{array}{ll} F_{dk}^{(\alpha)} & \text{for} \quad \alpha = \delta \\ 0 & \text{for} \quad \alpha \neq \delta \end{array} \right.$$

$$C_{\alpha d\gamma}^{a\delta c} = \left\{ \begin{array}{ll} C_{d\gamma}^{(\alpha)} & \text{for} \quad \alpha = \delta \\ 0 & \text{for} \quad \alpha \neq \delta \end{array} \right.$$

Using the above notation the rasing and lowering of the middle index of the connection coefficients are given by the following formulae:

(6.11)

$$\left[\begin{array}{c} F_{ijk} \\ F^a_{i\alpha k} \end{array}\right] \; = \left[\begin{array}{cc} g_{ij} & g^{\;d}_{j\delta} \\ g^a_{\alpha h} & g^{ad}_{\alpha \delta} \end{array}\right] \, \left[\begin{array}{c} F^h_{ik} \\ 0 \end{array}\right] \; \Leftrightarrow \left[\begin{array}{c} F^h_{ik} \\ 0 \end{array}\right] \; = \left[\begin{array}{cc} g^{hj} & g^{h\alpha}_a \\ g^{\delta j}_d & g^{\delta\alpha}_{d\alpha} \end{array}\right] \, \left[\begin{array}{c} F_{ijk} \\ F^a_{i\alpha k} \end{array}\right]$$

$$\left[\begin{array}{c} F^a_{\alpha ik} \\ F^{ab}_{\alpha \beta k} \end{array} \right] \; = \; \left[\begin{array}{c} g_{ih} & g^c_{i\delta} \\ g^b_{\delta h} & g^{bc}_{\beta \gamma} \end{array} \right] \; \left[\begin{array}{c} 0 \\ F^{a\gamma}_{\alpha ck} \end{array} \right] \; \Leftrightarrow \left[\begin{array}{c} 0 \\ F^{a\delta}_{\alpha dk} \end{array} \right] \; = \; \left[\begin{array}{c} g^{hi} & g^{h\beta}_{\ b} \\ g^{\delta i}_{\ d} & g^{\delta\beta}_{\ db} \end{array} \right] \; \left[\begin{array}{c} F^a_{\alpha ik} \\ F^{ab}_{\alpha \beta k} \end{array} \right]$$

$$\left[\begin{array}{c} C_{ij\gamma}^{\ c} \\ C_{i\alpha\gamma}^{\ ac} \end{array} \right] \ = \left[\begin{array}{c} g_{jl} & g_{j\delta}^{\ d} \\ g_{\alpha l}^{a} & g_{\alpha\delta}^{ad} \end{array} \right] \, \left[\begin{array}{c} C_{i\gamma}^{\ lc} \\ 0 \end{array} \right] \ \Leftrightarrow \left[\begin{array}{c} C_{i\gamma}^{\ hc} \\ 0 \end{array} \right] \ = \left[\begin{array}{c} g^{hj} & g^{h\alpha}_{\ a} \\ g^{\delta j}_{\ d} & g^{\delta\alpha}_{\ d\alpha} \end{array} \right] \, \left[\begin{array}{c} C_{ij\gamma}^{\ cc} \\ C_{i\alpha\gamma}^{\ ac} \end{array} \right]$$

$$\begin{bmatrix} \begin{array}{c} C_{\alpha i \gamma}^{a \ c} \\ C_{\alpha i \gamma}^{a b c} \end{array} \end{bmatrix} \ = \ \begin{bmatrix} \begin{array}{c} g_{ih} & g_{i\epsilon}^{\ e} \\ g_{\beta h}^{b} & g_{\beta \epsilon}^{be} \end{array} \end{bmatrix} \ \begin{bmatrix} \begin{array}{c} 0 \\ C_{\alpha e \gamma}^{a c c} \end{array} \end{bmatrix} \ \Leftrightarrow \ \begin{bmatrix} \begin{array}{c} 0 \\ C_{\alpha d \gamma}^{a b c} \end{array} \end{bmatrix} \ = \ \begin{bmatrix} \begin{array}{c} g^{hi} & g_{b}^{h\beta} \\ g_{d}^{\delta i} & g_{bd}^{\delta \beta} \end{array} \end{bmatrix} \ \begin{bmatrix} \begin{array}{c} C_{\alpha i \gamma}^{a \ c} \\ C_{\alpha \beta \gamma}^{a b c} \end{array} \end{bmatrix}$$

The above equations are well defined because they are compatible with (4.4)

Using (6.11), (6.6) may be written in the following form:

(a)
$$g_{ij|k} = \delta_k g_{ij} - F_{ijk} - F_{jik}$$

(b)
$$g_{ij}|_{\gamma}^{c} = \partial_{\gamma}^{c} g_{ij} - C_{ij\gamma}^{c} - C_{ji\gamma}^{c}$$

(6.12)

(c)
$$g_{i\beta|k}^{\ b} = \delta_k g_{i\beta}^{\ b} - F_{i\beta k}^{\ b} - F_{\beta ik}^{\ b}$$

(d)
$$g_{i\beta}^{\ b}|_{\gamma}^{c} = \partial_{\gamma}^{c}g_{i\beta}^{\ b} - C_{i\beta\gamma}^{\ bc} - C_{\beta i\gamma}^{\ bc}$$

(e)
$$g_{\alpha\beta|k}^{ab} = \delta_k g_{\alpha\beta}^{ab} - F_{\alpha\beta k}^{ab} - F_{\beta\alpha k}^{ba}$$

(f)
$$g_{\alpha\beta}^{ab}|_{\gamma}^{c} = \partial_{\gamma}^{c}g_{\alpha\beta}^{ab} - C_{\alpha\beta\gamma}^{abc} - C_{\beta\alpha\gamma}^{bac}$$

7. Strongly distingushed connection coefficients in the recurrent K- Hamilton space

Theorem 7.1. In the recurrent K-Hamilton space supplied with the metric tensor G, arbitrary torsion tensor T, the strongly distinguished connection coefficients are determined by (7.1) - (7.8):

$$(7.1) 2F_{ijk} = (\delta_k g_{ij} + \delta_i g_{jk} - \delta_j g_{ki}) - (\lambda_k g_{ij} + \lambda_i g_{jk} - \lambda_j g_{ki}) + A$$

where

$$A = (F_{ijk} - F_{kji}) + (F_{kij} - F_{jik}) + (F_{ikj} - F_{jki}) =$$

$$= g_j^h (F_{ik}^h - F_{ki}^h) + g_{ih} (F_{kj}^h - F_{jk}^h) + g_{kh} (F_{ij}^h - F_{ji}^h).$$

$$(7.2) 2F_{i\alpha k}^{a} = (\delta_{k}g_{i\alpha}^{a} + \delta_{i}g_{\alpha k}^{a} - \partial_{\alpha}^{a}g_{ik}) -$$

$$-(\lambda_{k}g_{\alpha i}^{a} + \lambda_{i}g_{\alpha k}^{a} - \lambda_{\alpha}^{a}g_{ik}) + B,$$

$$B = (F_{i\alpha k}^{a} - F_{k\alpha i}^{a}) - (F_{\alpha ki}^{a} - C_{ik\alpha}^{a})) - (F_{\alpha ik}^{a} - C_{ki\alpha}^{a}) +$$

$$= g_{\alpha h}^{a}(F_{ik}^{h} - F_{ki}^{h}) - (g_{k\gamma}^{c}F_{\alpha ci}^{a\gamma} - g_{kh}C_{i\alpha}^{ha}) -$$

$$- (g_{i\gamma}^{c}F_{\alpha ck}^{a\gamma} - g_{ik}C_{k\alpha}^{ha}).$$

(7.3)
$$2F_{\alpha ik}^{a} = (\delta_{k}g_{\alpha i}^{a} + \delta_{\alpha}^{a}g_{ik} - \delta_{i}g_{k\alpha}^{a}) - (\lambda_{k}g_{\alpha i}^{a} + \lambda_{\alpha}^{a}g_{ik} - \lambda_{i}g_{k\alpha}^{a}) - B.$$

(7.4)
$$2F_{\alpha\beta k}^{ab} = (\delta_k g_{\alpha\beta}^{ab} + \partial_{\alpha}^a g_{\beta k}^b - \partial_{\beta}^b g_{k\alpha}^a) - (\lambda_k g_{\alpha\beta}^{ab} + \lambda_{\alpha}^a g_{\beta k}^b - \lambda_{\beta}^b g_{k\alpha}^a + C,$$

where

$$C = (F_{\alpha\beta k}^{ab} - C_{k\beta a}^{ba}) - (F_{\beta\alpha k}^{ba} - C_{k\alpha\beta}^{ab}) + (C_{\alpha k\beta}^{ab} - C_{\beta k\alpha}^{ba}) =$$

$$= (g_{\beta\gamma}^{bc} F_{\alpha ck}^{a\gamma} - g_{\beta k}^{b} C_{k\alpha}^{ha}) - (g_{\alpha\gamma}^{ac} F_{\beta ck}^{b} - g_{\alpha h}^{a} C_{k\beta}^{hb}) + g_{k\epsilon}^{e} (C_{\alpha e\beta}^{a\epsilon b} - C_{\beta e\alpha}^{b\epsilon a})$$

$$(7.5) \qquad 2C_{ij\alpha}^{a} = (\partial_{\alpha}^{a} g_{ij} + \delta_{i} g_{j\alpha}^{a} - \delta_{j} g_{\alpha i}^{a}) - (\lambda_{\alpha}^{a} g_{ij} + \lambda_{i} g_{j\alpha}^{a} - \lambda_{j} g_{\alpha i}^{a}) + D,$$

where

$$\begin{split} D &= (F_{i\alpha j}^{~a} - F_{j\alpha i}^{~a}) + (F_{\alpha ij}^{a} - C_{ji\alpha}^{~a}) - (F_{\alpha ji}^{a} - C_{ij\alpha}^{~a}) - \\ &- g_{\alpha h}^{a} (F_{ij}^{h} - F_{ji}^{h}) + (g_{i\gamma}^{~c} F_{\alpha cj}^{a\gamma} - g_{ih} C_{j\alpha}^{~ha}) - \\ &- (g_{j\gamma}^{~c} F_{\alpha ci}^{a\gamma} - g_{jh} C_{i\alpha}^{~ha}). \end{split}$$

(7.6)
$$2C_{k\alpha\beta}^{ab} = (\partial_{\beta}^{b}g_{k\alpha}^{a} + \delta_{k}g_{\alpha\beta}^{ab} - \partial_{\alpha}^{a}g_{\beta k}^{b}) - (\lambda_{\beta}^{b}g_{k\alpha}^{a} + \lambda_{k}g_{\alpha\beta}^{ab} - \partial_{\alpha}^{a}g_{\beta k}^{b}) - E,$$

where

$$E = (F_{\alpha\beta k}^{ab} - C_{k\beta\alpha}^{ba}) + (F_{\beta\alpha k}^{ba} - C_{k\alpha\beta}^{ab}) + (C_{\alpha k\beta}^{ab} - C_{\beta k\alpha}^{ba}) =$$

$$= (g_{\beta\gamma}^{bc} F_{\alpha ck}^{a\gamma} - g_{\beta h}^{b} C_{k\alpha}^{ha}) + (g_{\alpha\gamma}^{ac} F_{\beta ck}^{b\gamma} - g_{\alpha h}^{a} C_{k\beta}^{hb}) + g_{k\epsilon}^{e} (C_{\alpha \epsilon\beta}^{aeb} - C_{\beta \epsilon\alpha}^{bea})$$

(7.7)
$$2C_{\alpha k\beta}^{ab} = (\partial_{\beta}^{b}g_{\alpha k}^{a} + \partial_{\alpha}^{a}g_{k\beta}^{b} - \delta_{k}g_{\beta\alpha}^{ba} - (\lambda_{\beta}^{b}g_{\alpha k}^{a} + \lambda_{\alpha}^{a}g_{k\beta}^{b} - \lambda_{k}g_{\beta\alpha}^{ba}) + E$$

(7.8)
$$2C_{\alpha\beta\gamma}^{abc} = (\partial_{\gamma}^{c} g_{\alpha\beta}^{ab} + \partial_{\alpha}^{a} g_{\beta\gamma}^{bc} - \partial_{\beta}^{b} g_{\gamma\alpha}^{ca}) - (\lambda_{\gamma}^{c} g_{\alpha\beta}^{ab} + \lambda_{\alpha}^{a} g_{\beta\gamma}^{bc} - \lambda_{\beta}^{b} g_{\gamma\alpha}^{ca}) + F,$$

where

$$F = (C_{\alpha\beta\gamma}^{abc} - C_{\gamma\beta\alpha}^{cba}) + (C_{\alpha\gamma\beta}^{acb} - C_{\beta\gamma\alpha}^{bca}) + (C_{\gamma\alpha\beta}^{cab} - C_{\beta\alpha\gamma}^{bac}).$$

Proof. The proof follows from (6.7), (6.11) and (6.12).

In (7.1) - (7.9) the expressions A, B, C, D, E, F and G are functions of torsion tensor T and the nonlinear connection N. The connection coefficients in the K recurrent Hamilton space can be determined using an arbitrary temsor. In the usual terminology, that the space is torsion free when $T(X,Y) = 0 \ \forall X,Y \in T(E^*)$ is used, then we have

Theorem 7.2. In the torsin free recurrent K-Hamilton space, supplied with the metric tensor G, the strongly distinguished connection coefficients are determined by (7.1) - (7.8) where A, B, C, D, E and F have the following value:

$$(7.9) \qquad A = 0,$$

$$B = -(g_{k(\alpha)}^{\ c}\partial_{\alpha}^{a}N_{ci}^{(\alpha)} + g_{i(\alpha)}^{\ c}\partial_{\alpha}^{a}N_{ck}^{(\alpha)},$$

$$C = g_{\beta k(\alpha)}^{bc}\partial_{\alpha}^{a}N_{ck}^{(\alpha)} - g_{\alpha(\beta)}^{ac}\partial_{\beta}^{b}N_{ck}^{(\beta)},$$

$$D = g_{i(\alpha)}^{\ c}\partial_{\alpha}^{a}N_{cj}^{(\alpha)} - g_{j(\alpha)}^{\ c}\partial_{\alpha}^{a}N_{ci}^{(\alpha)},$$

$$E = g_{\beta(\alpha)}^{bc}\partial_{\alpha}^{a}N_{ck}^{(\alpha)} - g_{\alpha(\beta)}^{ac}\partial_{\beta}^{b}N_{ck}^{(\beta)}$$
and $\delta_{i}N_{aj}^{\alpha} - \delta_{j}N_{ai}^{\alpha} = 0$ (which follows from (5.11 f)).

Proof. The proof follows from Theorem 7.1 and Theorem 7.2.

If in (7.1) - (7.8) the filed $\lambda = \lambda_k dx^k + \lambda_\alpha^a \delta p_a^\alpha$ is equal to zero, then the recurrent K- Hamilton space becomes a K- Hamilton space supplied with a strongly distinguished metric connection.

Theorem 7.3. In the torsion free K-Hamilton space supplied with the metric tensor G, the strongly distinguished metric connection coefficients are given by

$$(a) \ 2F_{ijk} = (\delta_k g_{ij} + \delta_i g_{jk} - \delta_j g_{ki})$$

$$(b) \ 2F_{i\alpha k}^{\ a} = (\delta_k g_{i\alpha}^{\ a} + \delta_i g_{\alpha k}^{\ a} - \partial_{\alpha}^{\ a} g_{ik}) - (g_{k(\alpha)}^{\ c} \partial_{\alpha}^{\ a} N_{ci}^{(\alpha)} + g_{i(\alpha)}^{\ c}) \partial_{\alpha}^{\ a} N_{ck}^{(\alpha)})$$

$$(c) \ 2F_{\alpha ik}^{\ a} = (\delta_k g_{\alpha i}^{\ a} + \partial_{\alpha}^{\ a} g_{ik} - \delta_i g_{k\alpha}^{\ a}) + (g_{k(\alpha)}^{\ c} \partial_{\alpha}^{\ a} N_{ci}^{(\alpha)} + g_{i(\alpha)}^{\ c}) \partial_{\alpha}^{\ a} N_{ck}^{(\alpha)})$$

$$(d) \ 2F_{\alpha \beta k}^{\ ab} = (\delta_k g_{\alpha \beta}^{\ ab} + \partial_{\alpha}^{\ a} g_{\beta k}^{\ b} - \partial_{\beta}^{\ b} g_{k\alpha}^{\ a}) + (g_{\beta(\alpha)}^{\ b} \partial_{\alpha}^{\ a} N_{ck}^{(\alpha)} - g_{\alpha(\beta)}^{\ c} \partial_{\beta}^{\ b} N_{ck}^{(\beta)})$$

$$(7.10)$$

$$(e) \ 2C_{ij\alpha}^{\ a} = (\partial_{\alpha}^{\ a} g_{ij} + \delta_i g_{j\alpha}^{\ a} - \delta_j g_{\alpha i}^{\ a}) + (g_{i(\alpha)}^{\ c} \partial_{\alpha}^{\ a} N_{cj}^{(\alpha)} - g_{j(\alpha)}^{\ c} \partial_{\alpha}^{\ a} N_{ci}^{(\alpha)})$$

$$(f) \ 2C_{k\alpha\beta}^{\ ab} = (\partial_{\beta}^{\ b} g_{k\alpha}^{\ a} + \delta_k g_{\alpha\beta}^{\ ab} - \partial_{\alpha}^{\ a} g_{\beta k}^{\ b}) - (g_{\beta(\alpha)}^{\ b} \partial_{\alpha}^{\ a} N_{ck}^{(\alpha)} + g_{\alpha(\beta)}^{\ a} \partial_{\alpha}^{\ a} N_{ck}^{(\beta)})$$

$$(g) \ 2C_{\alpha k\beta}^{\ ab} = (\partial_{\beta}^{\ b} g_{\alpha k}^{\ a} + \partial_{\alpha}^{\ a} g_{\beta \beta}^{\ b} - \delta_k g_{\beta \alpha}^{\ ba}) + (g_{\beta(\alpha)}^{\ b} \partial_{\alpha}^{\ a} N_{ck}^{(\alpha)} + g_{\alpha(\beta)}^{\ a} \partial_{\beta}^{\ b} N_{ck}^{(\beta)})$$

$$(h) \ 2C_{\alpha \beta \gamma}^{\ abc} = (\partial_{\gamma}^{\ c} g_{\alpha \beta}^{\ ab} + \partial_{\alpha}^{\ a} g_{\beta \gamma}^{\ bc} - \partial_{\beta}^{\ b} g_{\gamma \alpha}^{\ ca})$$

$$\delta_i N_{aj}^{\alpha} - \delta_j N_{ai}^{\alpha} = 0.$$

Proof. The proof follows from Theorem 7.1 and Theorem 7.2. \Box

Theorem 7.4. In the torsion free metric K-Hamilton space in which $T_H(E^*)$ is orthogonal to $T_V(E^*)$, i.e. where the metric tensor G has the property $[g_{k\alpha}^{\ a}] = 0 \ \forall \ \alpha = \overline{1, K}$ the strongly distinguished connection coefficients have the form

(a)
$$2F_{ijk} = (\delta_k g_{ij} + \delta_i g_{jk} - \delta_j g_{ki})$$

(b)
$$2F_{i\alpha k}^{\ a} = -\partial_{\alpha}^{a}g_{ik}$$

(c)
$$2F_{\alpha ik}^a = \partial_{\alpha}^a g_{ik}$$

$$(d) \ 2F^{ab}_{\alpha\beta k} = \delta_k g^{ab}_{\alpha\beta} + g^{b\ c}_{\beta(\alpha)} \partial^a_\alpha N^{(\alpha)}_{ck} - g^{a\ c}_{\alpha(\beta)} \partial^b_\beta N^{(\beta)}_{ck}$$

(e)
$$2C_{ij\alpha}^{\ a} = \partial_{\alpha}^{a}g_{ij}$$

$$(f) \ 2C_{k\alpha\beta}^{\ ab} = \delta_k g_{\alpha\beta}^{\ ab} - (g_{\beta(\alpha)}^{\ b\ c} \partial_{\alpha}^{\ a} N_{ck}^{(\alpha)} + g_{\alpha(\beta)}^{\ a\ c} \partial_{\beta}^{\ b} N_{ck}^{(\beta)})$$

$$(g) \ 2C_{\alpha k\beta}^{ab} = -\delta_k g_{\beta\alpha}^{ba} + g_{\beta(\alpha)}^{bc} \partial_\alpha^a N_{ck}^{(\alpha)} + g_{\alpha(\beta)}^{ac} \partial_\beta^b N_{ck}^{(\beta)})$$

$$(h) \ \ 2C^{abc}_{\alpha\beta\gamma} = \partial^c_{\gamma}g^{ab}_{\alpha\beta} + \partial^a_{\alpha}g^{bc}_{\beta\gamma} - \partial^b_{\beta}g^{ca}_{\gamma\alpha}.$$

Proof. The proof follows from (7.10) and the condition $[g_{k\alpha}^{\ a}] = 0$ for $\forall \alpha = \overline{1, K}$ \square .

It is obvious that in (7.11)

$$F^a_{\alpha ih} = -F^a_{i\alpha h}$$
 $C^{ab}_{k\alpha\beta} = -C^{ab}_{\alpha k\beta}$.

In the recurrent K-Hamilton space in which $T_H(E^*)$ is orthogonal to $T_V(E^*)$ i.e. where the metric tensor has the property $[g_{i\alpha}^{\ a}] = 0$ for $\forall \alpha \overline{1, K}$, from (6.11) we get

(a)
$$F_{ik}^{h} = g^{hj}F_{ijh}$$

(b)
$$0 = g_{ad}^{\alpha\delta} F_{i\alpha k}^{\ a}$$

(c)
$$0 = g^{hi}F^a_{\alpha ik}$$

(d)
$$F_{ak}^{(\alpha)} = g_{ab}^{(\alpha)\beta} F_{\alpha\beta k}^{ab}$$

(7.12)

(e)
$$C_i^{\ hc}_{\ \gamma} = g^{hj}C_{ij\gamma}^{\ c}$$

(f)
$$0 = g_{da}^{\delta\alpha} C_{i\alpha\gamma}^{\ ac}$$

(g)
$$0 = g^{hi}C^a_{\alpha i\gamma}^c$$

(h)
$$C^{(\alpha)}_{d\gamma}^{ac} = g^{(\alpha)\beta}_{db}C^{abc}_{\alpha\beta\gamma}$$
.

The strongly distinguished connection coefficients which appear on the right-hand side of (7.12) for the torsion free, metric K-Hamilton space are determined by (7.11).

For the torsion free strongly distinguished connection for which $(\partial_{\alpha}^{a} N_{bk}^{\beta} = 0 \ \forall \alpha, \beta = \overline{1, K}$ the relation (5.14) holds. These relations may be obtained for the above case also from (7.11) and (7.12).

References

- Anastasiei M., Models of Finsler and Lagrange geometry. The Proc. of IVth National Seminar of Finsler and Lagrange spaces. Univ. Brasov, (R.S.Romania) 1986, 43-56.
- [2] Aringazin A.K., Asanov G.S., Problems of Finslerian theory of gauge fields and gravitation. Rep. Math. Physics 25 (1988), 35-93.
- [3] Čomić I., A Generalization to d-connection, Tensor N.S. Vol. 48, 1989, 199-208.
- [4] Čomić I., Curvature theory of generalized Mirons d-connection. Differential Geometry and Its Application Proc. Conf. 1984 Brno, 17-26.
- [5] Comić I., Recurrent k-Lagrangian spaces with a strongly distinguished connection (to appear)
- [6] Čomić I., Kirhovits Sz. M. Generalized Mirons d-connection in a k-Lagrangian space, Proceedings of VI-th Conference on Finsler geometry, Brasov 1990, 39-58.
- [7] Comić I., Generalized Mirons d-connection in the recurrent K-Hamilton space (to appear in Publ. de l Inst. Math. Beograd)

- [8] Čomić I., Recurrent Hamilton spaces with generalized Mirons dconnection (to appear)
- [9] Kirhovits Sz.M., On k-Hamilton Geometry I (to appear in Math. Pannnonica).
- [10] Kirhovits Sz.M., On k-Hamilton Geometry II (to appear in Math. Pannonica).
- [11] Miron R., Anastasiei M., Fibrate vectoriale spatii Lagrange aplicatii, Bucuresti, 1987.
- [12] Miron R., On the Finslerian theory of relativity. Tensor N.S. 44, 1987, 63-81.
- [13] Miron R., Hamilton geometry. Seminarul de Mecanica, Univ. Timisoara (R.S.) Preprint Nr. 3., 1987, 54 pg.
- [14] Miron R., Anastasiei M., Vector bundles. Lagrange spaces. Applications to Relativity (in Romanian) Edit. Acad. R.S. Romania, 1987.
- [15] Miron R., Janus S., Anastasiei M., The geometry of the dual of a vector bundle, Publ. Inst. Math. (Beograd) 46(60) 1989, 145-162.
- [16] Opris D., Lagrange Hamilton geometry Seminarul de Mecanica Univ. din Timisoare 1988, 26 pp.
- [17] Sakaguchi T., Subspaces in Lagrange spaces Ph. Thesis, Univ. Iasi (R.S. Romania), 1988, 167 pg.

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