

NEW SPECIAL GEODESIC MAPPINGS OF GENERAL AFFINE CONNECTION SPACES

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Abstract

In this work¹ we define R_θ -projective geodesic mappings ($\theta = 1, \dots, 5$) of two general affine connection spaces and obtain some invariant geometric objects of these mappings, generalizing Weyl's tensor. Also, we define R_θ -projectively flat affine connection spaces $G\bar{A}_N$ and find necessary conditions for the space GA_N to be R_θ -projectively flat.

Introduction

Consider two N -dimensional differentiable manifolds GA_N and $G\bar{A}_N$ and differentiable mapping $f : GA_N \rightarrow G\bar{A}_N$. We can consider these manifolds in the common by this mapping system of local coordinates. If the connection coefficients $L_{jk}^i(x)$ and $\bar{L}_{jk}^i(x)$, for the connection introduced in GA_N and $G\bar{A}_N$ respectively, are non-symmetric in lower indices, we call GA_N and $G\bar{A}_N$ *general affine connection spaces*.

One says that reciprocal one valued mapping $f : GA_N \rightarrow G\bar{A}_N$ is *geodesic*, [5,6] if geodesics of the space GA_N pass to geodesics of the space $G\bar{A}_N$. In the corresponding points $M(x)$ and $\bar{M}(x)$ we can put

$$(0.1) \quad \bar{L}_{jk}^i(x) = L_{jk}^i(x) + P_{jk}^i(x), \quad (i, j, k = 1, \dots, N),$$

where $P_{jk}^i(x)$ is the *deformation tensor* of the connection L of GA_N according to the mapping $f : GA_N \rightarrow G\bar{A}_N$.

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A necessary and sufficient condition that the mapping f be geodesic [5] is that the deformaton tensor P_{jk}^i from (0.2) has the form

$$(0.2) \quad P_{jk}^i(x) = \delta_j^i \psi_k(x) + \delta_k^i \psi_j(x) + \xi_{jk}^i(x),$$

where

$$(0.3) \quad \psi_i = \frac{1}{N+1} (\bar{L}_{ip}^p - L_{ip}^p),$$

$$(0.4) \quad \xi_{jk}^i = \bar{L}_{jk}^i - L_{jk}^i$$

and \underline{jk} denotes symmetrisation and $\overset{\vee}{jk}$ -antisymmetrisation with respect to j, k . In GA_N can define four kinds of covariant derivatives [1,2,3]. For example, for a tensor a_j^i in GA_N we have

$$\begin{aligned} a_{j|_1^i}^i &= a_{j,m}^i + L_{pm}^i a_j^p - L_{jm}^p a_p^i, & a_{j|_2^i}^i &= a_{j,m}^i + L_{mp}^i a_j^p - L_{mj}^p a_p^i, \\ a_{j|_3^i}^i &= a_{j,m}^i + L_{pm}^i a_j^p - L_{mj}^p a_p^i, & a_{j|_4^i}^i &= a_{j,m}^i + L_{mp}^i a_j^p - L_{jm}^p a_p^i. \end{aligned}$$

Denote by $\left|, \right|_{\theta}$ a covariant derivative of the kind θ in GA_N and $G\bar{A}_N$ respectively.

In the case of the space GA_N we have five independent curvature tensors [4] (in [4] R_5 is denoted by \bar{R}_2):

$$\begin{aligned} R_1^i{}_{jmn} &= L_{jm,n}^i - L_{jn,m}^i + L_{jm}^p L_{pn}^i - L_{jn}^p L_{pm}^i, \\ R_2^i{}_{jmn} &= L_{mj,n}^i - L_{nj,m}^i + L_{mj}^p L_{np}^i - L_{nj}^p L_{mp}^i, \\ R_3^i{}_{jmn} &= L_{jm,n}^i - L_{nj,m}^i + L_{jm}^p L_{np}^i - L_{nj}^p L_{pm}^i + L_{nm}^p (L_{pj}^i - L_{jp}^i), \\ R_4^i{}_{jmn} &= L_{jm,n}^i - L_{nj,m}^i + L_{jm}^p L_{np}^i - L_{nj}^p L_{pm}^i + L_{mn}^p (L_{pj}^i - L_{jp}^i), \\ R_5^i{}_{jmn} &= \frac{1}{2} (L_{jm,n}^i + L_{mj,n}^i - L_{jn,m}^i - L_{nj,m}^i + L_{jm}^p L_{pn}^i + L_{mj}^p L_{np}^i \\ &\quad - L_{jn}^p L_{mp}^i - L_{nj}^p L_{pm}^i). \end{aligned}$$

By virtue of the geodesic mapping $f : GA_N \rightarrow G\bar{A}_N$ we obtain tensors $\bar{R}_{\theta}^i{}_{jmn}$ ($\theta = 1, \dots, 5$), where for example

$$(0.5) \quad \bar{R}_1^i{}_{jmn} = \bar{L}_{jm,n}^i - \bar{L}_{jn,m}^i + \bar{L}_{jm}^p \bar{L}_{pn}^i - \bar{L}_{jn}^p \bar{L}_{pm}^i.$$

In the case of geodesic mapping $f : A_N \rightarrow \bar{A}_N$ of the symmetric affine connection spaces A_N and \bar{A}_N we have an invariant geometric object

$$(0.6) \quad W^i_{jmn} = R^i_{jmn} + \frac{2}{N+1} \delta_j^i R_{[mn]} + \frac{1}{N^2-1} [\delta_m^i (NR_{jn} + R_{nj}) - \delta_n^i (NR_{jm} + R_{mj})],$$

where R^i_{jmn} is Riemann-Cristoffel's curvature tensor of the space A_N , and R_{jm} Richi's tensor.

The object W^i_{jmn} is called Weil's tensor, or a tensor of projective curvature [8]. Having a geodesic mapping of two general affine connection spaces, we can not find a generalization of Weil's tensor as an invariant of geodesic mapping in general case. For that reason we define a special geodesic mapping.

1. \bar{R} -projective mappings

For the first kind curvature tensors of the spaces GA_N and $G\bar{A}_N$ respectively we find the relation

$$(1.1) \quad \begin{aligned} \bar{R}^i_{jmn} &= R^i_{jmn} + \delta_j^i (\psi_{mn} - \psi_{nm}) + \delta_m^i \psi_{jn} - \delta_n^i \psi_{jm} \\ &\quad - \delta_m^i \xi_{jn}^p \psi_p + \delta_n^i \xi_{jm}^p \psi_p + \xi_{jm|n}^i - \xi_{jn|m}^i + 2\psi_j \xi_{mn}^i \\ &\quad + \xi_{jm}^p \xi_{pn}^i - \xi_{jn}^p \xi_{pm}^i + 2L_{mn}^i \psi_j + 2L_{mn}^p \psi_p \delta_j^i + 2L_{mn}^p \xi_{jp}^i, \\ \psi_{mn} &= \psi_{m|n} - \psi_m \psi_n, \quad (\theta = 1, 2). \end{aligned}$$

Contracting with respect to i and n from (1.1) we get

$$(1.2) \quad \begin{aligned} \bar{R}_{jm} &= R_{jm} - \psi_{[jm]} - (N-1)\psi_{jm} + (N-1)\xi_{jm}^p \psi_p + \xi_{jm|p}^p \\ &\quad - \xi_{jp|m}^p + \psi_j \bar{L}_{mp}^p + \xi_{jm}^p \xi_{pq}^p - \xi_{jq}^p \xi_{pm}^q + 2L_{mj}^p \psi_p + 2L_{mq}^p \xi_{jp}^q \end{aligned}$$

Here \bar{R}_{jm} and R_{jm} are the first kind Ricci tensors of the spaces $G\bar{A}_N$ and GA_N respectively and $[jm]$ denotes an alternation without a division. From (1.2) we obtain

$$(1.3) \quad \begin{aligned} \bar{R}_{[jm]} &= R_{[jm]} - 2\psi_{[jm]} - (N-1)\psi_{[jm]} + 2(N-1)\xi_{jm}^p \psi_p \\ &\quad + \xi_{jm|p}^p - \xi_{jp|m}^p + \xi_{mp|j}^p + 2\psi_j \bar{L}_{mp}^p - 2\psi_m \bar{L}_{jp}^p + 2\xi_{jm}^p \xi_{pq}^p \\ &\quad - \xi_{jq}^p \xi_{pm}^q + \xi_{mq}^p \xi_{pj}^q + 4L_{mj}^p \psi_p + 2L_{mq}^p \xi_{jp}^q - 2L_{jq}^p \xi_{mp}^q, \end{aligned}$$

from where

$$(1.4) \quad \begin{aligned} (N+1)\psi_{1[jm]} &= R_{1[jm]} - \bar{R}_{1[jm]} + 2(N-1)\xi_{jm}^p \psi_p + \xi_{jm|p}^p - \xi_{jp|m}^p \\ &+ \xi_{mp|j}^p + 2\psi_j \bar{L}_{mp}^p - 2\psi_m \bar{L}_{jp}^p + 2\xi_{jm}^p \xi_{pq}^q + 4L_{mj}^p \psi_p + 2L_{mq}^p \xi_{jp}^q - 2L_{jq}^p \xi_{mp}^q. \end{aligned}$$

Substituting (1.4) in (1.2), we get

$$(1.5) \quad \begin{aligned} (N-1)\psi_{jm} &= R_{jm} - \bar{R}_{jm} - \frac{1}{N+1} [R_{1[jm]} - \bar{R}_{1[jm]}] + 2(N-1)\xi_{jm}^p \psi_p \\ &+ 2\xi_{jm|p}^p - \xi_{jp|m}^p + \xi_{mp|j}^p + 2\psi_j \bar{L}_{mp}^p - 2\psi_m \bar{L}_{jp}^p + 2\xi_{jm}^p \xi_{pq}^q \\ &+ 4L_{mj}^p \psi_p + 2L_{mq}^p \xi_{jp}^q - 2L_{jq}^p \xi_{mp}^q] + (N-1)\xi_{jm}^p \psi_p + \xi_{jm|p}^p - \xi_{jp|m}^p \\ &+ 2\psi_j \bar{L}_{mp}^p + \xi_{jm}^p \xi_{pq}^q - \xi_{jq}^p \xi_{pm}^q + 2L_{mj}^p \psi_p + 2L_{mq}^p \xi_{jp}^q. \end{aligned}$$

Let us denote

$$(1.6) \quad \begin{aligned} \mathcal{D}_{jm} &= \frac{N-1}{N+1} \xi_{jm}^p \psi_p + \frac{2}{N^2-1} (N\psi_j \bar{L}_{mp}^p - \psi_m \bar{L}_{jp}^p + NL_{mq}^p \xi_{jp}^q + L_{jq}^p \xi_{mp}^q) \\ &+ \frac{1}{N+1} \xi_{jm}^p \xi_{pq}^q - \frac{1}{N-1} \xi_{jq}^p \xi_{pm}^q + \frac{2}{N+1} L_{mj}^p \psi_p. \end{aligned}$$

Now, (1.5) we can express in the form

$$(1.7) \quad \begin{aligned} \psi_{jm} &= \frac{1}{N-1} \{ R_{jm} - \bar{R}_{jm} - \frac{1}{N+1} [R_{1[jm]} - \bar{R}_{1[jm]}] \\ &+ 2\xi_{jm|p}^p - \xi_{jp|m}^p + \xi_{mp|j}^p \} + \xi_{jm|p}^p - \xi_{jp|m}^p + \mathcal{D}_{jm}, \end{aligned}$$

i.e.

$$(1.8) \quad \begin{aligned} \psi_{jm} &= \frac{1}{N-1} (R_{jm} - \bar{R}_{jm}) - \frac{1}{N^2-1} (R_{1[jm]} - \bar{R}_{1[jm]}) \\ &+ \frac{1}{N+1} (\bar{L}_{jm|p}^p - L_{jm|p}^p) - \frac{N}{N+1} (\bar{L}_{jp|m}^p - L_{jp|m}^p) \\ &- \frac{1}{N^2-1} (\bar{L}_{mp|j}^p - L_{mp|j}^p) + \mathcal{D}_{jm}. \end{aligned}$$

Substituting (1.8) in (1.1) one obtains

$$\begin{aligned}
 \bar{R}_{jmn}^i &= R_{jmn}^i + \frac{1}{N-1} \delta_j^i (R_{[mn]} - \bar{R}_{[mn]}) \\
 &- \frac{1}{N^2-1} \delta_j^i (2R_{[mn]} - 2\bar{R}_{[mn]}) + \frac{2}{N+1} \delta_j^i (\bar{L}_{\sqrt{1}n|p}^p - L_{\sqrt{1}n|p}^p) \\
 (1.9) \quad &- \frac{N}{N+1} \delta_j^i (\bar{L}_{\sqrt{1}n|p}^p - \bar{L}_{\sqrt{1}n|m}^p - L_{\sqrt{1}n|p}^p + L_{\sqrt{1}n|m}^p) - \frac{1}{N^2-1} \delta_j^i (\bar{L}_{\sqrt{1}m|p}^p - \bar{L}_{\sqrt{1}m|n}^p \\
 &- L_{\sqrt{1}m|p}^p + L_{\sqrt{1}m|n}^p) + \delta_j^i \mathcal{D}_{mn} + \frac{1}{N-1} \delta_m^i (R_{jn} - \bar{R}_{jn}) \\
 &- \frac{1}{N^2-1} \delta_m^i (R_{[jn]} - \bar{R}_{[jn]}) + \frac{1}{N+1} \delta_m^i (\bar{L}_{\sqrt{1}j|p}^p - L_{\sqrt{1}j|p}^p) \\
 &- \frac{N}{N+1} \delta_m^i (\bar{L}_{\sqrt{1}j|n}^p - L_{\sqrt{1}j|n}^p) - \frac{1}{N^2-1} \delta_m^i (\bar{L}_{\sqrt{1}j|p}^p - L_{\sqrt{1}j|p}^p) + \delta_m^i \mathcal{D}_{jn} \\
 &- \frac{1}{N-1} \delta_n^i (R_{jm} - \bar{R}_{jm}) + \frac{1}{N^2-1} \delta_n^i (R_{[jm]} - \bar{R}_{[jm]}) \\
 (1.9) \quad &- \frac{1}{N+1} \delta_n^i (\bar{L}_{\sqrt{1}j|p}^p - L_{\sqrt{1}j|p}^p) - \frac{N}{N+1} \delta_m^i (\bar{L}_{\sqrt{1}j|m}^p - L_{\sqrt{1}j|m}^p) \\
 &+ \frac{1}{N^2-1} \delta_n^i (\bar{L}_{\sqrt{1}m|j}^p - L_{\sqrt{1}m|j}^p) - \delta_n^i \mathcal{D}_{jm} - \delta_m^i \xi_{jn}^p \psi_p + \delta_n^i \xi_{jm}^p \psi_p + \xi_{jm|n}^i \\
 &- \xi_{jn|m}^i + 2\psi_j \xi_{mn}^i + \xi_{jm}^p \xi_{pn}^i - \xi_{jn}^p \xi_{pm}^i + 2L_{\sqrt{1}mn}^i \psi_j + 2L_{\sqrt{1}mn}^p \psi_p \delta_j^i + 2L_{\sqrt{1}mn}^p \xi_{jp}^i.
 \end{aligned}$$

Introducing the condition

$$\begin{aligned}
 (1.10) \quad &\delta_j^i \mathcal{D}_{[mn]} + \delta_m^i \mathcal{D}_{jn} - \delta_n^i \mathcal{D}_{jm} - \delta_m^i \xi_{jn}^p \psi_p + \delta_n^i \xi_{jm}^p \psi_p + 2\psi_j \xi_{mn}^i \\
 &+ \xi_{jm}^p \xi_{pn}^i - \xi_{jn}^p \xi_{pm}^i + 2L_{\sqrt{1}mn}^i \psi_j + 2L_{\sqrt{1}mn}^p \psi_p \delta_j^i + 2L_{\sqrt{1}mn}^p \xi_{jp}^i = 0,
 \end{aligned}$$

(1.9) can be expressed in the form

$$(1.11) \quad \bar{W}(R)_1^i{}_{jmn} = W(R)_1^i{}_{jmn}$$

where

$$\begin{aligned}
 W(R)_1^i{}_{jmn} &= R_{jmn}^i + \frac{1}{N+1} \delta_j^i R_{[mn]} + \frac{1}{N^2-1} [(NR_{1jn} + R_{1nj}) \delta_m^i \\
 &- (NR_{1jm} + R_{1mj}) \delta_n^i] - \frac{2}{N+1} \delta_j^i L_{\sqrt{1}mn|p}^p + \frac{N^2-N-1}{N^2-1} \delta_j^i (L_{\sqrt{1}n|p}^p - L_{\sqrt{1}n|m}^p) \\
 (1.12) \quad &- \frac{1}{N+1} \delta_m^i L_{\sqrt{1}n|p}^p + \frac{1}{N+1} \delta_m^i (NL_{\sqrt{1}j|n}^p + \frac{1}{N-1} L_{\sqrt{1}n|j}^p) + \frac{1}{N+1} \delta_n^i L_{\sqrt{1}j|m}^p \\
 &- \frac{1}{N+1} \delta_n^i (NL_{\sqrt{1}j|m}^p + \frac{1}{N-1} L_{\sqrt{1}m|j}^p) - L_{\sqrt{1}j|m}^i + L_{\sqrt{1}j|n}^i.
 \end{aligned}$$

Definition 1.1. The geodesic mapping $f : GA_N \rightarrow G\bar{A}_N$ is R -projective if the condition (1.10) is satisfied.

Definition 1.2. The space GA_N is R -projectively flat if there exists an R -projective mapping of the space GA_N to a flat space (i.e. to a space, whose connection coefficients in special coordinates are zero). So, we proved

Theorem 1.1. The tensor (1.12) is an invariant of an R -projective mapping.

Theorem 1.2. If GA_N is R -projectively flat, then

$$(1.13) \quad \mathcal{W}(\bar{R})^i_{jmn} = 0.$$

Proof. $G\bar{A}_N$ is a flat space. Then by (1.12) we get $\bar{\mathcal{W}}(\bar{R})^i_{jmn} = 0$, and using (1.11) we can see that (1.13) holds.

2. R_2 -projective mappings

Definition 2.1. The geodesic mapping $f : GA_N \rightarrow G\bar{A}_N$ is R_2 -projective if the following condition is satisfied

$$(2.1) \quad \begin{aligned} \delta_j^i D_{2[mn]} + \delta_m^i D_{2jn} - \delta_n^i D_{2jm} - \delta_m^i \xi_{nj}^p \psi_p + \delta_n^i \xi_{mj}^p \psi_p + 2\psi_j \xi_{nm}^i \\ + \xi_{mj}^p \xi_{np}^i - \xi_{nj}^p \xi_{mp}^i + 2L_{nm}^i \psi_j + 2L_{nm}^p \psi_p \delta_j^i + 2L_{nm}^p \xi_{pj}^i = 0. \end{aligned}$$

where

$$(2.2) \quad \begin{aligned} D_{2jm} = \frac{N-1}{N+1} \xi_{mj}^p \psi_p + \frac{2}{N^2-1} (N\psi_j \bar{L}_{pm}^p - \psi_m \bar{L}_{pj}^p + NL_{qm}^p \xi_{pj}^q \\ + L_{qv}^p \xi_{pm}^q) + \frac{1}{N+1} \xi_{mj}^p \xi_{qp}^q - \frac{1}{N-1} \xi_{qj}^p \xi_{mp}^q + \frac{2}{N+1} L_{jm}^p \psi_p, \end{aligned}$$

For the curvature tensors R and \bar{R} of the space GA_N and $G\bar{A}_N$ we have the relation

$$(2.3) \quad \begin{aligned} \bar{R}_{2jmn}^i = R_{2jmn}^i + \delta_j^i (\psi_{mn} - \psi_{nm}) + \delta_m^i \psi_{jn} - \delta_n^i \psi_{jm} \\ - \delta_m^i \xi_{nj}^p \psi_p + \delta_n^i \xi_{mj}^p \psi_p + \xi_{mj}^i |_{2n} - \xi_{nj}^i |_{2m} + 2\psi_j \xi_{nm}^i \\ + \xi_{mj}^p \xi_{np}^i - \xi_{nj}^p \xi_{mp}^i + 2L_{nm}^i \psi_j + 2L_{nm}^p \psi_p \delta_j^i + 2L_{nm}^p \xi_{pj}^i. \end{aligned}$$

Analogously to the previous case, we get an invariant of an \bar{R}_2 -projective mapping $f : GA_N \rightarrow G\bar{A}_N$

$$(2.4) \quad \begin{aligned} \mathcal{W}(\bar{R}_2)^i{}_{jmn} &= \bar{R}_2^i{}_{jmn} + \frac{1}{N+1} \delta_j^i \bar{R}_2^i{}_{[mn]} + \frac{1}{N^2-1} [(N\bar{R}_2^{jn} + \bar{R}_2^{nj}) \delta_m^i \\ &- (N\bar{R}_2^{jm} + \bar{R}_2^{mj}) \delta_n^i] - \frac{2}{N+1} \delta_j^i L_{\sqrt{2}}^p{}_{n|m} + \frac{N^2-N-1}{N^2-1} \delta_j^i (L_{\sqrt{2}}^p{}_{m|n} - L_{\sqrt{2}}^p{}_{n|m}) \\ &- \frac{1}{N+1} \delta_m^i L_{\sqrt{2}}^p{}_{n|j} + \frac{1}{N+1} \delta_m^i (NL_{\sqrt{2}}^p{}_{pj|n} + \frac{1}{N-1} L_{\sqrt{2}}^p{}_{pn|j}) + \frac{1}{N+1} \delta_n^i L_{\sqrt{2}}^p{}_{m|j} \\ &- \frac{1}{N+1} \delta_n^i (NL_{\sqrt{2}}^p{}_{pj|m} + \frac{1}{N-1} L_{\sqrt{2}}^p{}_{pm|j}) - L_{\sqrt{2}}^i{}_{mj|n} + L_{\sqrt{2}}^i{}_{nj|m}. \end{aligned}$$

Consequently, the next theorems are valid

Theorem 2.1. *The tensor (2.4) is an invariant of an \bar{R}_2 -projective mapping $f : GA_N \rightarrow G\bar{A}_N$.*

Theorem 2.2. *If GA_N is \bar{R}_2 -projectively flat then we have*

$$\mathcal{W}(\bar{R}_2)^i{}_{jmn} = 0.$$

3. \bar{R}_3 -projective mappings

Definition 3.1. The geodesic mapping $f : GA_N \rightarrow G\bar{A}_N$ is \bar{R}_3 -projective if the following condition holds

$$(3.1) \quad \begin{aligned} &\delta_j^i \mathcal{D}_3^i{}_{[mn]} + \delta_m^i \mathcal{D}_3^i{}_{jn} - \delta_n^i \mathcal{D}_3^i{}_{jm} + 2\delta_j^i L_{\sqrt{3}}^p{}_{m\nu} \psi_p + 2\delta_m^i L_{\sqrt{3}}^p{}_{jn} \psi_p \\ &+ (\delta_n^i \xi_{jm}^p - \delta_m^i \xi_{nj}^p) \psi_p + \xi_{jm}^p \xi_{np}^i - \xi_{jn}^p \xi_{pm}^i + 2\psi_n (L_{\sqrt{3}}^i{}_{mj} + \xi_{mj}^i) \\ &+ 2\psi_m (L_{\sqrt{3}}^i{}_{nj} + \xi_{nj}^i) + 2\xi_{nm}^p (L_{\sqrt{3}}^i{}_{pj} + \xi_{pj}^i) = 0. \end{aligned}$$

where

$$\begin{aligned} \mathcal{D}_3^i{}_{jm} &= \xi_{jm}^p \psi_p + \frac{1}{N+1} \xi_{jm}^p \xi_{qp}^q - \frac{2}{N^2-1} [2\psi_p (L_{\sqrt{3}}^p{}_{mj} + \xi_{mj}^p) \\ &+ \psi_m (L_{\sqrt{3}}^p{}_{pj} + \xi_{pj}^p) - \psi_j (L_{\sqrt{3}}^p{}_{pm} + \xi_{pm}^p) + \xi_{qm}^p L_{\sqrt{3}}^q{}_{pj} - \xi_{qj}^p L_{\sqrt{3}}^q{}_{pm}] \\ &+ \frac{1}{N-1} [2\psi_p (L_{\sqrt{3}}^p{}_{mj} + \xi_{mj}^p) + 2\psi_m (L_{\sqrt{3}}^p{}_{pj} + \xi_{pj}^p) + 2\xi_{qm}^p (L_{\sqrt{3}}^q{}_{pj} + \xi_{pj}^q) - \xi_{jq}^p \xi_{pm}^q] \end{aligned}$$

Definition 3.2. The space GA_N is \bar{R}_3 -projectively flat if there exists an \bar{R}_3 -projective mapping of the space GA_N into a flat space.

In the case of curvature tensors of the third kind of the spaces GA_N and \bar{GA}_N we get the relation

$$(3.2) \quad \begin{aligned} \bar{R}_3^i{}_{jmn} &= R_3^i{}_{jmn} + \delta_j^i (\psi_{mn} - \psi_{nm}) + \delta_m^i \psi_{jn} - \delta_n^i \psi_{jm} \\ &+ \psi_p (\delta_n^i \xi_{jm}^p - \delta_m^i \xi_{nj}^p) + \xi_{jm|n}^i - \xi_{nj|m}^i + \xi_{jm}^p \xi_{np}^i - \xi_{jn}^p \xi_{pm}^i \\ &+ 2\psi_n (L_{m\check{v}}^i + \xi_{mj}^i) + 2\psi_m (L_{n\check{v}}^i + \xi_{nj}^i) + 2\xi_{nm}^p (L_{p\check{v}}^i + \xi_{pj}^i). \end{aligned}$$

Also, it holds

$$(3.3) \quad \psi_{mn} = \psi_{mn} + 2L_{m\check{v}}^p \psi_p.$$

From (3.2) and (3.3) we get

$$(3.4) \quad \begin{aligned} \bar{R}_3^i{}_{jmn} &= R_3^i{}_{jmn} + \delta_j^i (\psi_{mn} - \psi_{nm}) + \delta_m^i \psi_{jn} - \delta_n^i \psi_{jm} + 2\delta_j^i L_{m\check{v}}^p \psi_p \\ &+ 2\delta_m^i L_{j\check{v}}^p \psi_p + (\delta_n^i \xi_{jm}^p - \delta_m^i \xi_{nj}^p) \psi_p + \xi_{jm|n}^i - \xi_{nj|m}^i + \xi_{jm}^p \xi_{np}^i - \xi_{jn}^p \xi_{pm}^i \\ &+ 2\psi_n (L_{m\check{v}}^i + \xi_{mj}^i) + 2\psi_m (L_{n\check{v}}^i + \xi_{nj}^i) + 2\xi_{nm}^p (L_{p\check{v}}^i + \xi_{pj}^i). \end{aligned}$$

Contracting (3.4) with respect to i and n we get

$$(3.5) \quad \begin{aligned} \bar{R}_3^i{}_{jm} &= R_3^i{}_{jm} - \psi_{[jm]} - (N-1)\psi_{jm} + (N+1)\xi_{jm}^p \psi_p + \xi_{jm|p}^p - \xi_{pj|m}^p \\ &+ \xi_{jm}^p \xi_{qp}^q - \xi_{jq}^p \xi_{pm}^q + 2\psi_p (L_{m\check{v}}^p + \xi_{mj}^p) + 2\psi_m (L_{p\check{v}}^p + \xi_{pj}^p) + 2\xi_{qm}^p (L_{p\check{v}}^q + \xi_{pj}^q), \end{aligned}$$

hence

$$(3.6) \quad \begin{aligned} (N+1)\psi_{[jm]} &= R_3^i{}_{[jm]} - \bar{R}_3^i{}_{[jm]} + 2(N+1)\xi_{jm}^p \psi_p + 2\xi_{jm|p}^p - \xi_{pj|m}^p + \xi_{pm|j}^p \\ &+ 2\xi_{jm}^p \xi_{qp}^q - \xi_{jq}^p \xi_{pm}^q + \xi_{mq}^p \xi_{pj}^q + 4\psi_p (L_{m\check{v}}^p + \xi_{mj}^p) + 2\psi_m (L_{p\check{v}}^p + \xi_{pj}^p) \\ &- 2\psi_j (L_{p\check{v}}^p + \xi_{pm}^p) + 2\xi_{qm}^p (L_{p\check{v}}^q + \xi_{pj}^q) - 2\xi_{qj}^p (L_{p\check{v}}^q + \xi_{pm}^q). \end{aligned}$$

Substituting (3.6) in (3.5) we get

$$(3.7) \quad \begin{aligned} (N-1)\psi_{jm} &= R_3^i{}_{jm} - \bar{R}_3^i{}_{jm} - \frac{1}{N+1} [R_3^i{}_{[jm]} - \bar{R}_3^i{}_{[jm]} + 2\xi_{jm|p}^p \\ &- \xi_{pj|m}^p + \xi_{pm|j}^p] + \xi_{jm|p}^p - \xi_{pj|m}^p + (N-1)\mathcal{D}_{jm}^i, \end{aligned}$$

Now, from (3.4,6,7) by condition (3.1) we get

$$(3.8) \quad \overline{\mathcal{W}}(R)_3^i{}_{jmn} = \mathcal{W}(R)_3^i{}_{jmn}$$

where

$$\begin{aligned} \mathcal{W}(R)_3^i{}_{jmn} = & R_3^i{}_{jmn} + \frac{1}{N+1} \delta_j^i R_{[mn]} + \frac{1}{N^2-1} [(NR_{3jn} + R_{3nj})\delta_m^i \\ & - (NR_{3jm} + R_{3mj})\delta_n^i] + \frac{2}{N^2-1} \delta_j^i (2L_{\sqrt{2}m|p}^p - L_{\sqrt{1}pn}^p + L_{\sqrt{1}m|p}^p - L_{\sqrt{2}mn|p}^p) \\ & + \frac{1}{N-1} \delta_j^i (L_{\sqrt{1}pm|n}^p - L_{\sqrt{1}pn|m}^p) + \frac{1}{N^2-1} \delta_m^i (2L_{\sqrt{2}jn|p}^p - L_{\sqrt{1}pj|n}^p + L_{\sqrt{1}pn|j}^p) \\ & - \frac{1}{N-1} \delta_m^i (L_{\sqrt{2}jn|p}^p - L_{\sqrt{1}pj|n}^p) - \frac{1}{N^2-1} \delta_n^i (2L_{\sqrt{2}jm|p}^p - L_{\sqrt{1}pj|m}^p + L_{\sqrt{1}pm|j}^p) \\ & + \frac{1}{N-1} \delta_n^i (L_{\sqrt{2}jm|p}^p - L_{\sqrt{1}pj|m}^p) - L_{\sqrt{2}jm|n}^i + L_{\sqrt{1}nj|m}^i \end{aligned}$$

Consequently, the next theorems hold:

Theorem 3.1. *The tensor $\mathcal{W}(R)_3^i{}_{jmn}$ is an invariant of an R_3 -projective mapping.*

Theorem 3.2. *If GA_N is R_3 -projectively flat then*

$$\mathcal{W}(R)_3^i{}_{jmn} = 0.$$

4. R_4 -projective mappings

Definition 4.1. A geodesic mapping $f : GA_N \rightarrow \overline{GA}_N$ is R_4 -projective if the following condition is satisfied

$$(4.1) \quad \begin{aligned} & \delta_j^i \mathcal{D}_4^i{}_{[mn]} + \delta_m^i \mathcal{D}_4^i{}_{jn} - \delta_n^i \mathcal{D}_4^i{}_{jm} + 2\delta_j^i L_{\sqrt{2}mn}^p \psi_p + 2\delta_m^i L_{\sqrt{2}jn}^p \psi_p \\ & + (\delta_n^i \xi_{jm}^p - \delta_m^i \xi_{nj}^p) \psi_p + \xi_{jm}^p \xi_{np}^i - \xi_{jn}^p \xi_{pm}^i + 2\psi_n (L_{mj}^i + \xi_{mj}^i) \\ & + 2\psi_m (L_{nj}^i + \xi_{nj}^i) + 2\xi_{mn}^p (L_{pj}^i + \xi_{pj}^i) = 0. \end{aligned}$$

where

$$\begin{aligned} \mathcal{D}_4^i{}_{jm} = & \xi_{jm}^p \psi_p + \frac{1}{N+1} \xi_{jm}^p \xi_{qp}^q - \frac{2}{N^2-1} [2\psi_p (L_{mj}^p + \xi_{mj}^p) \\ & + \psi_m (L_{pj}^p + \xi_{pj}^p) - \psi_j (L_{pm}^p + \xi_{pm}^p) + \xi_{mq}^p L_{pj}^q - \xi_{jq}^p L_{pm}^q] \\ & + \frac{1}{N-1} [2\psi_p (L_{mj}^p + \xi_{mj}^p) + 2\psi_m (L_{pj}^p + \xi_{pj}^p) + 2\xi_{mq}^p (L_{pj}^q + \xi_{pj}^q) - \xi_{jq}^p \xi_{pm}^q] \end{aligned}$$

Definition 4.2. The space GA_N is R_4 -projectively flat if there exists an R_4 -projective mapping of the space GA_N into a flat space.

For curvature tensors of the fourth kind we get

$$(4.2) \quad \begin{aligned} \bar{R}_{4jmn}^i &= R_{4jmn}^i + \delta_j^i (\psi_{mn} - \psi_{nm}) + \delta_m^i \psi_{jn} - \delta_n^i \psi_{jm} \\ &+ \psi_p (\delta_n^i \xi_{jm}^p - \delta_m^i \xi_{nj}^p) + \xi_{jm|n}^i - \xi_{nj|m}^i + \xi_{jm}^p \xi_{np}^i - \xi_{jn}^p \xi_{pm}^i \\ &+ 2\psi_n (L_{mj}^i + \xi_{mj}^i) + 2\psi_m (L_{nj}^i + \xi_{nj}^i) + 2\xi_{mn}^p (L_{pj}^i + \xi_{pj}^i). \end{aligned}$$

Now analogously to the previous cases, we get an invariant of an R_4 -projective mapping in the form

$$\begin{aligned} \mathcal{W}(R_4)^i{}_{jmn} &= R_{4jmn}^i + \frac{1}{N+1} \delta_j^i R_{4[mn]} + \frac{1}{N^2-1} [(NR_{4jn} + R_{4nj}) \delta_m^i \\ &- (NR_{4jm} + R_{4mj}) \delta_n^i] + \frac{2}{N^2-1} \delta_j^i (2L_{\sqrt{2}n|p}^p - L_{\sqrt{2}n|p}^p + L_{\sqrt{2}n|p}^p \\ &- L_{\sqrt{2}n|p}^p) + \frac{1}{N-1} \delta_j^i (L_{\sqrt{2}m|n}^p - L_{\sqrt{2}m|n}^p) + \frac{1}{N^2-1} \delta_m^i (2L_{\sqrt{2}jn|p}^p \\ &- L_{\sqrt{2}jn|p}^p + L_{\sqrt{2}jn|p}^p) - \frac{1}{N-1} \delta_m^i (L_{\sqrt{2}jn|p}^p - L_{\sqrt{2}jn|p}^p) - \frac{1}{N^2-1} \delta_n^i (2L_{\sqrt{2}jm|p}^p \\ &- L_{\sqrt{2}jm|p}^p + L_{\sqrt{2}jm|p}^p) + \frac{1}{N-1} \delta_n^i (L_{\sqrt{2}jm|p}^p - L_{\sqrt{2}jm|p}^p) - L_{\sqrt{2}jm|p}^i + L_{\sqrt{2}jn|p}^i \end{aligned}$$

i.e. the next theorems hold:

Theorem 4.1. The tensor $\mathcal{W}(R_4)^i{}_{jmn}$ is an invariant of an R_4 -projective mapping.

Theorem 4.2. If GA_N is R_4 -projectively flat then

$$\mathcal{W}(R_4)^i{}_{jmn} = 0.$$

5. R_5 -projective mappings

Definition 5.1. A geodesic mapping $f : GA_N \rightarrow \overline{GA}_N$ is R_5 -projective if the following condition is satisfied

$$(5.1) \quad \begin{aligned} &\frac{2}{N+1} \delta_j^i \xi_{mn}^p \xi_{pq}^q + \frac{1}{N+1} \delta_m^i \xi_{jn}^p \xi_{pq}^q - \frac{1}{N+1} \delta_n^i \xi_{jm}^p \xi_{pq}^q \\ &- \frac{1}{N+1} \delta_m^i \xi_{jq}^p \xi_{np}^q + \frac{1}{N-1} \delta_n^i \xi_{jq}^p \xi_{mp}^q + \xi_{jm}^p \xi_{pn}^i - \xi_{jn}^p \xi_{mp}^i = 0 \end{aligned}$$

Definition 5.2. The space GA_N is R_5 -projectively flat if there exists an R_5 -projective mapping of the space GA_N into a flat space.

For curvature tensors of the fifth kind (0.9) of the spaces GA_N and $G\bar{A}_N$ we find the relation

$$(5.2) \quad \begin{aligned} \bar{R}_{5jmn}^i &= R_{5jmn}^i + \frac{1}{2}(\psi_{1mn} - \psi_{2nm} + \psi_{2mn} - \psi_{1nm}) \\ &+ \frac{1}{2}\delta_m^i(\psi_{1jn} + \psi_{2jn}) - \frac{1}{2}\delta_n^i(\psi_{1jm} + \psi_{2jm}) + \frac{1}{2}(\xi_{jm|n}^i - \xi_{jn|m}^i) \\ &+ \xi_{mj|n}^i - \xi_{nj|m}^i + \xi_{jm}^p \xi_{pn}^i - \xi_{jn}^p \xi_{mp}^i + \xi_{mj}^p \xi_{np}^i - \xi_{nj}^p \xi_{pm}^i. \end{aligned}$$

Putting

$$\psi_{12mn} = \frac{1}{2}(\psi_{1mn} + \psi_{2mn})$$

we get from (5.2)

$$(5.3) \quad \begin{aligned} \bar{R}_{5jmn}^i &= R_{5jmn}^i + (\psi_{12mn} - \psi_{12}) + \delta_m^i \psi_{12jn} - \delta_n^i \psi_{12jm} + \frac{1}{2}(\xi_{jm|n}^i - \xi_{jn|m}^i) \\ &+ \xi_{mj|n}^i - \xi_{nj|m}^i + \xi_{jm}^p \xi_{pn}^i - \xi_{jn}^p \xi_{mp}^i + \xi_{mj}^p \xi_{np}^i - \xi_{nj}^p \xi_{pm}^i, \end{aligned}$$

Eliminating ψ_{12mn} from (5.3), analogously to the previous cases, we get

$$(5.4) \quad \bar{W}(R_5)^i{}_{jmn} = W(R_5)^i{}_{jmn}$$

where

$$\begin{aligned} W(R_5)^i{}_{jmn} &= R_{5jmn}^i + \frac{1}{N+1}\delta_j^i R_{5[mn]} + \frac{1}{N^2-1}[(NR_5^{jn} + R_{5nj})\delta_m^i \\ &- (NR_5^{jm} + R_{5mj})\delta_n^i] - \frac{1}{N+1}\delta_j^i (L_{\sqrt{v}_3}^p{}_{mn|p} + L_{\sqrt{v}_4}^p{}_{nm|p}) \\ &- \frac{1}{2(N+1)}\delta_j^i (L_{\sqrt{v}_4}^p{}_{np|m} - L_{\sqrt{v}_4}^p{}_{mp|n} + L_{\sqrt{v}_3}^p{}_{pn|m} - L_{\sqrt{v}_3}^p{}_{pm|n}) \\ &- \frac{1}{2(N+1)}\delta_m^i (L_{\sqrt{v}_3}^p{}_{jn|p} + L_{\sqrt{v}_4}^p{}_{nj|p}) + \frac{1}{2(N^2-1)}\delta_m^i (L_{\sqrt{v}_4}^p{}_{np|j} + L_{\sqrt{v}_3}^p{}_{pn|j}) \\ &+ \frac{N}{2(N^2-1)}\delta_m^i (L_{\sqrt{v}_4}^p{}_{jp|n} + L_{\sqrt{v}_3}^p{}_{pj|n}) + \frac{1}{2(N+1)}\delta_n^i (L_{\sqrt{v}_3}^p{}_{jm|p} + L_{\sqrt{v}_4}^p{}_{mj|p}) \\ &- \frac{1}{2(N^2-1)}\delta_n^i (L_{\sqrt{v}_4}^p{}_{mp|j} + L_{\sqrt{v}_3}^p{}_{pm|j}) - \frac{N}{2(N^2-1)}\delta_n^i (L_{\sqrt{v}_4}^p{}_{jp|m} + L_{\sqrt{v}_3}^p{}_{pj|m}) \\ &- \frac{1}{2}(L_{\sqrt{v}_3}^i{}_{jm|n} - L_{\sqrt{v}_4}^i{}_{jn|m} + L_{\sqrt{v}_4}^i{}_{mj|n} - L_{\sqrt{v}_3}^i{}_{nj|m}). \end{aligned}$$

Hence:

Theorem 5.1. *The tensor $\mathcal{W}(R_5^i)_{jmn}$ is an invariant of an R_5 -projective mapping.*

Theorem 5.2. *If GA_N is R_5 -projectively flat then*

$$\mathcal{W}(R_5^i)_{jmn} = 0.$$

In the case of generalized Riemannian space (GR_N) the connection coefficients are defined by means of a non-symmetric basic tensor [1]-[3], [7] and they are non-symmetric too. The tensors $W(R_\theta^i)_{jmn}$ [9] obtained as invariants of a map $f : GR_N \rightarrow G\bar{R}_N$ are particular cases of obtained here tensors $\mathcal{W}(R_\theta^i)_{jmn}$ ($\theta = 1, \dots, 5$). For example

$$\begin{aligned} W(R_1^i)_{jmn} &= R_1^i{}_{jmn} + \frac{1}{N+1} \delta_j^i R_{1[mn]} + \frac{1}{N^2-1} [(NR_1{}_{jn} + R_{nj})\delta_m^i \\ &\quad - (NR_1{}_{jm} + R_{mj})\delta_n^i] - \frac{2}{N+1} \delta_j^i L_{m\nu_1}^p - \frac{1}{N+1} \delta_m^i L_{j\nu_1}^p \\ &\quad + \frac{1}{N+1} \delta_n^i L_{j\nu_1}^p - L_{j\nu_1}^i{}_{m} + L_{j\nu_1}^i{}_{m}. \end{aligned}$$

When GA_N (GR_N) reduces to the Riemannian space, the magnitudes $\mathcal{W}(R_\theta^i)$ ($W(R_\theta^i)$) ($\theta = 1, \dots, 5$) reduce to the Weil's tensor [8]

$$W^i{}_{jmn} = R^i{}_{jmn} + \frac{1}{N-1} (\delta_m^i R_{jn} - \delta_n^i R_{jm}).$$

References

- [1] L.P. Eisenhart, *Generalized Riemannian spaces I*, Proc. Nat. Acad. Sci. USA **37** (1951), 311-315.
- [2] S.M. Minčić, *Ricci identities in the space of non-symmetric affine connection*, Mat. Vesnik **10**(25) (1973), 161-172.
- [3] S.M. Minčić, *New commutation formulas in the non-symmetric affine connection space*, Publ. Inst. Math. (Beograd) (N. S) **22**(36) (1977), 189-199.
- [4] S.M. Minčić, *Independent curvature tensors and pseudotensors of spaces with non-symmetric affine connection*, Coll. Math. Soc. János Bolyai **31** (1979), 445-460.
- [5] S.M. Minčić and M.S. Stanković, *On geodesic mappings of general affine connection spaces and of generalized Riemannian spaces*, Mat. Vesnik **49** (1997), 27-33.
- [6] S.M. Minčić and M.S. Stanković, *Equitortion geodesic mappings of generalized Riemannian spaces*, Publ. Inst. Math. (Beograd) (NS) **61**(75) (1997), 97-104.

- [7] K.D. Singh, *On generalized Riemann spaces*, Riv. Mat. Univ. Parma 7 (1956), 125–138.
- [8] N.S. Sinyukov, *Geodesic mappings of Riemannian spaces*, Nauka, Moskva, 1979.
- [9] M.S. Stanković and S.M. Minčić, *New special geodesic mappings of generalized Riemannian spaces*, Publ. Inst. Math. (Beograd) 67(81) (2000), 92–102.

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