

ON A CANONIC ALMOST GEODESIC MAPPINGS OF THE SECOND TYPE OF AFFINE SPACES

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Abstract. In this paper we investigate two kinds of almost geodesic mappings of the second type of affine spaces. Also we find some invariant geometric objects for the canonic almost geodesic mappings.

0. Introduction

Let GA_N be an N -dimensional space with an affine connection L given with the aid of components L_{jk}^i in each local map V on a differentiable manifold. Generally it is $L_{jk}^i \neq L_{kj}^i$.

Generalizing conception of a geodesic mappings for Riemannian and affine spaces Sinyukov introduced [5] following notations:

The curve $l : x^h = x^h(t)$ is called the almost geodesic line if its tangential vector $\lambda^h(t) = dx^h/dt \neq 0$ satisfies the equations

$$\bar{\lambda}_{(2)}^h = \bar{a}(t)\lambda^h + \bar{b}(t)\bar{\lambda}_{(1)}^h, \quad \bar{\lambda}_{(1)}^h = \lambda_{||\alpha}^h \lambda^\alpha, \quad \bar{\lambda}_{(2)}^h = \bar{\lambda}_{(1)||\alpha}^h \lambda^\alpha,$$

where $\bar{a}(t)$ and $\bar{b}(t)$ are functions of a parameter t , and $||$ denotes the covariant derivative with respect to the connection in \bar{A}_N .

A mapping f of the affine space A_N onto a space \bar{A}_N is called the almost geodesic mapping if any geodesic line of the space A_N turns into almost geodesic line of the space \bar{A}_N .

Sinjukov [5] singled out the three types of the almost geodesic mappings, π_1, π_2, π_3 for spaces without torsion. In the present work we investigate the

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mappings of the type π_2 for spaces with torsion. In a differentiable manifold with nonsymmetric affine connection L_{jk}^i , for a vector exist two kinds of covariant derivative:

$$\lambda_{1m}^h = \lambda_{,m}^h + L_{pm}^h \lambda^p, \quad \lambda_{2m}^h = \lambda_{,m}^h + L_{mp}^h \lambda^p.$$

Thus, in the case of the space with nonsymmetric affine connection we can define two kinds of almost geodesic lines and two kinds of almost geodesic mappings.

1. Almost geodesic mappings of affine spaces

In an affine spaces GA_N (with nonsymmetric affine connection L) one can define four kinds of a covariant derivative [1,2]. Let us denote by $\bar{\lambda}_{\theta}^h$ a covariant derivative of the kind θ in GA_N and $G\bar{A}_N$ respectively. A curve $l: x^h = x^h(t)$ is called *almost geodesic line of the first kind* if its tangential vector $\lambda^h(t) = dx^h/dt \neq 0$ satisfies the equations

$$(1.1) \quad \bar{\lambda}_{1(2)}^h = \bar{a}_1(t) \lambda^h + \bar{b}_1(t) \bar{\lambda}_{1(1)}^h, \quad \bar{\lambda}_{1(1)}^h = \lambda^h|_{1\alpha} \lambda^\alpha, \quad \bar{\lambda}_{1(2)}^h = \bar{\lambda}_{1(1)}^h|_{1\alpha} \lambda^\alpha,$$

where $\bar{a}_1(t)$ and $\bar{b}_1(t)$ are functions of a parameter t . A curve $l: x^h = x^h(t)$ is called *almost geodesic line of the second kind* if its tangential vector $\lambda^h(t) = dx^h/dt \neq 0$ satisfies the equations

$$(1.2) \quad \bar{\lambda}_{2(2)}^h = \bar{a}_2(t) \lambda^h + \bar{b}_2(t) \bar{\lambda}_{2(1)}^h, \quad \bar{\lambda}_{2(1)}^h = \lambda^h|_{2\alpha} \lambda^\alpha, \quad \bar{\lambda}_{2(2)}^h = \bar{\lambda}_{2(1)}^h|_{2\alpha} \lambda^\alpha,$$

where $\bar{a}_2(t)$ and $\bar{b}_2(t)$ are functions of a parameter t .

A mapping f of the space GA_N onto a space with nonsymmetric affine connection $G\bar{A}_N$ is called *almost geodesic mapping of the first kind* if any geodesic line of the space GA_N turns into almost geodesic line of the first kind of the space $G\bar{A}_N$. A mapping f is called *almost geodesic mapping of the second kind* if any geodesic line of the space GA_N turns into almost geodesic line of the second kind of the space $G\bar{A}_N$.

We can put

$$(1.3) \quad \bar{L}_{ij}^h(x) = L_{ij}^h(x) + P_{ij}^h(x),$$

where $L_{ij}^h(x)$, $\bar{L}_{ij}^h(x)$ are connection coefficients of the space GA_N and $G\bar{A}_N$, ($N > 2$), and $P_{ij}^h(x)$ is a deformation tensor. From [7] it follows that the following results hold:

Theorem 1.1. *The mapping f of the space GA_N onto $G\bar{A}_N$ is the almost geodesic mapping of the first kind if and only if the deformation tensor $P_{ij}^h(x)$ satisfies the conditions*

$$(1.4) \quad (P_{\alpha\beta}^h|_{\gamma} + P_{\delta\alpha}^h P_{\beta\gamma}^{\delta})\lambda^{\alpha} \lambda^{\beta} \lambda^{\gamma} = b_1 P_{\alpha\beta}^h \lambda^{\alpha} \lambda^{\beta} + a_1 \lambda^h$$

identically, where a_1 and b_1 are invariants.

Theorem 1.2. *The mapping f of the space GA_N onto $G\bar{A}_N$ is the almost geodesic mapping of the second kind if and only if the deformation tensor $P_{ij}^h(x)$ satisfies the conditions*

$$(1.4') \quad (P_{\alpha\beta}^h|_{\gamma} + P_{\alpha\delta}^h P_{\beta\gamma}^{\delta})\lambda^{\alpha} \lambda^{\beta} \lambda^{\gamma} = b_2 P_{\alpha\beta}^h \lambda^{\alpha} \lambda^{\beta} + a_2 \lambda^h$$

identically, where a_1 and b_1 are invariant.

2. Almost geodesic mappings of the second type

In [5] Sinyukov introduced almost geodesic mapping of the second type π_2 for affine spaces without torsion by condition:

$$b = \frac{b_{\gamma\delta} \lambda^{\gamma} \lambda^{\delta}}{\sigma_{\alpha} \lambda^{\alpha}}.$$

Analogously, almost geodesic mappings of the first kind of affine space is the second type π_2 if for the function b_1 satisfied the following condition:

$$b_1 = \frac{b_{\gamma\delta} \lambda^{\gamma} \lambda^{\delta}}{\sigma_{\alpha} \lambda^{\alpha}},$$

where $\sigma_{\alpha} \lambda^{\alpha} \neq 0$.

Let there be

$$(2.1) \quad P_{\alpha\beta}^h \lambda^{\alpha} \lambda^{\beta} = 2\sigma_{\alpha} \lambda^{\alpha} F_{\beta}^h \lambda^{\beta} + 2\psi_{\alpha} \lambda^{\alpha} \lambda^h.$$

Then

$$(P_{\alpha\beta}^h - 2\sigma_{\alpha} F_{\beta}^h - 2\psi_{\alpha} \delta_{\beta}^h) \lambda^{\alpha} \lambda^{\beta} \equiv 0,$$

wherefrom

$$(2.2) \quad \underline{P}_{ij}^h = \psi_i \delta_j^h + \psi_j \delta_i^h + \sigma_i F_j^h + \sigma_j F_i^h.$$

We can put $P_{ij}^h = \xi_{ij}^h$. In this paper $\underset{\vee}{ij}$ denotes symmetrization with division and $\underset{\wedge}{ij}$ denotes antisymmetrization with division. Then

$$(2.3) \quad P_{ij}^h(x) = \psi_i(x) \delta_j^h + \psi_j(x) \delta_i^h + \sigma_i(x) F_j^h(x) + \sigma_j(x) F_i^h(x) + \xi_{ij}^h(x).$$

In (2.3) ψ_i, σ_i are vectors, F_i^h is a tensor and ξ_{ij}^h is an antisymmetric tensor. Substituting (2.3) in (1.4) we obtain:

$$(2.4) \quad \begin{aligned} F_{i|j}^h + F_{j|i}^h + F_\delta^h F_i^\delta \sigma_j + F_\delta^h F_j^\delta \sigma_i + \xi_{\delta i}^h F_j^\delta + \xi_{\delta j}^h F_i^\delta \\ = \mu_i F_j^h + \mu_j F_i^h + \nu_i \delta_j^h + \nu_j \delta_i^h, \end{aligned}$$

where μ_i, ν_i are covariant vectors.

Conditions (2.3) and (2.4) are *the basic equations* of the mapping π_1 .

Almost geodesic mapping of the second kind is the second type π_2 if satisfied the following condition for the function b_2 in (1.4')

$$b_2 = \frac{b_{2\gamma\delta} \lambda^\gamma \lambda^\delta}{\sigma_\alpha \lambda^\alpha}$$

where $\sigma_\alpha \lambda^\alpha \neq 0$. Using the methods from the previous case we get

$$(2.4') \quad \begin{aligned} F_{i|j}^h + F_{j|i}^h + F_\delta^h F_i^\delta \sigma_j + F_\delta^h F_j^\delta \sigma_i + \xi_{i\delta}^h F_j^\delta + \xi_{j\delta}^h F_i^\delta \\ = \mu_i F_j^h + \mu_j F_i^h + \nu_i \delta_j^h + \nu_j \delta_i^h, \end{aligned}$$

where μ_i, ν_i are covariant vectors.

Conditions (2.3) and (2.4') are *the basic equations* of almost geodesic mappings of π_2 .

3. The property of reciprocity of almost geodesic mappings of the second type

The mapping π_1 has *the property of reciprocity*, if its inverse mapping π_1^{-1} is of π_2 type, and π_1^{-1} corresponds to the same affinor F_i^h also. Since the inverse mapping $\pi_1^{-1} : G\bar{A}_N \rightarrow GA_N$ satisfies

$$\bar{P}_{ij}^h = -P_{ij}^h,$$

we can put in (2.3) the following:

$$\bar{\psi}_i = -\psi_i, \quad \bar{\sigma}_i = -\sigma_i, \quad \bar{F}_i^h = F_i^h, \quad \bar{\xi}_{ij}^h = -\xi_{ij}^h.$$

The mapping π_2 has the property of reciprocity if and only if the affnor F_i^h of the space $G\bar{A}_N$ satisfies the equation of the form (2.4), i.e.

$$(3.1) \quad F_{i||j}^h - F_\alpha^h F_{(i}^\alpha \sigma_{j)} - \xi_{\alpha(i}^h F_{j)}^\alpha = \bar{\mu}_{(i} F_{j)}^h + \bar{\nu}_{(i} \delta_{j)}^h,$$

where (ij) is a symmetrization without division with respect to i and j , and $||$ is a covariant derivative of the first kind in $G\bar{A}_N$. Crossing in (3.1) to a covariant derivative of the first kind in GA_N , we get

$$(3.2) \quad F_\alpha^h F_{(i}^\alpha \sigma_{j)} + \xi_{\alpha(i}^h F_{j)}^\alpha = \bar{\mu}_{(i} F_{j)}^h + \bar{\nu}_{(i} \delta_{j)}^h,$$

where the vectors $\bar{\mu}_i, \bar{\nu}_i$ are expressed by $\mu_i, \nu_i, \bar{\mu}_i, \bar{\nu}_i, \psi_i, \sigma_i, F_i^h$. Since $\sigma \neq 0$, we get

$$(3.3) \quad F_\alpha^h F_i^\alpha = p\delta_i^h + qF_i^h,$$

where p i q are invariants.

On the base of the facts given above, we have:

Theorem 3.1. *Relation (3.3) expressed necessary and sufficient condition that a mapping $\pi_2 : GA_N \rightarrow G\bar{A}_N$ has the property of reciprocity.*

The equations (2.3) and (2.4) are invariant to the mapping of the affnor:

$$(3.4) \quad \tilde{F}_i^h = rF_i^h + s\delta_i^h \quad (r \neq 0).$$

Then we have

$$(3.5) \quad \tilde{F}_\alpha^h \tilde{F}_i^\alpha = \tilde{p}\delta_i^h + \tilde{q}\tilde{F}_i^h,$$

where are

$$\tilde{p} = r^2p - s^2 - srq, \quad \tilde{q} = 2s + rq.$$

Here we can select invariants r and s such that

$$\tilde{q} \equiv 0, \quad \tilde{p} = \tilde{e} (= \pm 1, 0).$$

In this case we have

$$(3.6) \quad \tilde{F}_\alpha^h \tilde{F}_i^\alpha = \tilde{\epsilon} \delta_i^h.$$

On the base of the facts given above we can put

$$(3.7) \quad F_\alpha^h F_i^\alpha = e \delta_i^h \quad (e = \pm 1, 0)$$

Using (3.7) in the condition (2.4) we get:

$$(3.8) \quad F_{(i|j)}^h + \xi_{\alpha(i}^h F_j^\alpha) = \mu_{(i} F_j^h) + \nu_{(i} \delta_j^h)$$

Hence, the almost geodesic mapping $\pi_2 : GA_N \rightarrow G\bar{A}_N$ which satisfies the property of reciprocity is characterized by equations (2.3), (3.7) and (3.8). This mapping is denoted

$$\pi_2(e) : GA_N \rightarrow G\bar{A}_N.$$

The case of the mappings π_2 can be investigated analogously.

4. Canonic almost geodesic mappings

The mapping $\pi_1(e) : GA_N \rightarrow G\tilde{A}_N$ is *canonic* if in by mapping common coordinate system between connection coefficients L_{ij}^h and \tilde{L}_{ij}^h satisfies the relation

$$(4.1) \quad \tilde{L}_{ij}^h = L_{ij}^h + \sigma_i F_j^h + \sigma_j F_i^h + \xi_{ij}^h,$$

where ξ_{ij}^h is an antisymmetric tensor, σ_i a vector, F_i^h is an affinor, and F_i^h satisfies the conditions (3.7) and (3.8).

Theorem 4.1. *Every mapping π_1 is canonic or it can be present in form of the product of a geodesic and a canonic mapping.*

Proof. If $\psi_i \equiv 0$ then this mapping is canonic. Let $\psi_i \neq 0$. Suppose that for the mapping $f : GA_N \rightarrow G\bar{A}_N$ satisfies

$$\bar{L}_{ij}^h = L_{ij}^h + \sigma_{(i} F_j^h) + \bar{\xi}_{ij}^h,$$

and for mapping $g : G\bar{A}_N \rightarrow G\tilde{A}_N$ satisfies

$$\tilde{L}_{ij}^h = \bar{L}_{ij}^h + \psi_{(i} \delta_j^h) + \tilde{\xi}_{ij}^h.$$

Then for the product of the mappings f and g we have

$$\tilde{L}_{ij}^h = L_{ij}^h + \psi_{(i}\delta_{j)}^h + \sigma_{(i}F_{j)}^h + \bar{\xi}_{ij}^h + \tilde{\xi}_{ij}^h.$$

i.e. the mapping π_2 is presented as the product of the geodesic f and the canonic mapping g . \square

Theorem 4.2. *Geometric objects of the space GA_N*

$$(4.2) \quad T_{ij}^h = L_{ij}^h + \frac{1}{e - F^2} [(FL_{\alpha j}^\alpha - F_j^\alpha L_{\beta\alpha}^\beta)F_i^h + (FL_{\alpha i}^\alpha - F_i^\alpha L_{\beta\alpha}^\beta)F_j^h],$$

are invariant of the canonic mapping $\pi_2(e)$ ($e = \pm 1$), where F_i^h is an affnor, such that $F = F_\alpha^\alpha$, $e - F^2 \neq 0$ and ij is a symmetrization with division.

Proof. Let us contract in (4.1) with respect to h, i :

$$(4.3) \quad \tilde{L}_{\alpha j}^\alpha = L_{\alpha j}^\alpha + F\tilde{\sigma}_j + \xi_{\alpha j}^\alpha$$

where

$$(4.4) \quad \tilde{\sigma}_i = \sigma_\alpha F_i^\alpha, \quad F = F_\alpha^\alpha.$$

From (3.7) and (4.4) we have

$$(4.5) \quad F_j^\beta (\tilde{L}_{\alpha\beta}^\alpha - L_{\alpha\beta}^\alpha) = e\sigma_j + F\tilde{\sigma}_j.$$

Substituting (4.3) in (4.5) we get

$$(4.6) \quad \sigma_j = \frac{1}{e - F^2} [F_j^\beta (\tilde{L}_{\alpha\beta}^\alpha - L_{\alpha\beta}^\alpha) - F(\tilde{L}_{\alpha j}^\alpha - L_{\alpha j}^\alpha)].$$

From (4.6) and (4.1) we have

$$\begin{aligned} \tilde{L}_{ij}^h &+ \frac{1}{e - F^2} [(F\tilde{L}_{\alpha j}^\alpha - F_j^\alpha \tilde{L}_{\beta\alpha}^\beta)F_i^h + (F\tilde{L}_{\alpha i}^\alpha - F_i^\alpha \tilde{L}_{\beta\alpha}^\beta)F_j^h] \\ &= L_{ij}^h + \frac{1}{e - F^2} [(FL_{\alpha j}^\alpha - F_j^\alpha L_{\beta\alpha}^\beta)F_i^h + (FL_{\alpha i}^\alpha - F_i^\alpha L_{\beta\alpha}^\beta)F_j^h], \end{aligned}$$

i.e.

$$(4.7) \quad \tilde{T}_{ij}^h(x) = T_{ij}^h(x).$$

Theorem is proved. \square

From (4.1) and (3.7) we get

$$(4.8) \quad F_{i|j}^h = F_{i|j}^h + \tilde{\sigma}_i F_j^h - e \sigma_i \delta_j^h + \xi_{\alpha j}^h F_i^\alpha - \xi_{ij}^\alpha F_\alpha^h,$$

where $\underset{1}{|}$ is a covariant derivative of the first kind in a space $G\hat{A}_N$. From here we get

$$(4.9) \quad F_\alpha^h (F_{i|j}^\alpha - F_{i|j}^\alpha) = e(\tilde{\sigma}_i \delta_j^h - \sigma_i F_j^h) + \xi_{\beta j}^\alpha F_i^\beta F_\alpha^h - e \xi_{ij}^h,$$

i.e.

$$(4.10) \quad \sigma_i F_j^h + \sigma_j F_i^h = e F_\alpha^h (F_{i|j}^\alpha - F_{i|j}^\alpha) + \tilde{\sigma}_{(i} \delta_{j)}^h + e \xi_{\beta(j}^\alpha F_{i)}^\beta F_\alpha^h$$

Substituting (4.10) in (4.1) we get

$$(4.11) \quad \begin{aligned} \tilde{L}_{ij}^h + e F_\alpha^h F_{(i|j)}^\alpha - e \tilde{L}_{\beta(j}^\alpha F_{i)}^\beta F_\alpha^h \\ = L_{ij}^h + e F_\alpha^h F_{(i|j)}^\alpha - e L_{\beta(j}^\alpha F_{i)}^\beta F_\alpha^h + \tilde{\sigma}_{(i} \delta_{j)}^h. \end{aligned}$$

We can put the following relations in (4.11):

$$(4.12) \quad \begin{aligned} \hat{L}_{ij}^h &= \tilde{L}_{ij}^h + e F_\alpha^h F_{(i|j)}^\alpha - e \tilde{L}_{\beta(j}^\alpha F_{i)}^\beta F_\alpha^h, \\ \hat{L}_{ij}^h &= L_{ij}^h + e F_\alpha^h F_{(i|j)}^\alpha - e L_{\beta(j}^\alpha F_{i)}^\beta F_\alpha^h. \end{aligned}$$

Then

$$(4.13) \quad \hat{\tilde{L}}_{ij}^h = \hat{L}_{ij}^h + \tilde{\sigma}_{(i} \delta_{j)}^h.$$

We can see that the magnitudes $\hat{\tilde{L}}_{ij}^h$ and \hat{L}_{ij}^h are connection objects of the spaces $\hat{\tilde{A}}_N$ and \hat{A}_N without torsion [5]. From the condition (4.13) we can conclude that the space \hat{A}_N geodesic maps to the space $\hat{\tilde{A}}_N$. Then projective parameters of this spaces are equal, [3] i.e.

$$(4.14) \quad \hat{\tilde{T}}_{ij}^h(x) = \hat{T}_{ij}^h(x).$$

They order geometric objects of the spaces GA_N and $G\bar{A}_N$, which are invariant to the canonic almost geodesic mappings. Then we have

$$\begin{aligned} \hat{T}_{ij}^h(x) &= \hat{L}_{ij}^h - \frac{1}{1+N}(\hat{L}_{i\alpha}^\alpha \delta_j^h + \hat{L}_{j\alpha}^\alpha \delta_i^h) \\ &= L_{ij}^h + eF_\alpha^h F_{i|j}^\alpha - eL_{\beta(j}^\alpha F_i^\beta F_\alpha^h \\ &\quad - \frac{1}{1+N} [(L_{i\alpha}^\alpha + eF_\alpha^\beta F_{\beta|j}^\alpha - eL_{\beta(j}^\alpha F_i^\beta F_\alpha^\gamma) \delta_j^h \\ &\quad + (L_{j\alpha}^\alpha + eF_\alpha^\beta F_{\beta|i}^\alpha - eL_{\beta(j}^\alpha F_i^\beta F_\alpha^\gamma) \delta_i^h] \\ &= T_{ij}^h + eF_\alpha^h (F_{i|j}^\alpha - L_{\beta(j}^\alpha F_i^\beta) \\ &\quad - \frac{e}{1+N} F_\alpha^\beta [F_{\beta|j}^\alpha - L_{\gamma(\beta}^\alpha F_i^\gamma) \delta_j^h + (F_{\beta|i}^\alpha - L_{\gamma(\beta}^\alpha F_j^\gamma) \delta_i^h], \end{aligned}$$

i.e.

$$(4.15) \quad \begin{aligned} \hat{T}_{ij}^h(x) &= T_{ij}^h + eF_\alpha^h (F_{i|j}^\alpha - L_{\beta(j}^\alpha F_i^\beta) \\ &\quad - \frac{e}{1+N} F_\alpha^\beta [F_{\beta|j}^\alpha - L_{\gamma(\beta}^\alpha F_i^\gamma) \delta_j^h + (F_{\beta|i}^\alpha - L_{\gamma(\beta}^\alpha F_j^\gamma) \delta_i^h], \end{aligned}$$

where T_{ij}^h are objects of the projective connection [3]. On the base of the facts given above, we have

Theorem 4.3. *Geometric objects of the space GA_N (4.15) are invariant of the canonic mapping $\pi_2(e)$ ($e = \pm 1$) where F_i^h is an affinor, $F = F_\alpha^\alpha$ and ij is an antisymmetrization with division.*

Remark. From (4.13) it follows that the space \hat{A}_N geodesic maps to the space \tilde{A}_N . Also from (4.12) we can see that the spaces \hat{A}_N and \tilde{A}_N are with symmetric affine connection.

Now, for example the Weil's tensor [4,6] is an invariant of this mapping, i.e.

$$(4.16) \quad \hat{W}_2^i{}_{jmn}(x) = \tilde{W}_2^i{}_{jmn}(x)$$

where

$$(4.17) \quad \begin{aligned} \hat{W}_2^i{}_{jmn} &= \hat{R}^i{}_{jmn} + \frac{1}{1+N} \delta_j^i \hat{R}_{[mn]} + \frac{1}{N^2-1} [(N \hat{R}_{jn} + \hat{R}_{nj}) \\ &\quad - (N \hat{R}_{jm} + \hat{R}_{mj})], \end{aligned}$$

\hat{R}^i_{jmn} is a curvature tensor and \hat{R}_{jm} is a Ricci tensor of the space \hat{A}_N .

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